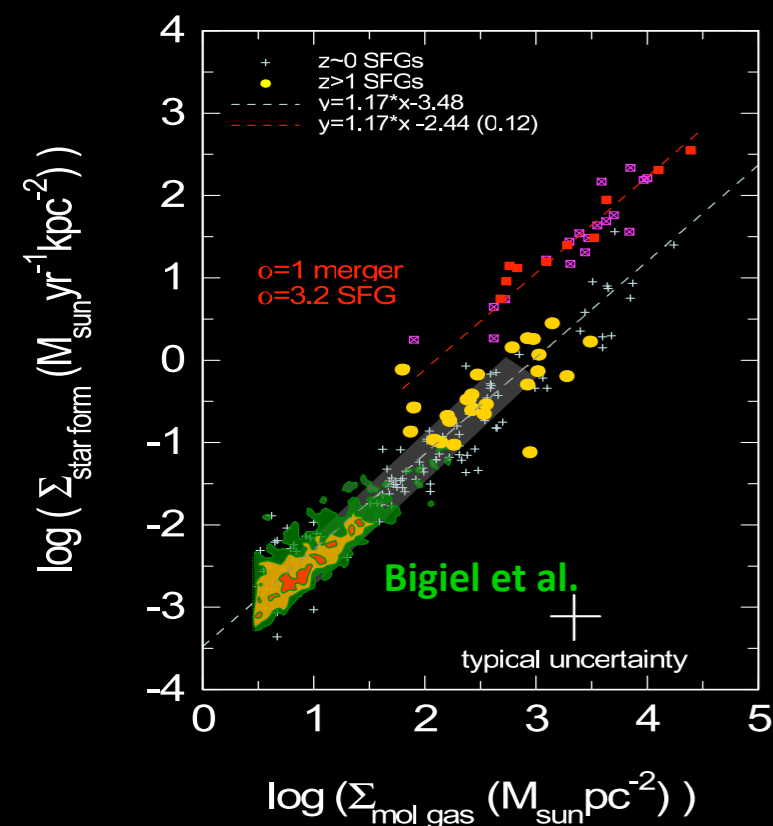
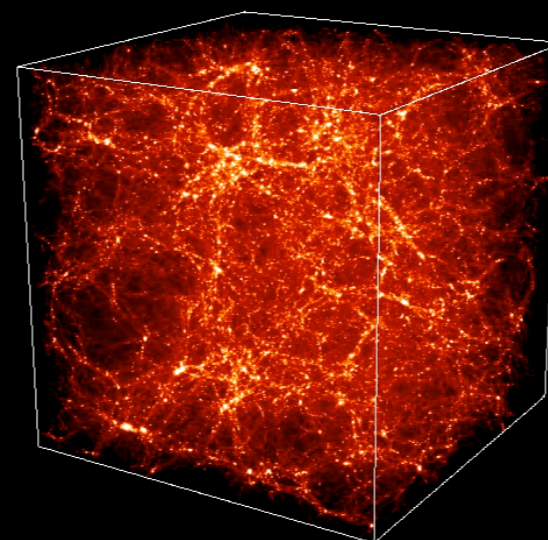
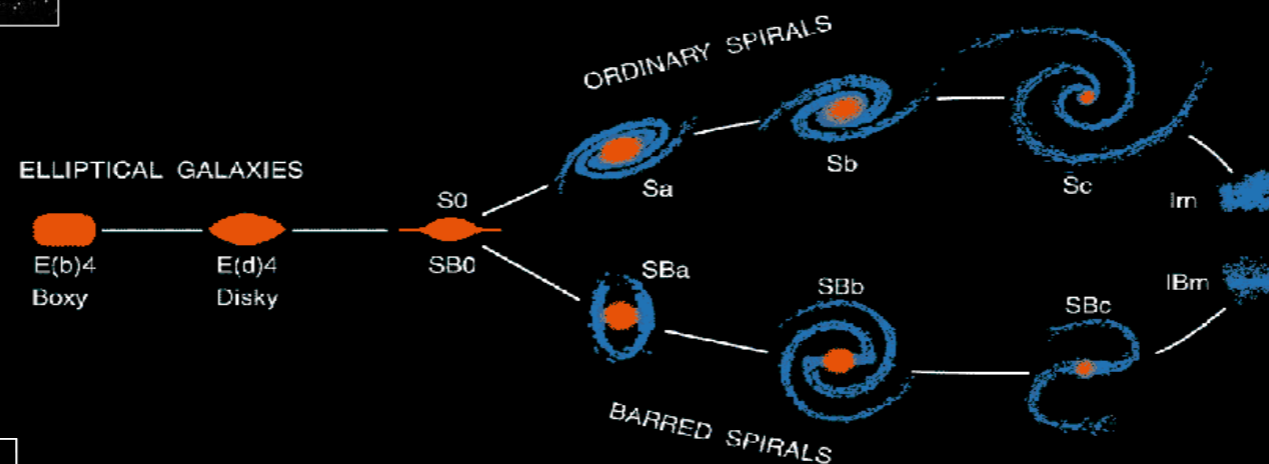


Galaxy formation modes and their relation to structure

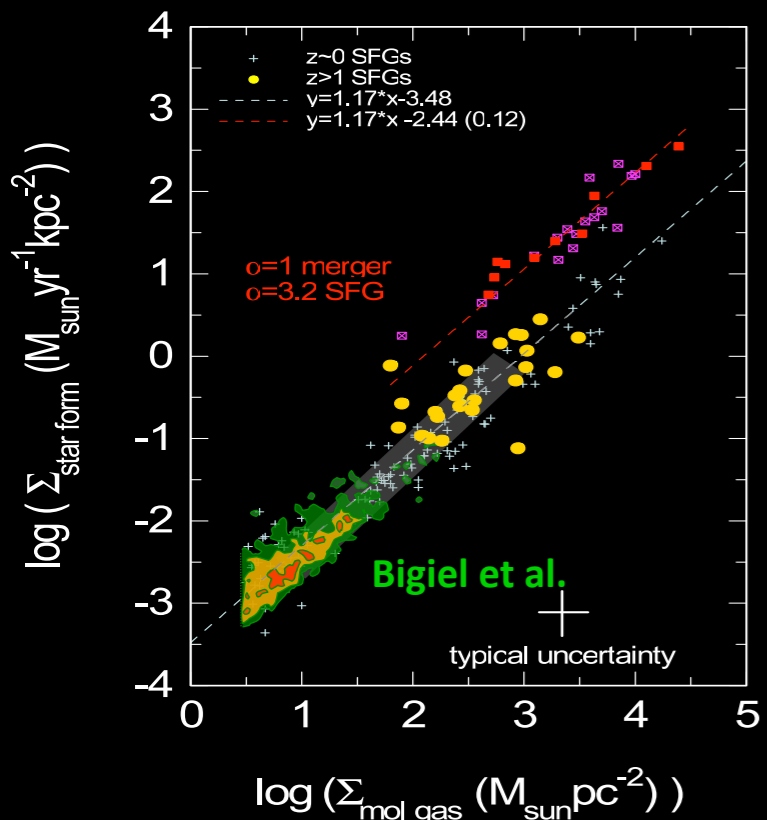
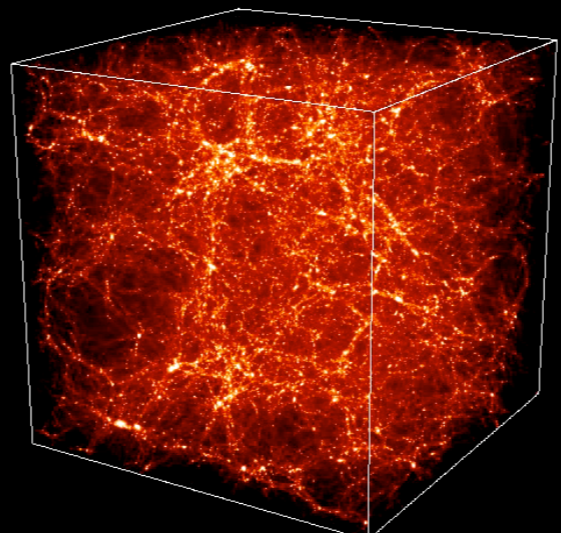
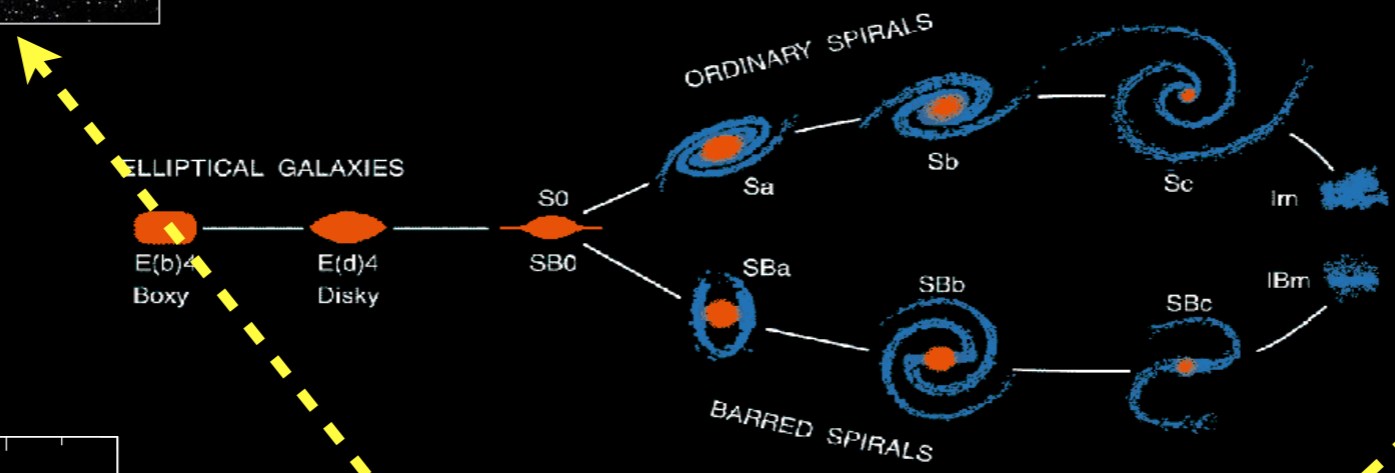
Andreas Burkert
CAST, University of Munich



Galaxy formation modes and their relation to structure

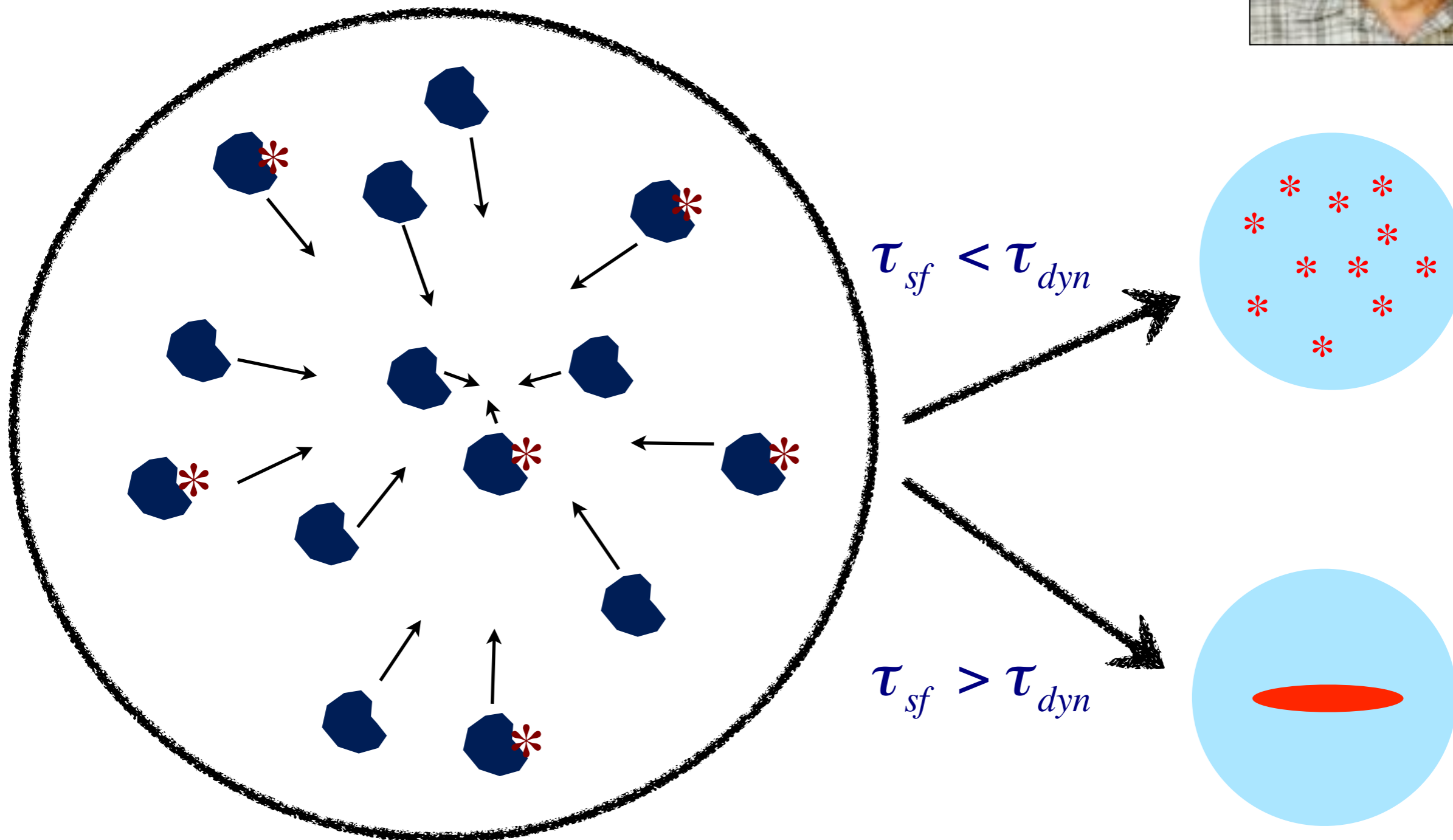
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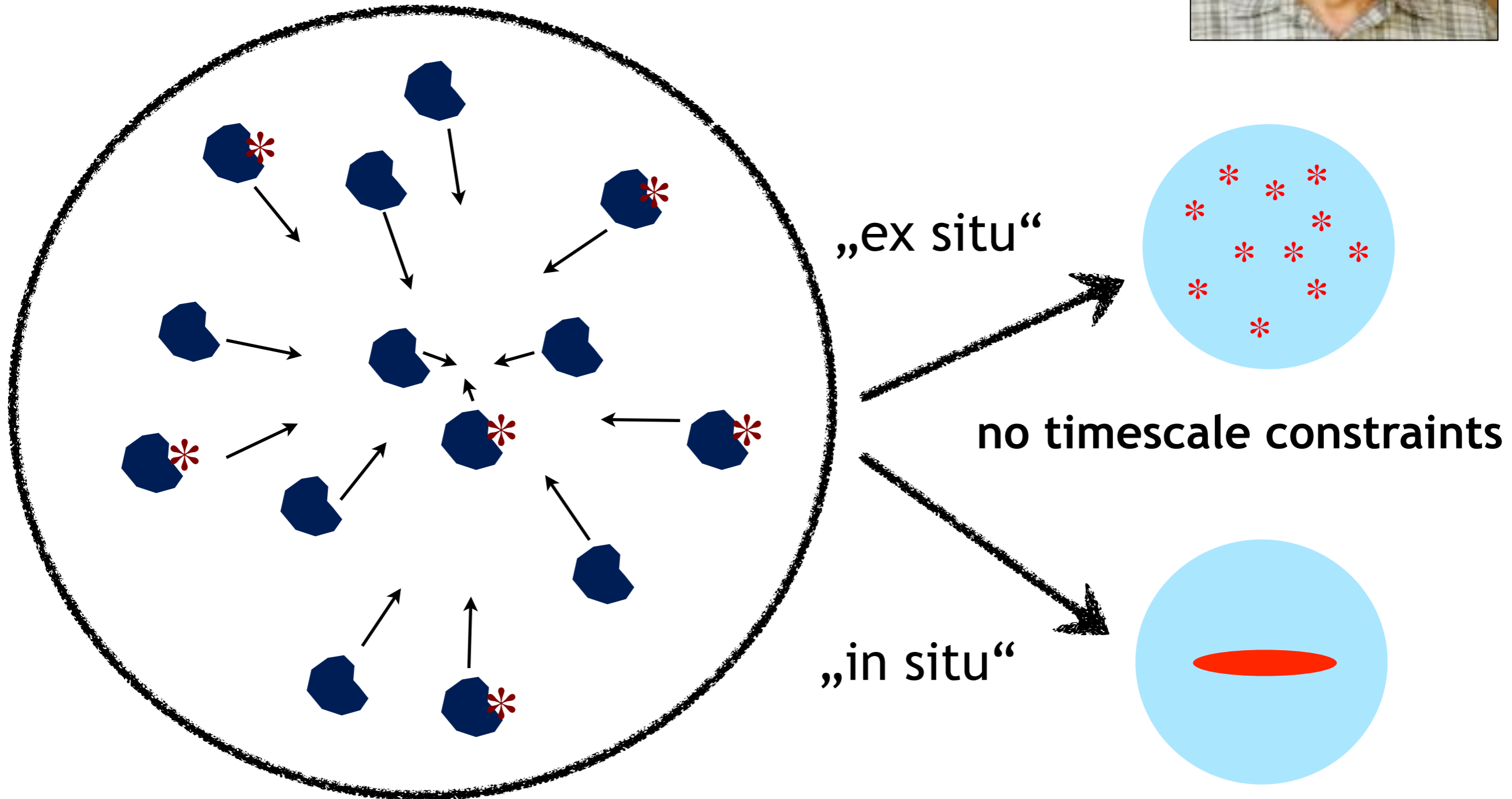
The Larson models

(Larson 69, 74, 75, 76)

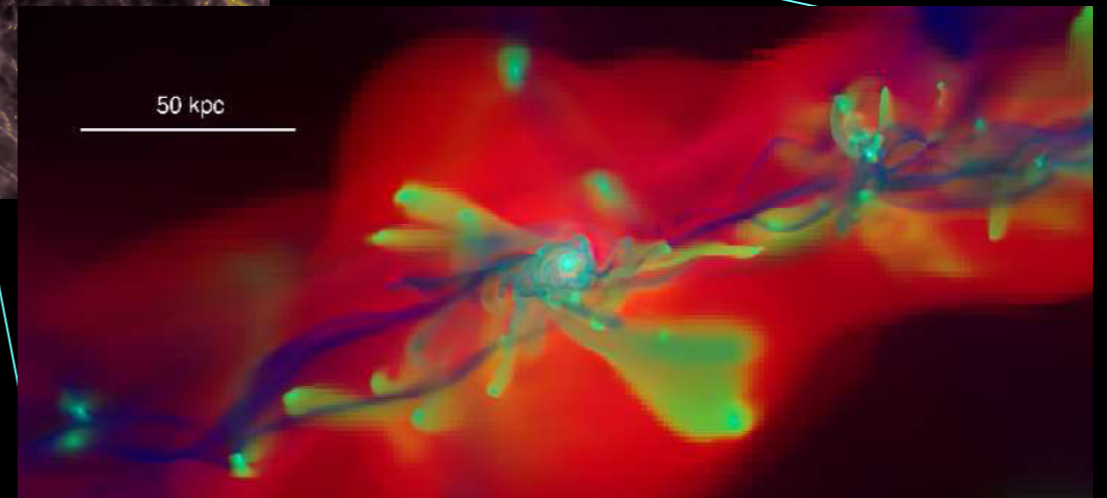
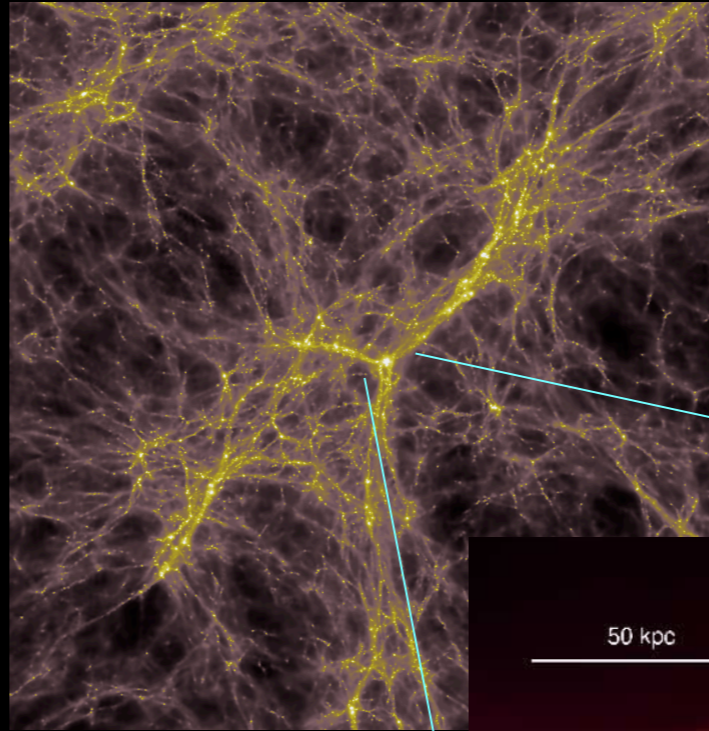
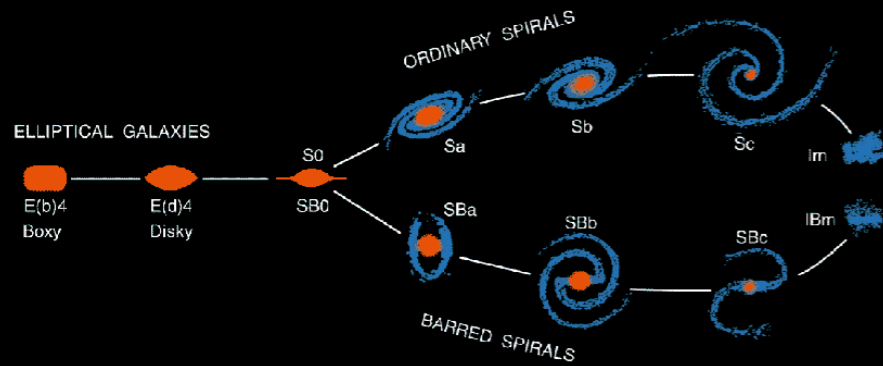


The Larson models

(Larson 69, 74, 75, 76)



Origin of the Hubble Sequence

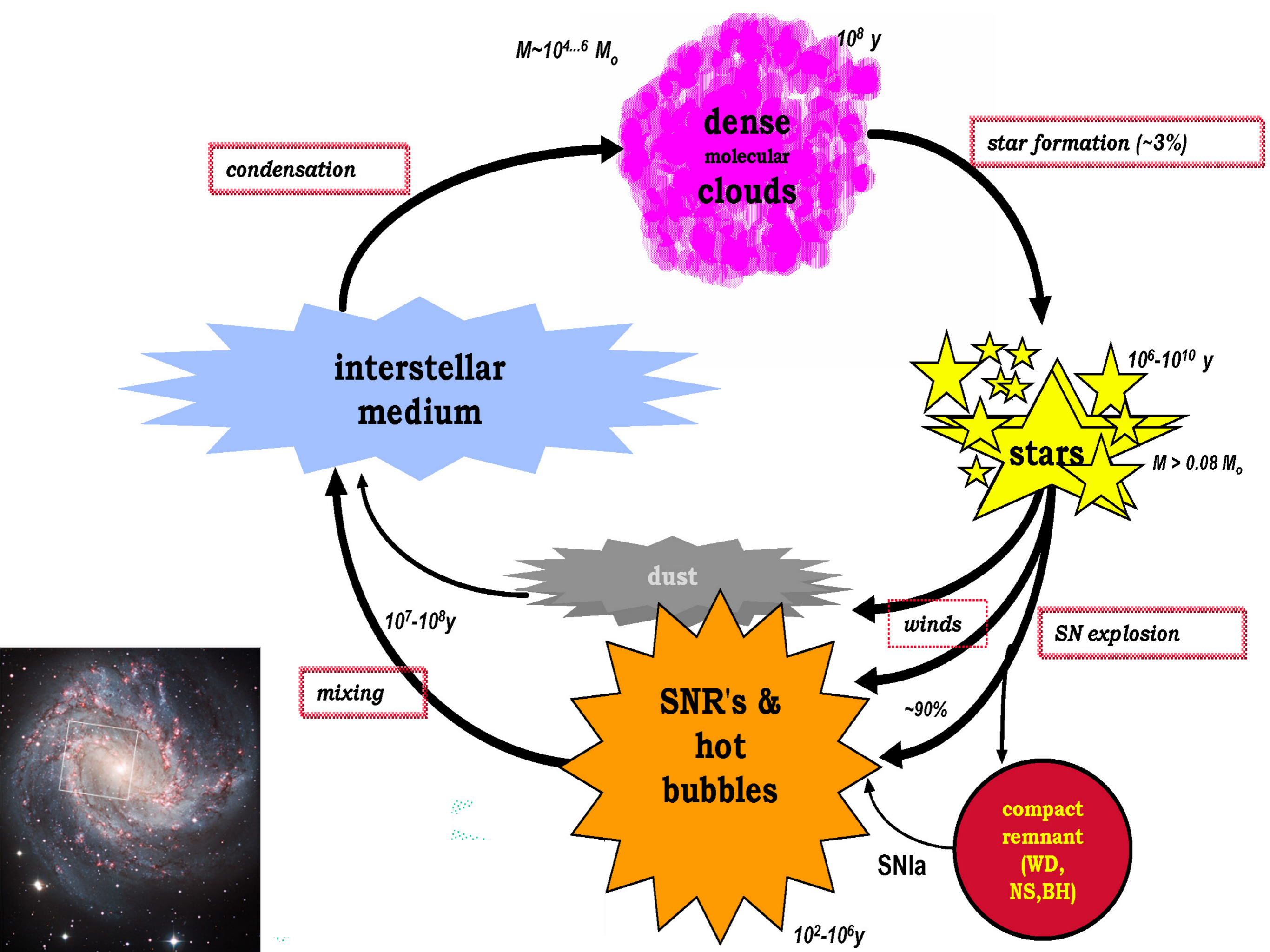


Gas infall forms disk galaxies

Major mergers form ellipticals

Toomre & Toomre 1972, Hernquist 1989-2011, Springel, Hopkins et al. 2003-2011, Robertson & Bullock 2008, Naab et al. 2003-2011

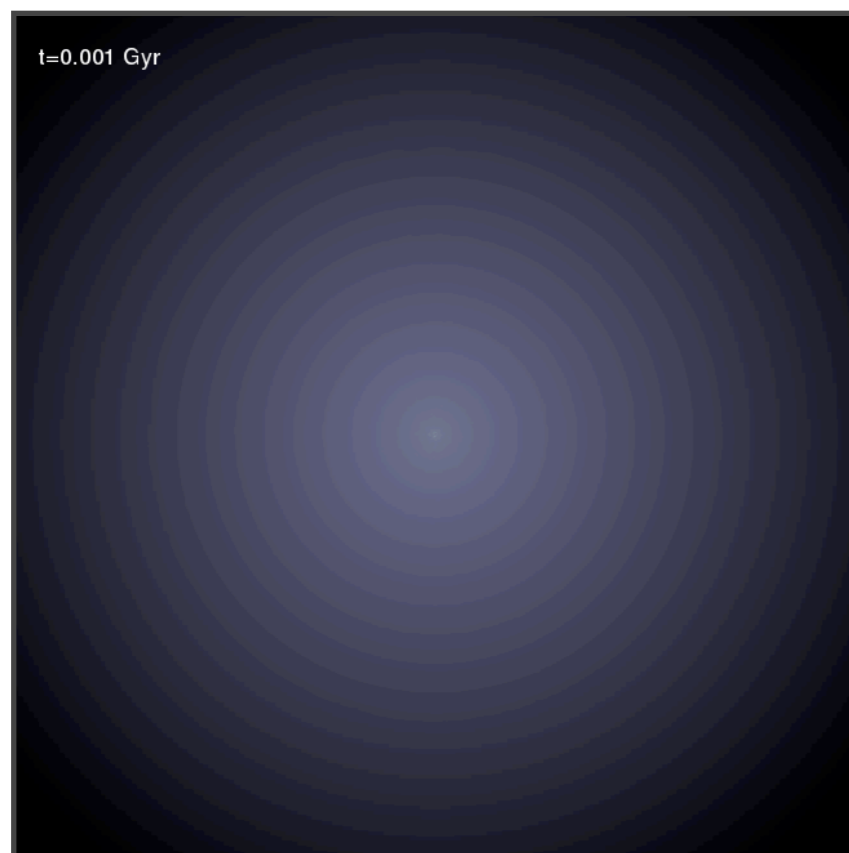
Dekel & Birnboim 2003,2006,2009, Keres et al. 2005, 2009, Davé 2007, Dekel et al. 2009, Agertz et al. 2009, Ceverino et al. 2010, Genel et al. 2010



Dobbs, Burkert & Pringle 11a,b, 12a,b

Feedback puffs up disk

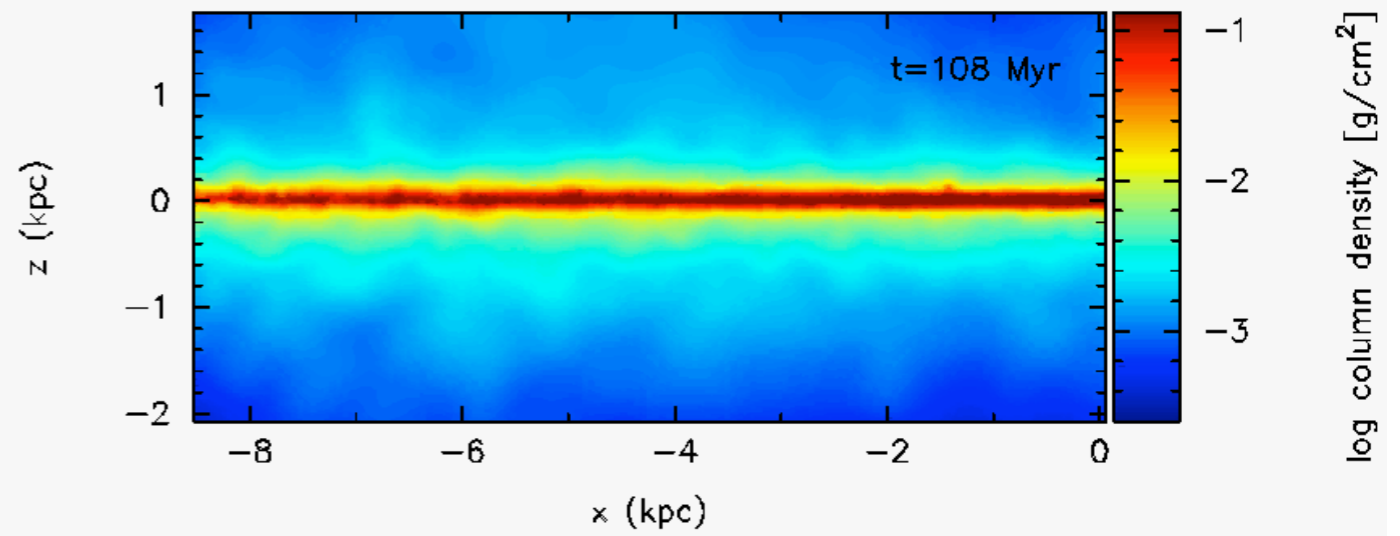
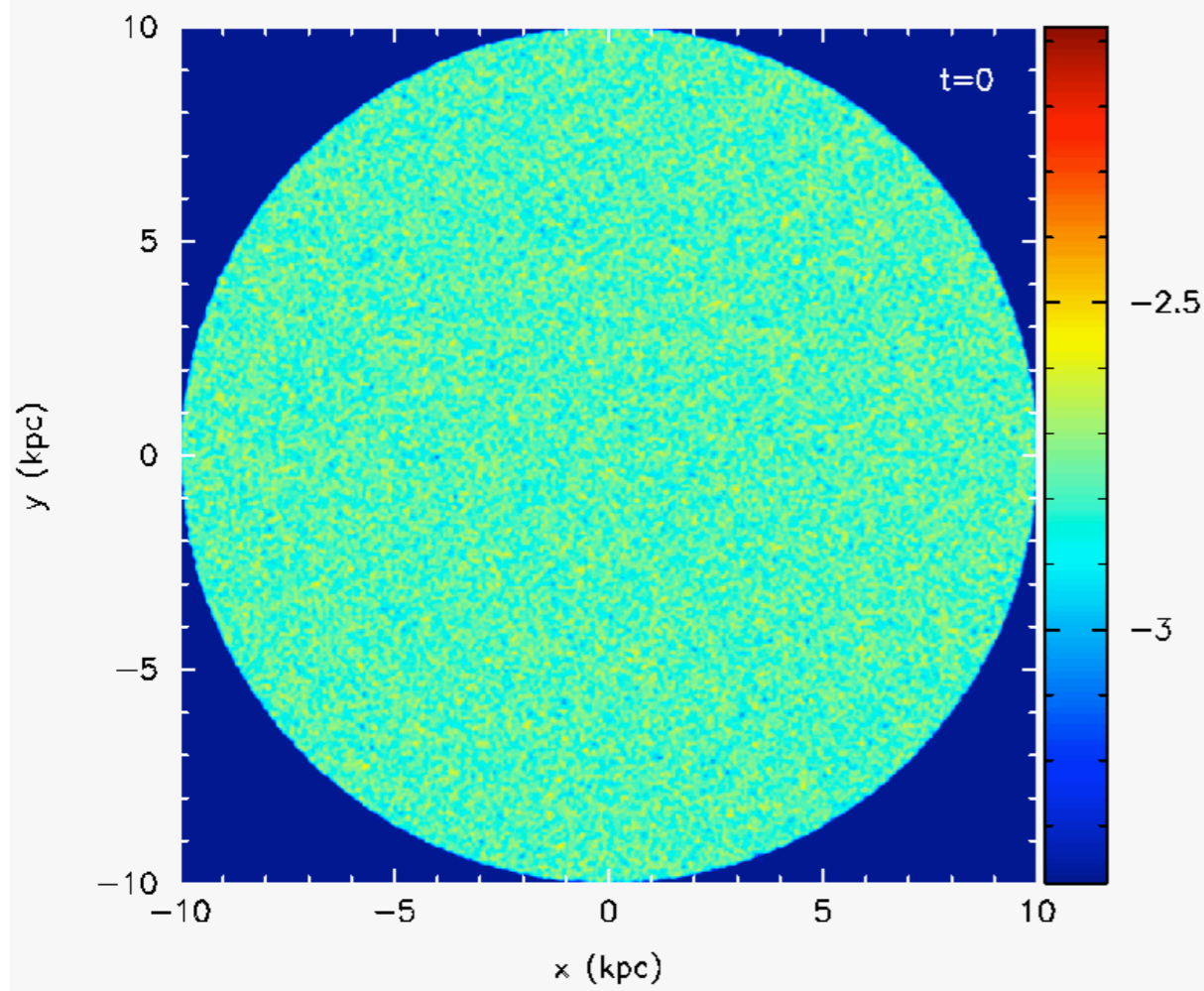
Filamentary interarm features (spurs)



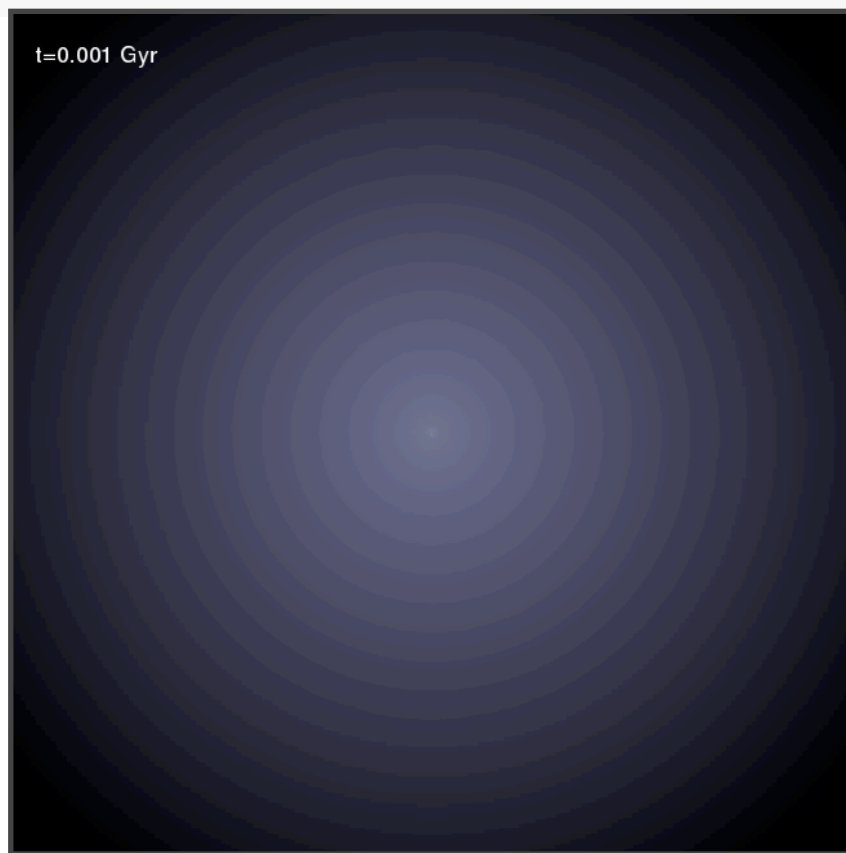
Agertz

Dobbs, Burkert & Pringle 11a,b, 12a,b

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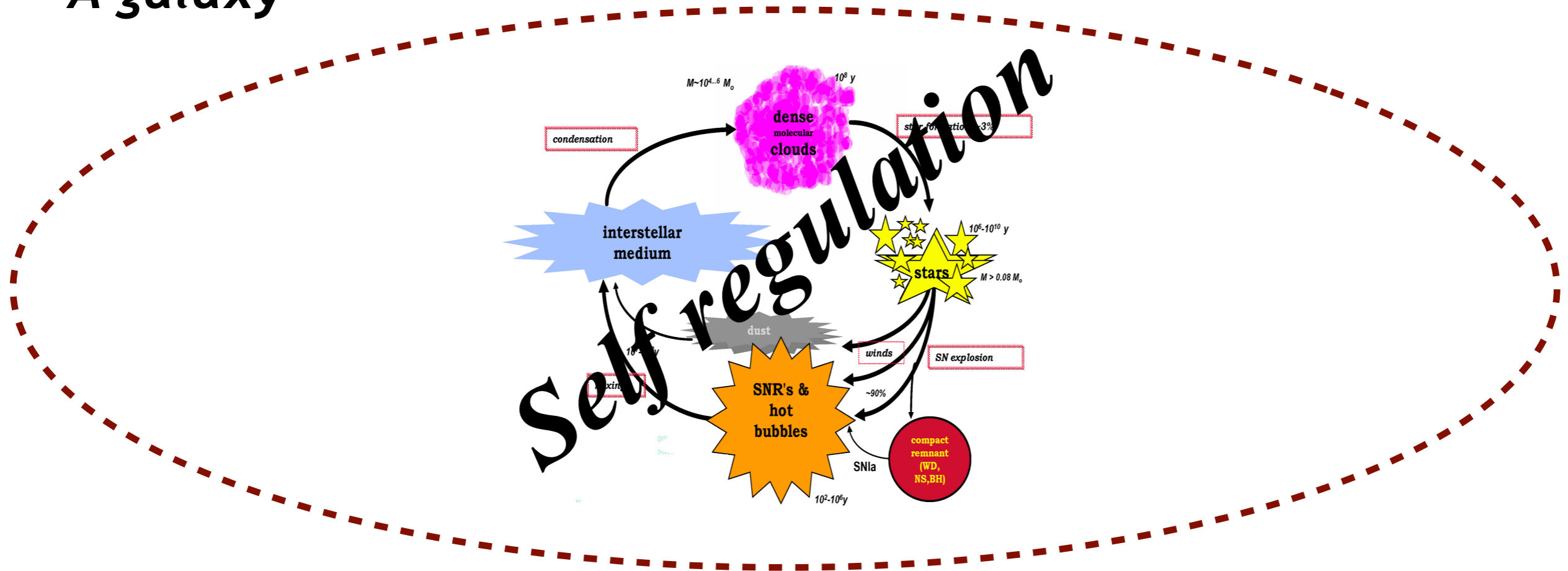


Agertz

The closed box model of self-regulated disk galaxy formation

(Wada+04; Dobbs+11a,b,12a,b; Bonnell+13; Brunner+13)

A galaxy



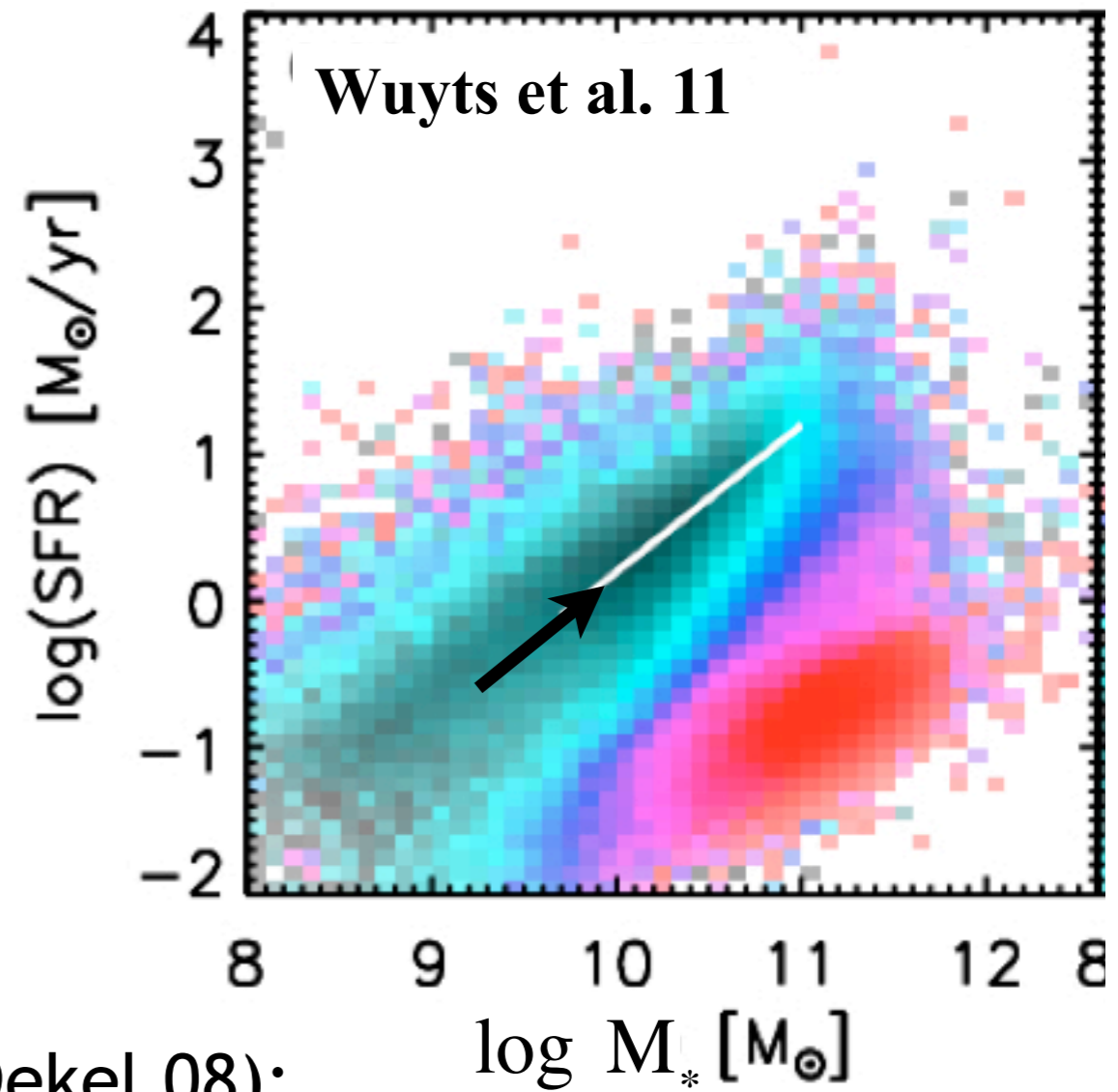
Evolution:

$$SFR(t) \downarrow \quad \text{and} \quad M_* \uparrow$$

The galaxy main sequence

Galaxy main sequence (Noeske et al. 07; Daddi et al. 07, Peng et al. 10, Bouche et al. 10, Wuyts et al. 11):

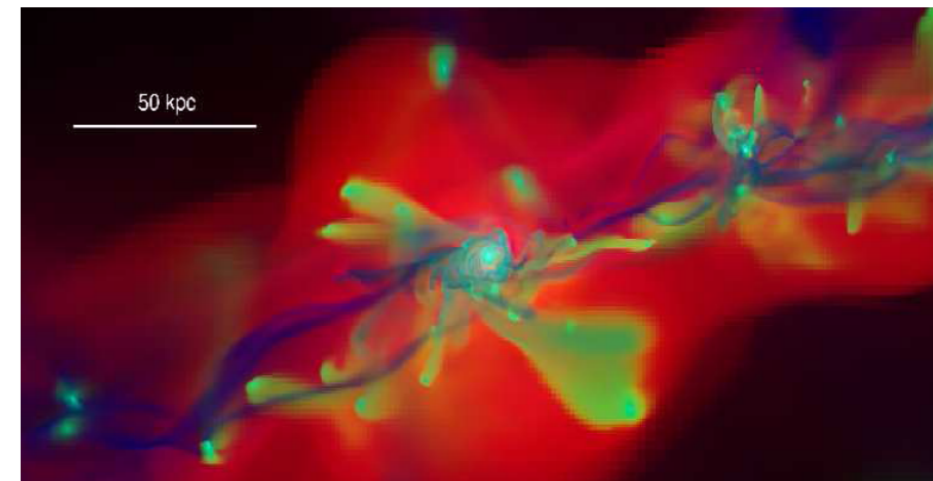
$$SFR \approx 6 \left(\frac{M_*}{10^{11} M_\odot} \right) (1+z)^{2.5} \frac{M_\odot}{yr}$$



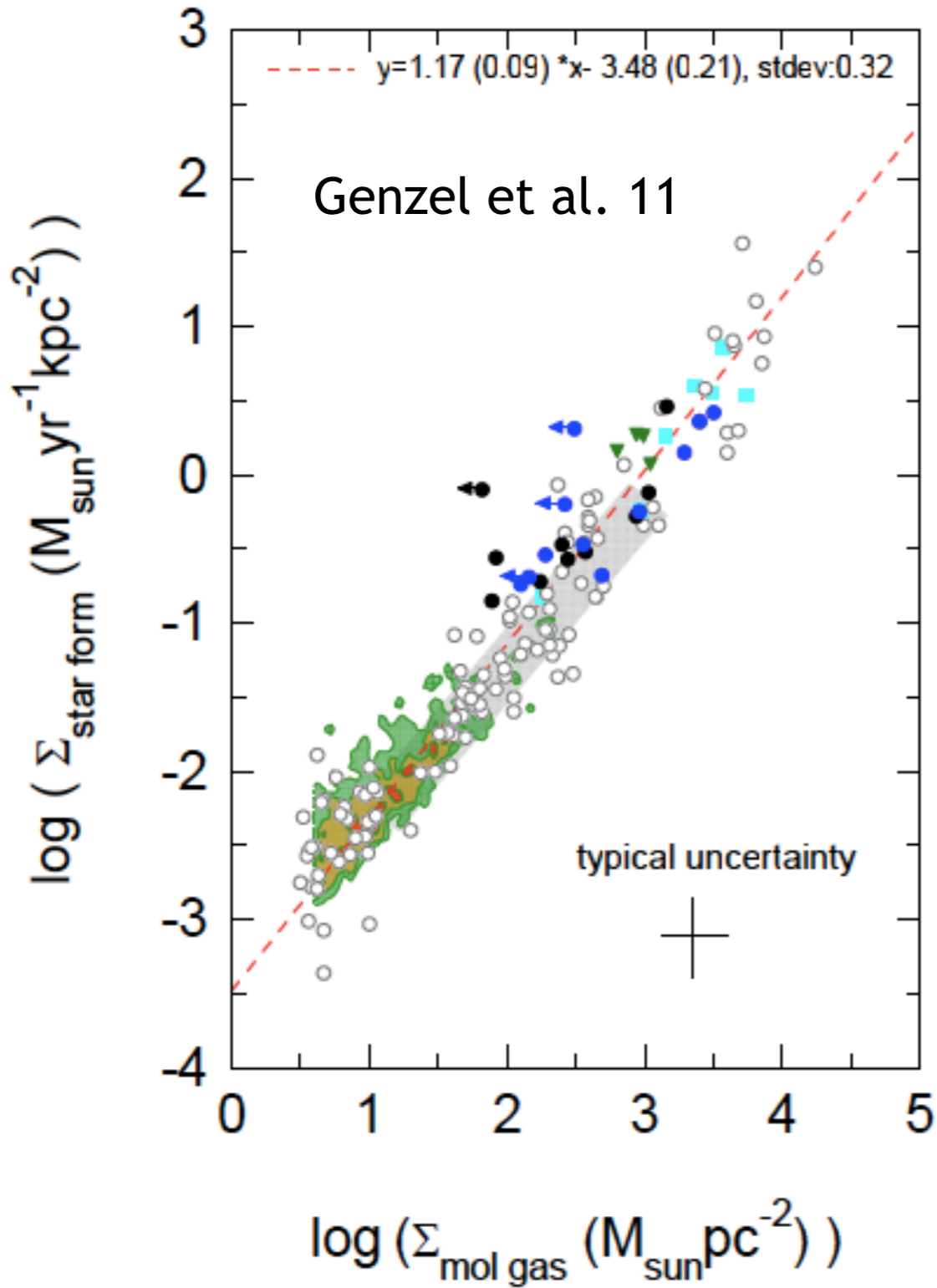
Cosmic baryonic accretion rate (Neistein & Dekel 08):

$$\left(\frac{dM_g}{dt} \right)_{acc} \approx 7 \cdot \epsilon_g \left(\frac{M_{DM}}{10^{12} M_\odot} \right)^{1.1} (1+z)^{2.2} \frac{M_\odot}{yr}$$

(Birnboim & Dekel 03; Dekel & Birnboim 06; Ceverino et al. 10, 12)



The universal gas depletion timescale



$$SFR = \frac{M_{H_2}}{\tau_{sf}} \text{ with } \tau_{sf} \approx 1 - 2 \cdot 10^9 \text{ yrs}$$

- Central limit theorem
- τ_{sf} is almost independent of redshift
- Gas depletion timescale **50 times** greater than local free-fall timescale.

$$\tau_{ff} \ll \tau_{sf} < \tau_{\text{Hubble}}$$



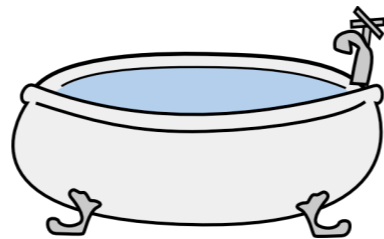
continuous replenishment

(Bouché et al. 07, McKee & Ostriker 08, Genzel et al. 10,11, Daddi et al. 10, Dave 11,12, Krumholz+ 12, Lilly et al. 13, Forbes et al. 13)

What determines SFR?

(Bouche et al. 10; Davé et al. 11a,b; Forbes et al. 13)

$$\frac{dM_g}{dt} = \left(\frac{dM_g}{dt} \right)_{acc} - \frac{M_g}{\tau_{sf}} (1 - R + \alpha_{wind})$$



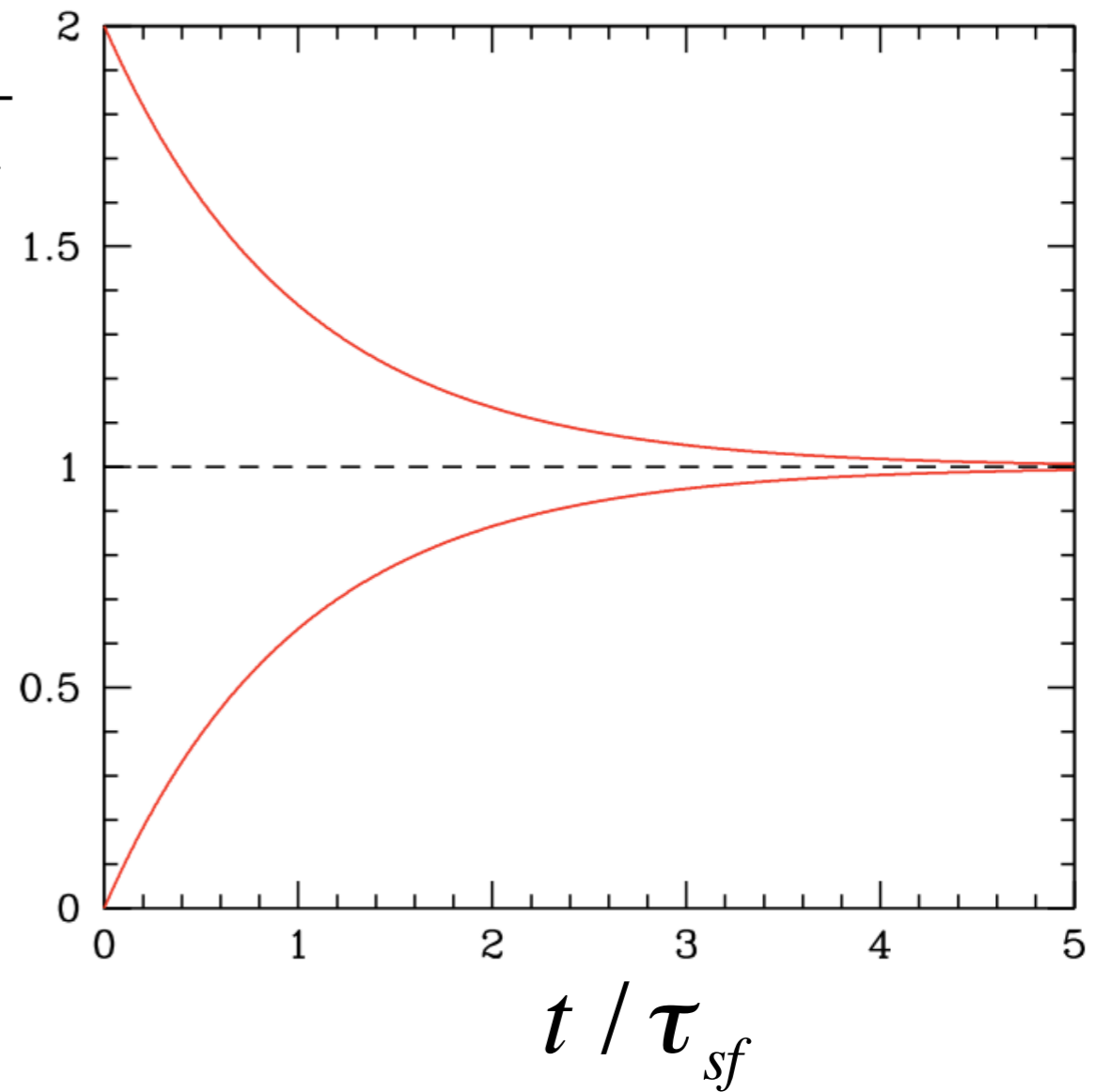
$$\dot{M}_{acc,eff}$$

$$SFR = \frac{M_g}{\tau_{sf}} = \frac{1}{1 - R + \alpha_{wind}} \left(\frac{dM_g}{dt} \right)_{acc} \left(1 - \exp\left(-\frac{t}{\tau_{sf}} \right) \right)$$



$$SFR = \dot{M}_{acc,eff}$$

$$\frac{SFR}{\dot{M}_{acc,eff}}$$



- τ_{sf} does not determine SFR

What's about metallicity?

$$Z_g = Z_{IGM} + y \frac{R}{\alpha_{wind} + R}$$

(Everett+ 8,10, Brook+ 11, Hopkins+ 12, Dalla Vecchia+ 12, Bolatto+ 13, Hirschmann+13, von Glasow+ 13, Hanz+ 13, Agertz+ 13)

What's about the (molecular) gas mass?

$$M_g = \dot{M}_{acc,eff} \cdot \tau_{sf}$$

Implications of the bathtub model

(Forbes+, astro-ph/1311.1509; Burkert+14)

$$M_g = SFR \cdot \tau_{sf} \sim M_* \cdot \tau_{sf} \longrightarrow$$

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$$M_g = SFR \cdot \tau_{sf} \sim M_* \cdot \tau_{sf} \quad \tau_{sf} \sim \frac{M_g}{M_*} \quad \longrightarrow$$

$$\tau_{sf} \approx 3 \text{Gyr} (1+z)^{-1} \left(\frac{M_{vir}}{10^{12} M_\odot} \right)^{-0.5}$$

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Implications of the bathtub model

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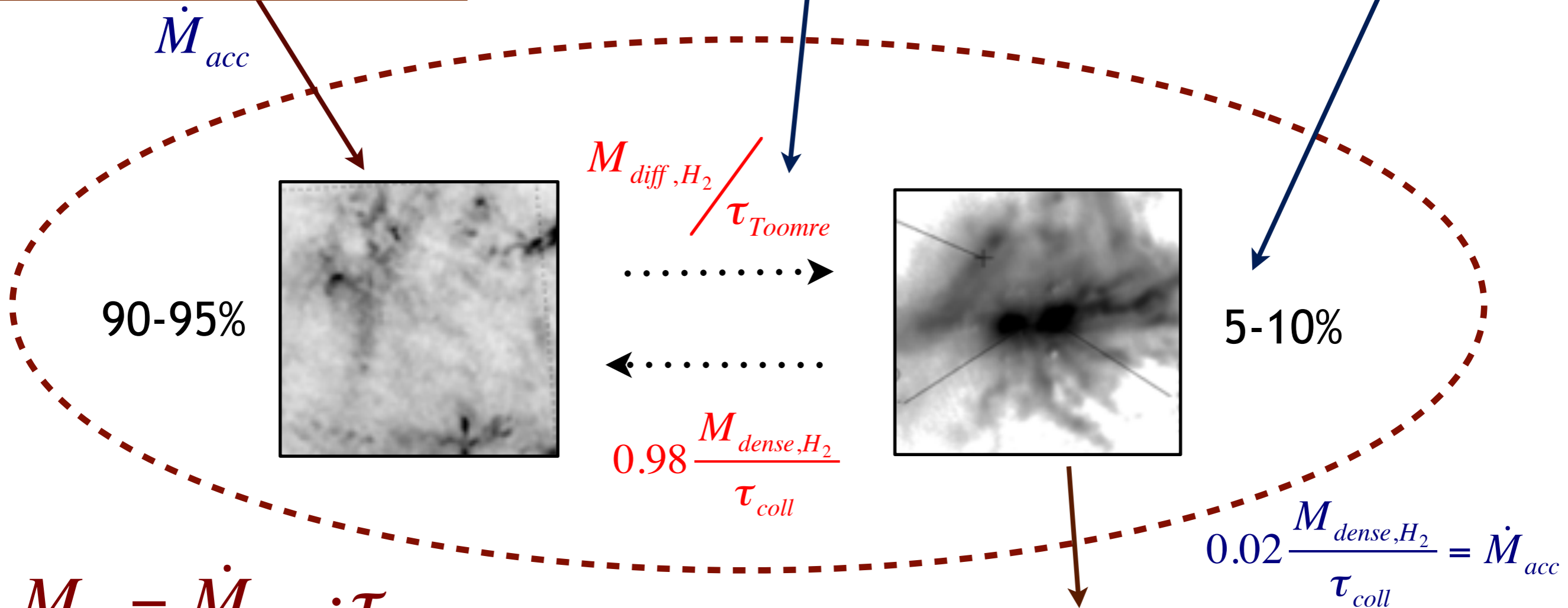
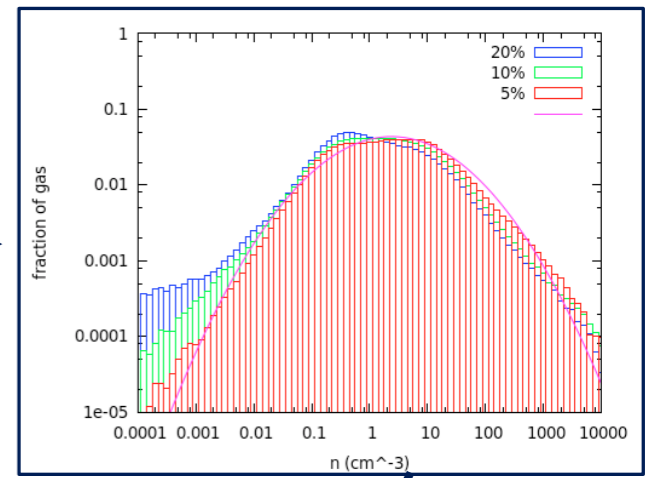
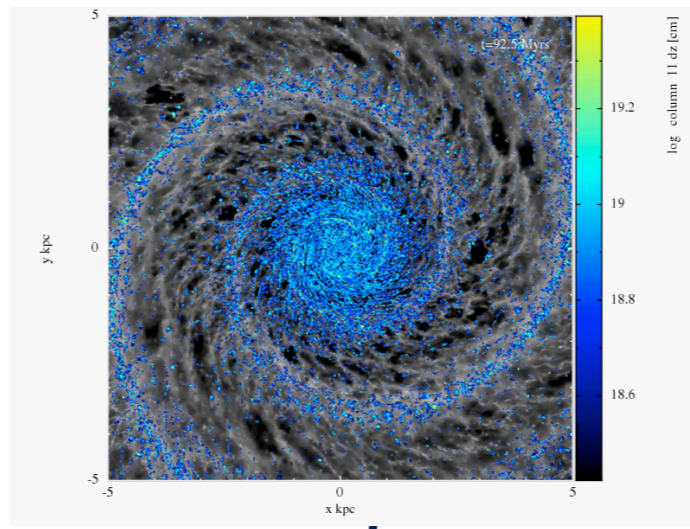
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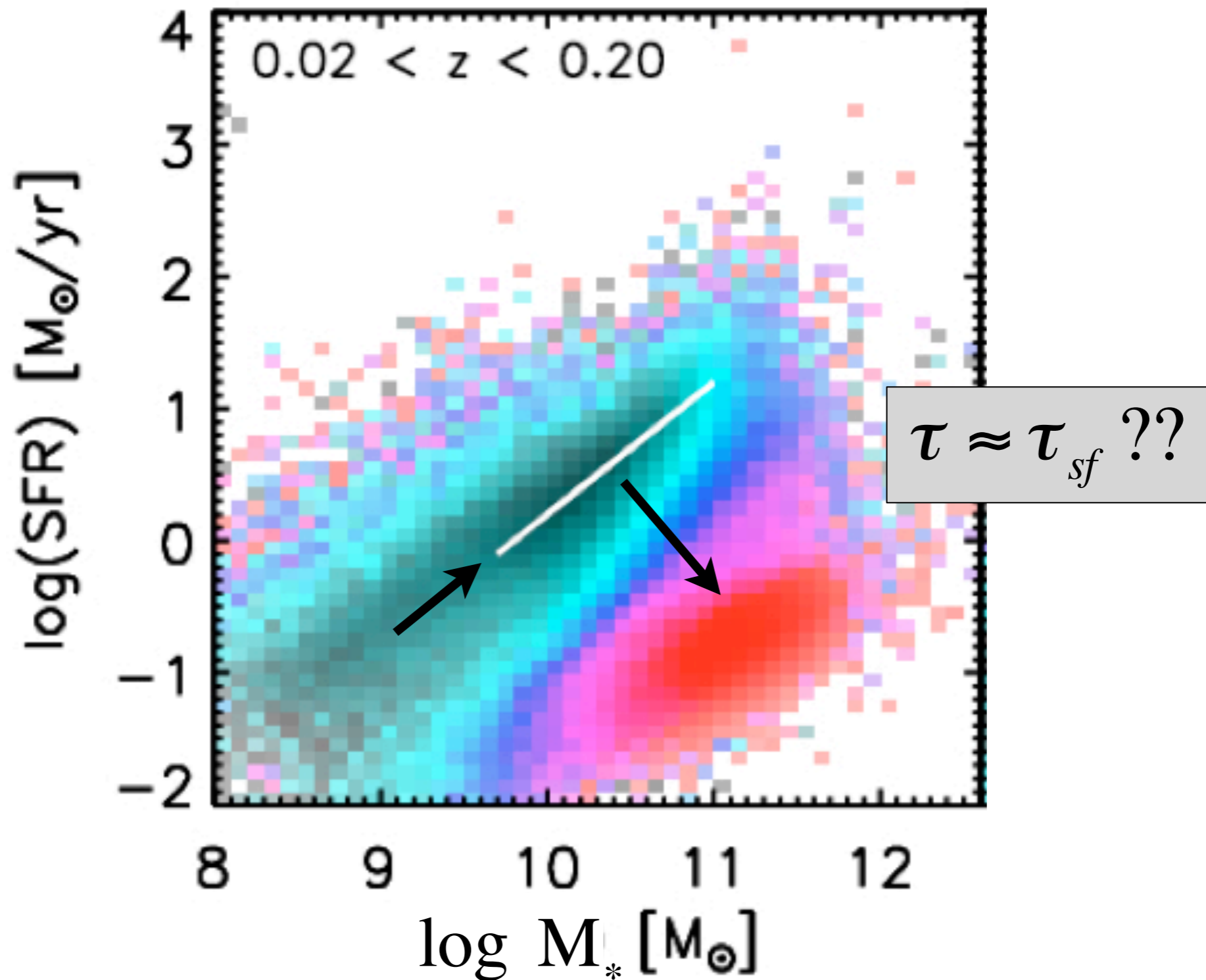
$$\alpha_{wind} = \left(\frac{M_{vir}}{10^{12} M_\odot} \right)^{-2/3}$$

This is consistent with recent models of **cosmic-ray driven** galactic winds (Wadepuhl&Springel 11; Salem&Bryan 13, Booth+13)



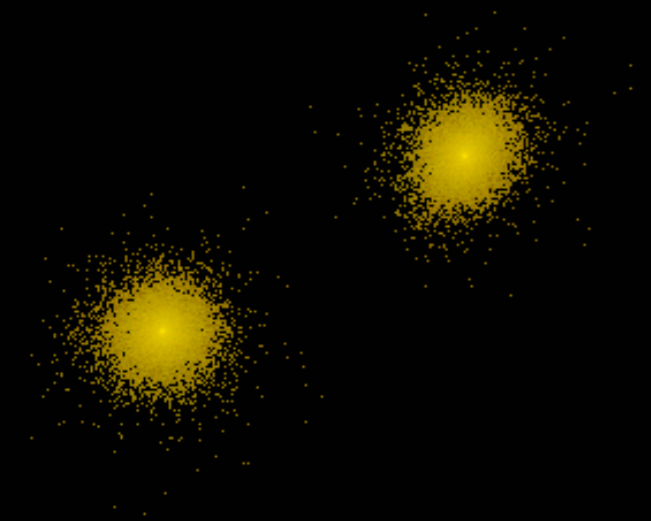
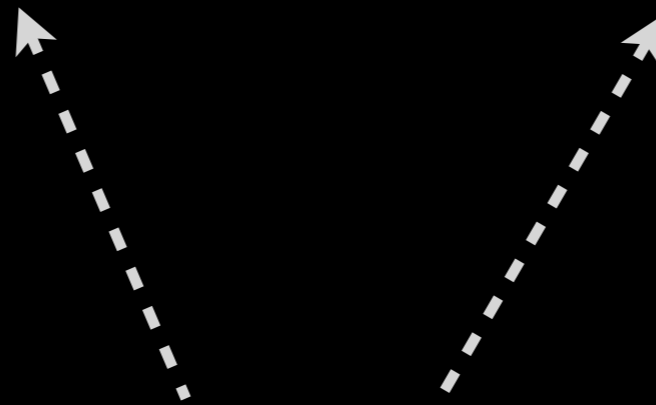
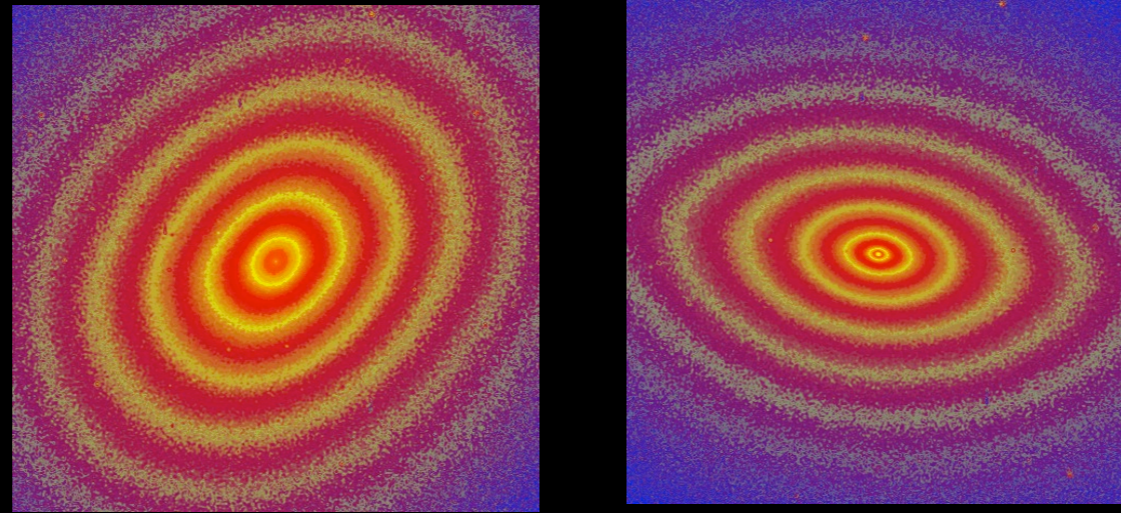
Self regulated evolution of disk galaxies

Evolution off the galaxy main sequence and the formation of red and dead galaxies



Major mergers clearly happen

(e.g. Hopkins+ 03-11; Naab+ 03-11, Johansson+09-11; Remus+12, see poster by Schauer+13)



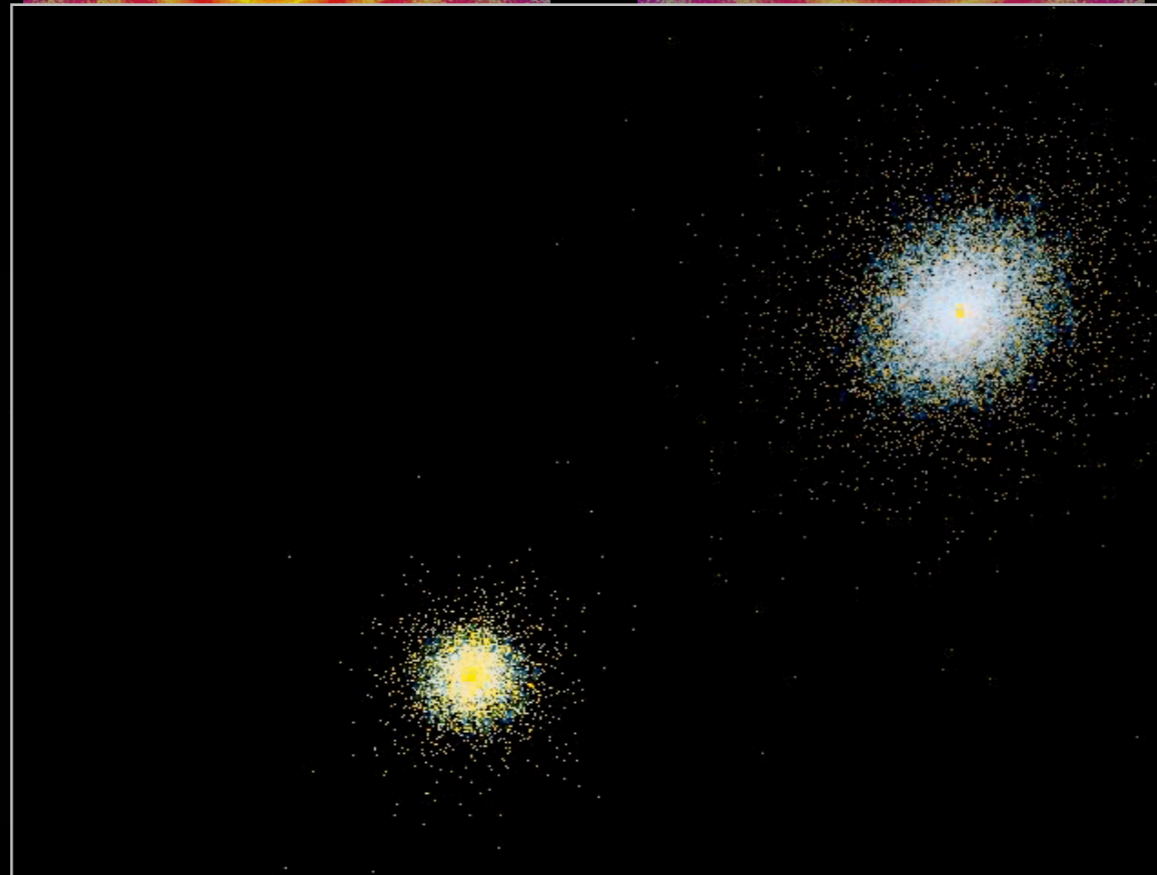
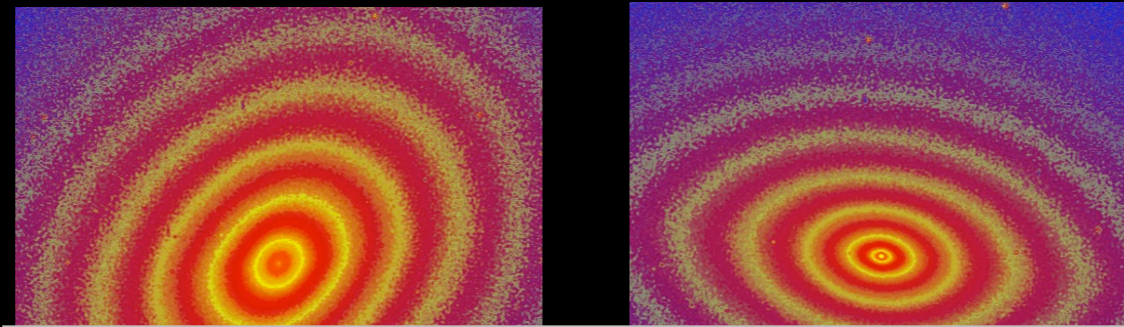
(Naab)



(Springel)

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$T = 160$ Myr



(Naab)

(Springel)



Disks can sometimes be quite robust

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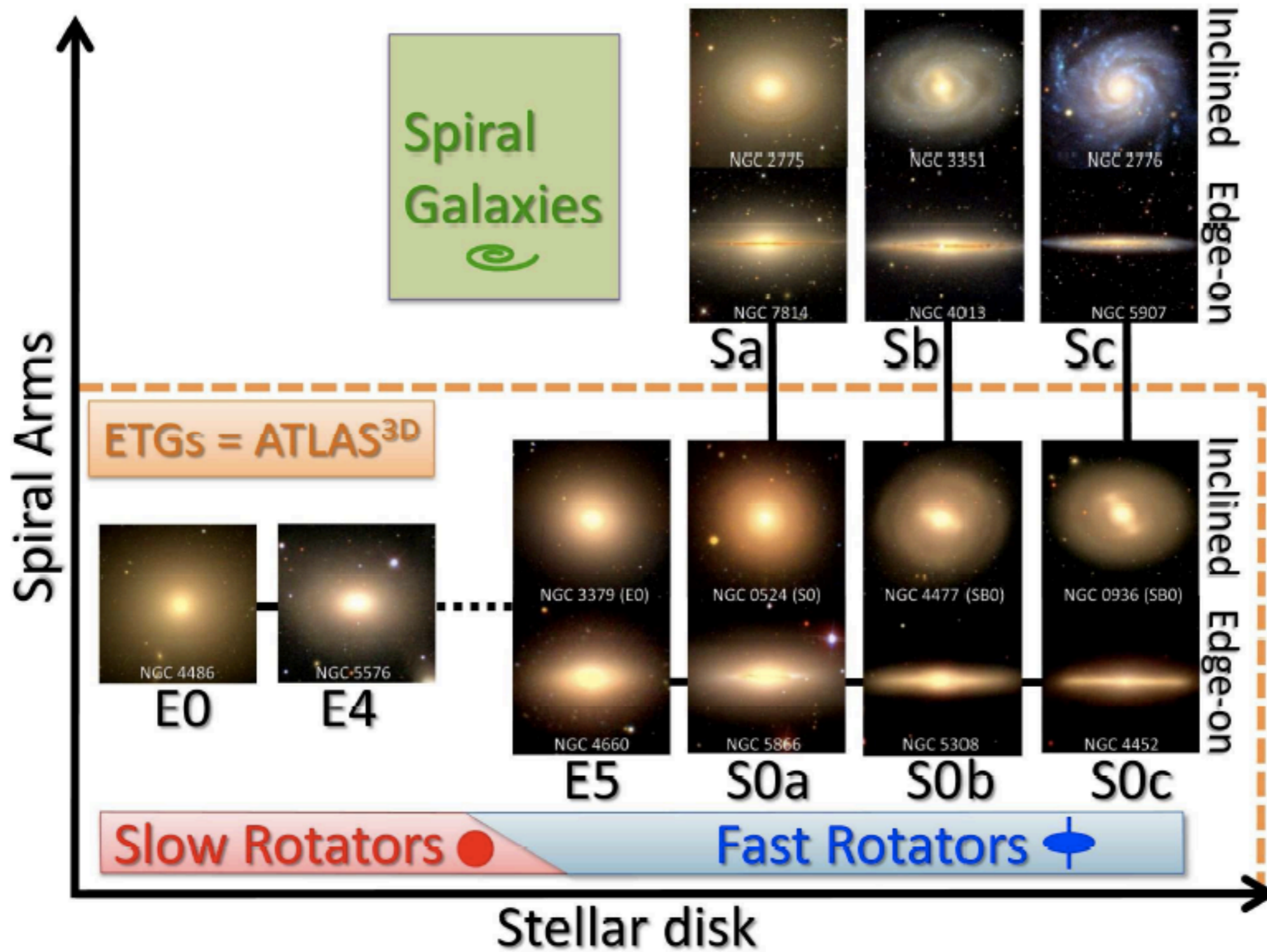
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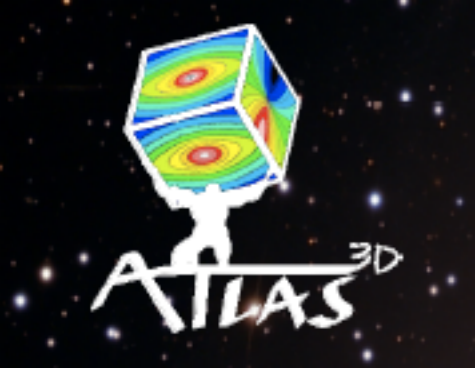
Copyright R. Teysaler (2008)

Teyssier 08

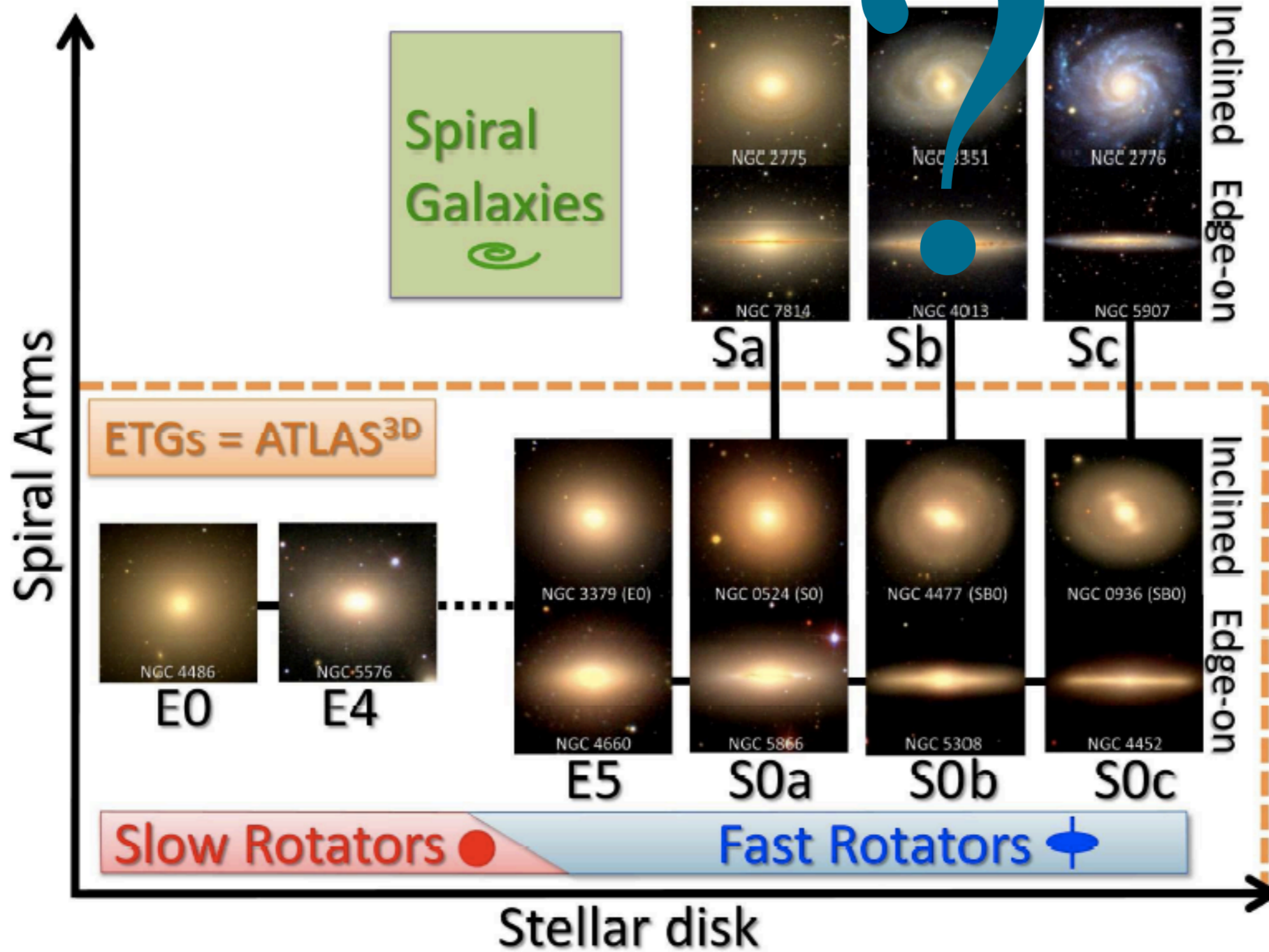


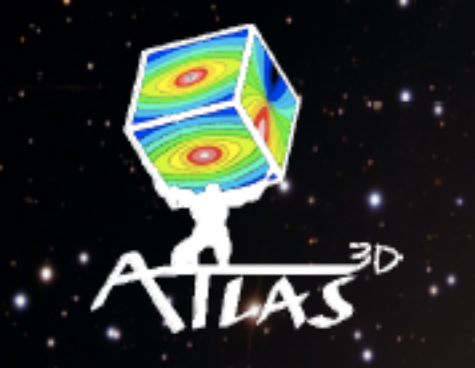
Cappellari+11,13



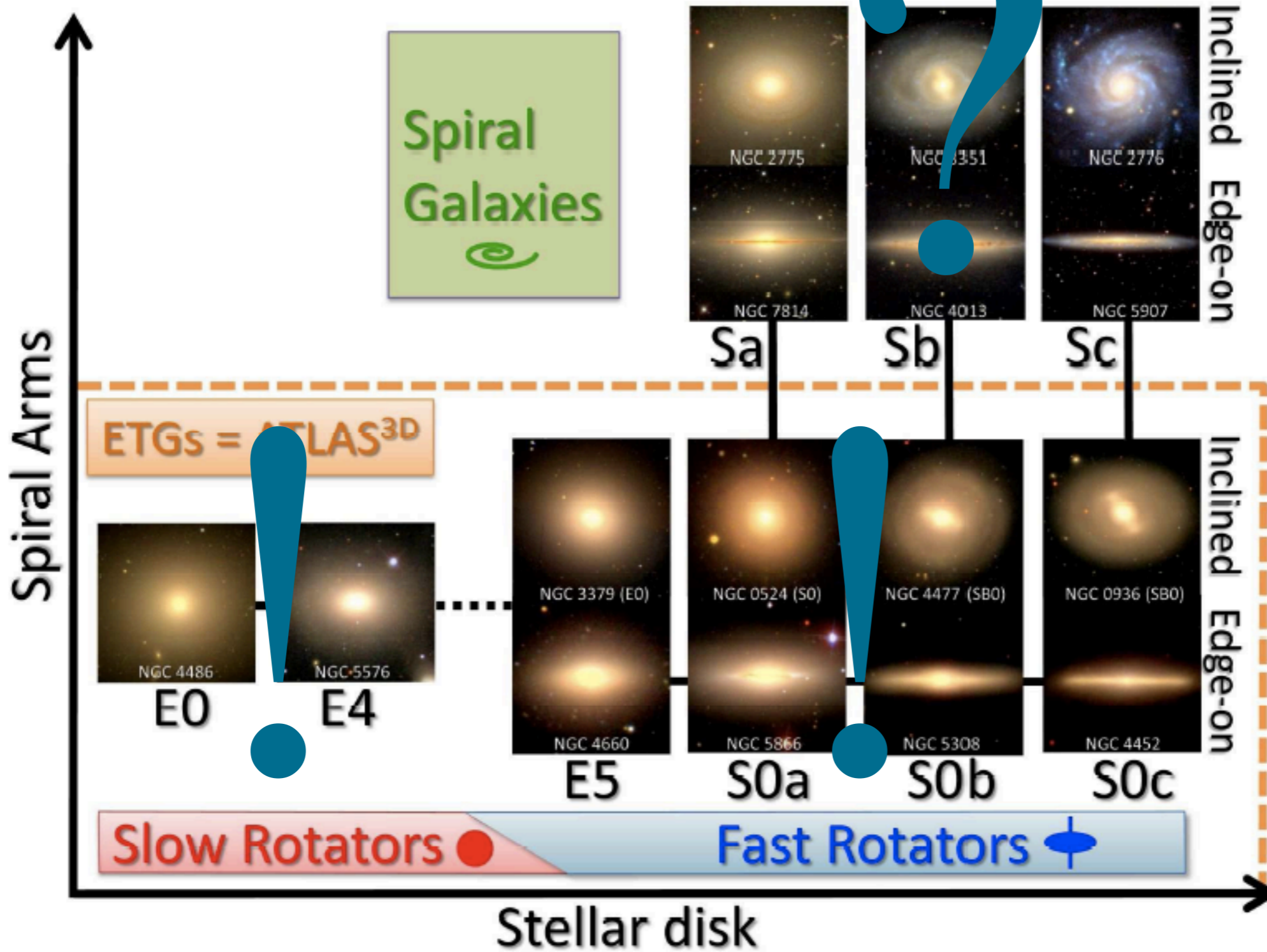


Cappellari+11,13





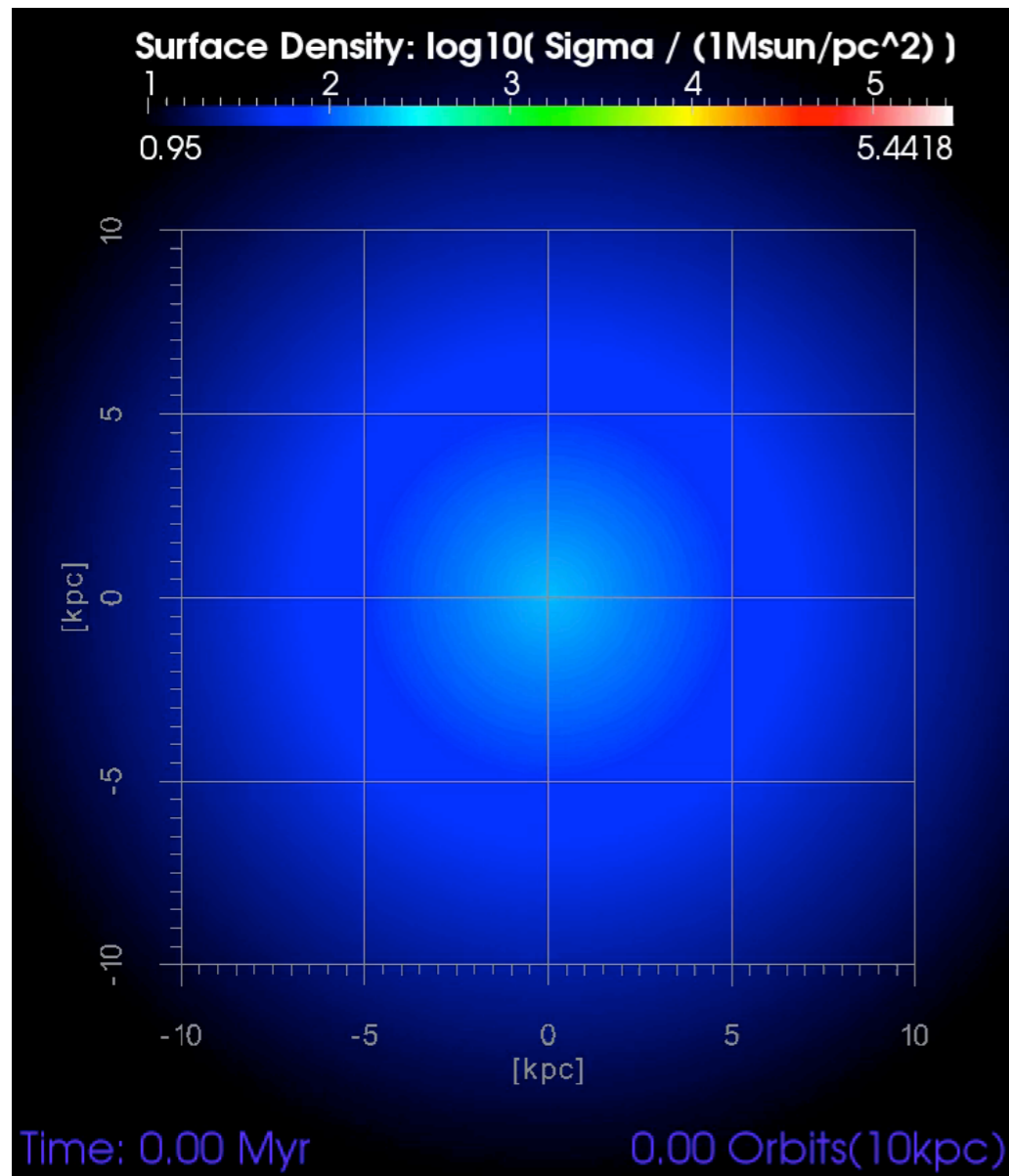
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Connecting high-z and low-z

- High-z compact ellipticals are flattened and **disky** (Bezanson+09, van der Wel+11, 13; Chang+ 13)
- They might have formed from extended high-z massive gas disks, going through **violent disk instability** (Dekel & Burkert, astro-ph/1310.1074)

Behrendt+13



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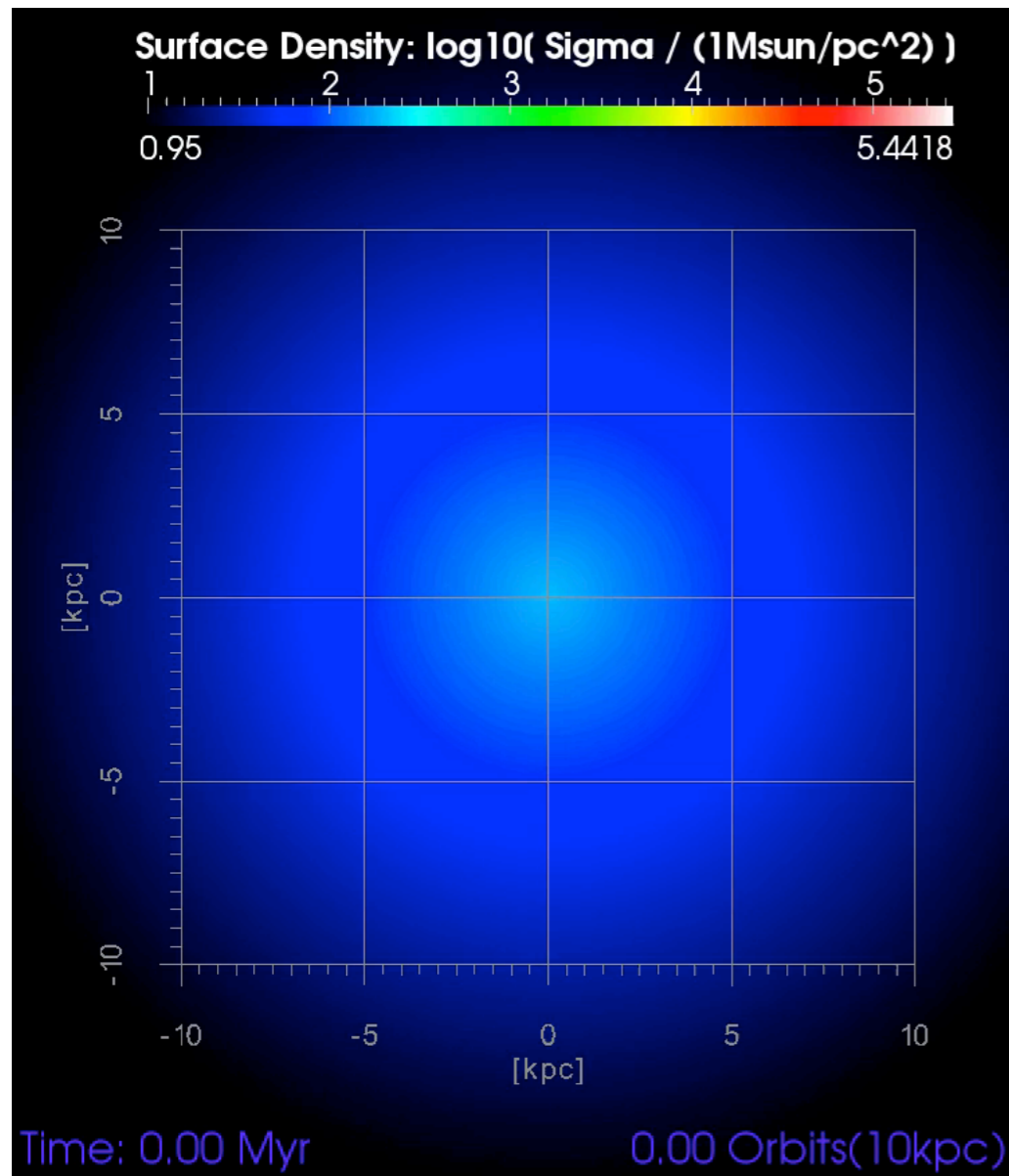
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High $\dot{M}_{acc} \sim (1+z)^{2.2} \rightarrow$ High $\delta_g \equiv \frac{M_g}{M_*}$

$$Q = 1 \rightarrow \sigma = \delta_g \cdot v_{rot}$$

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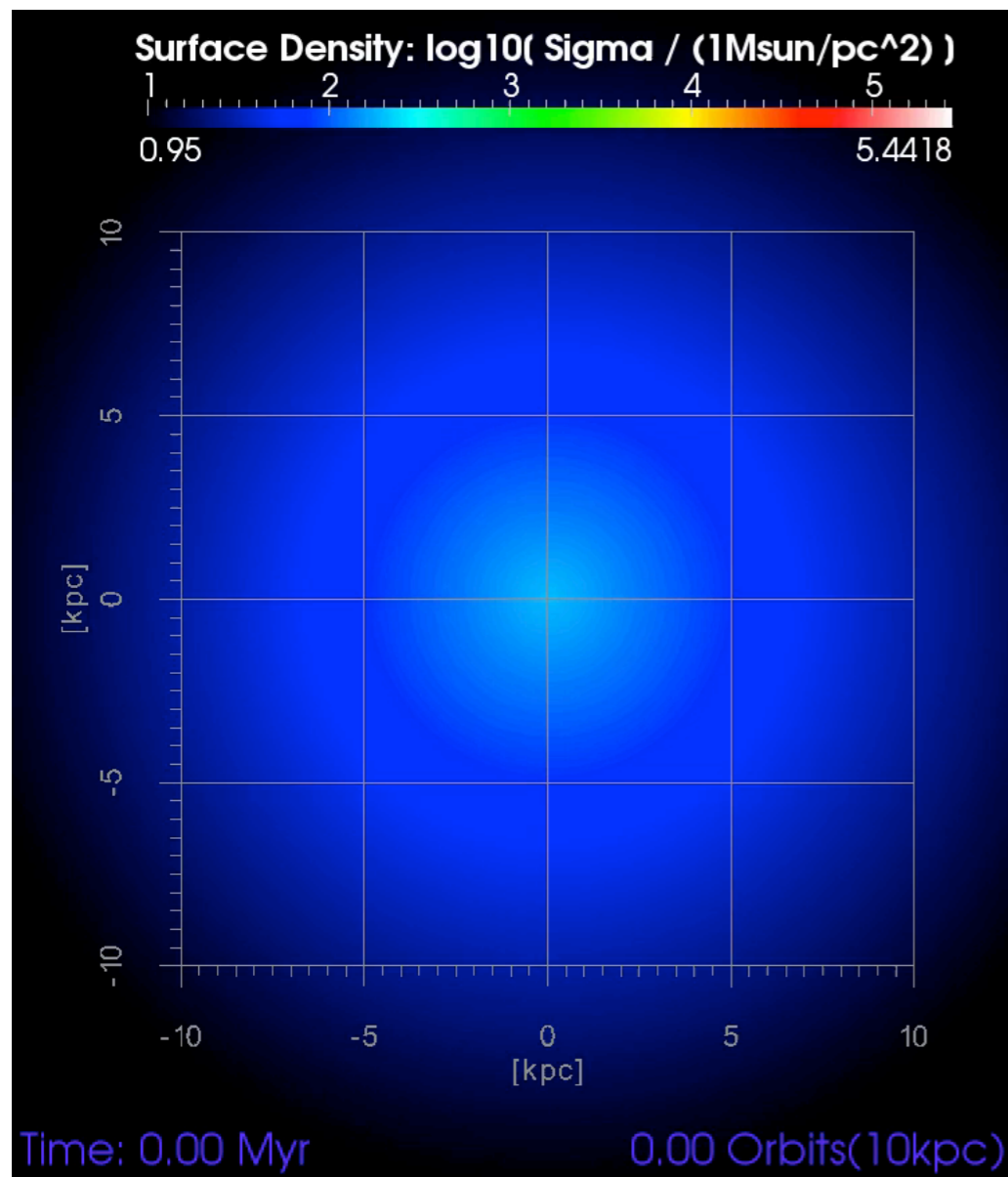
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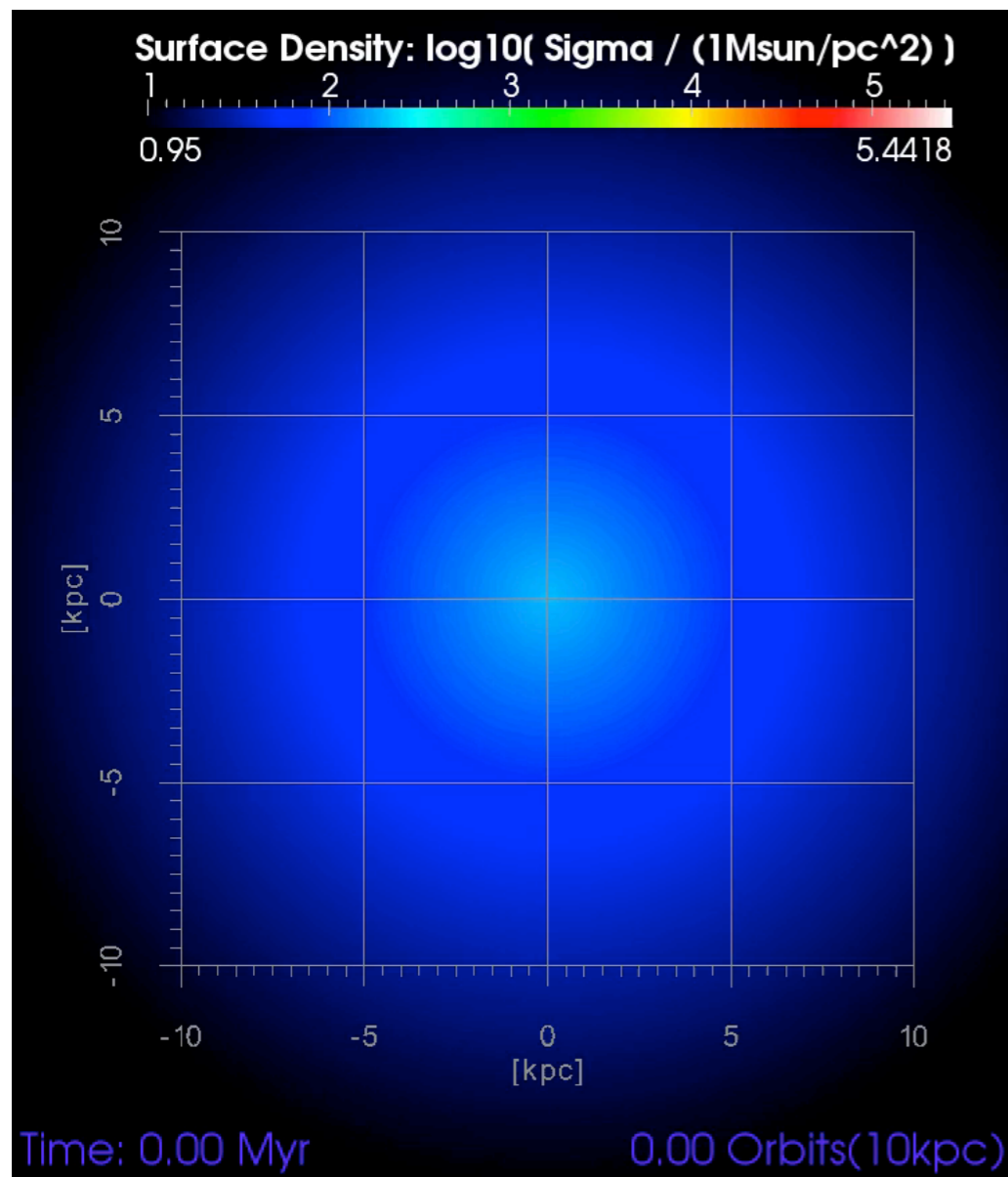
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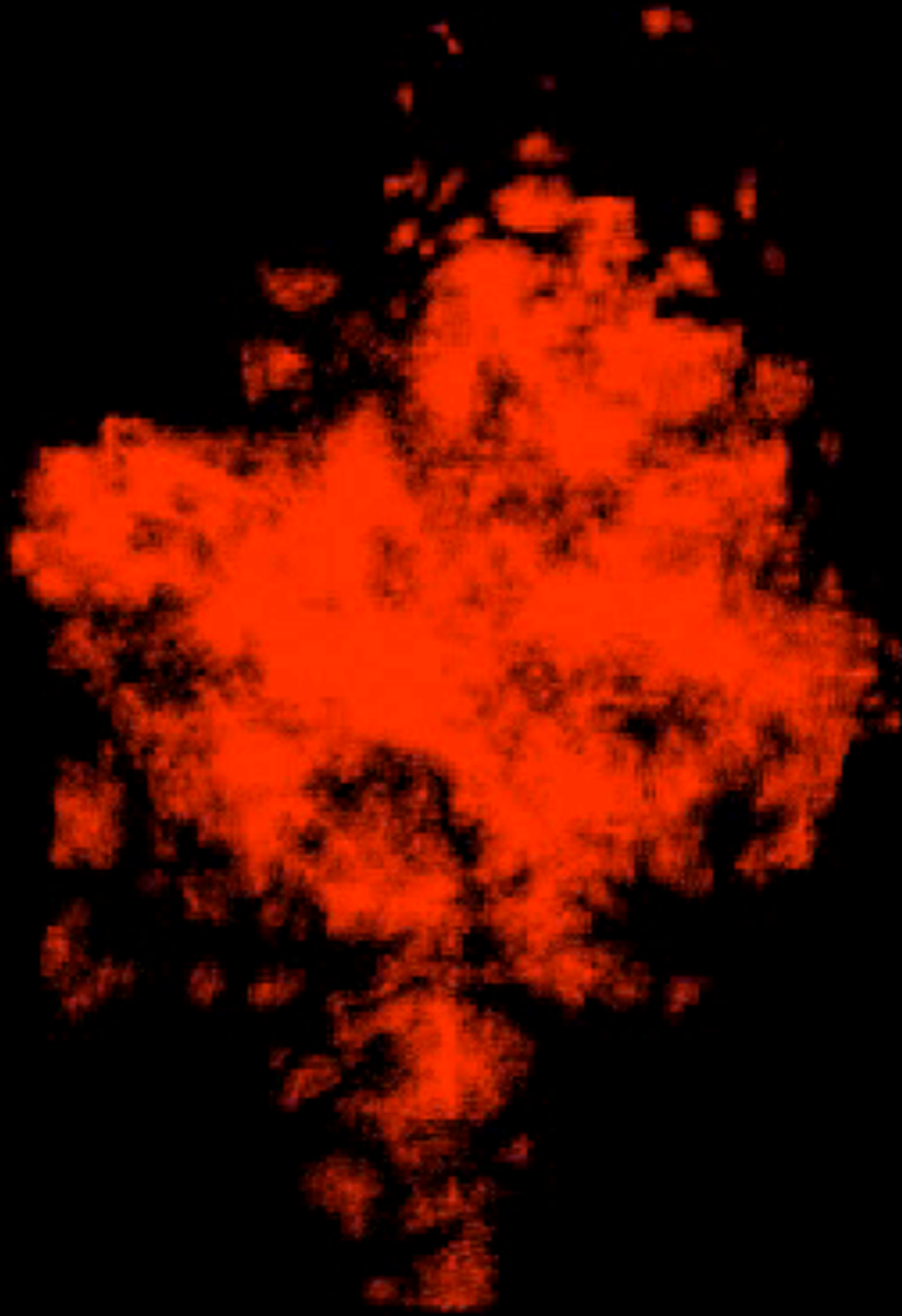
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VDI: (Dekel&Burkert 13)

$$\tau_{sf} < \tau_{inflow} \rightarrow \lambda_{disk} \leq 0.05$$

Behrendt+13

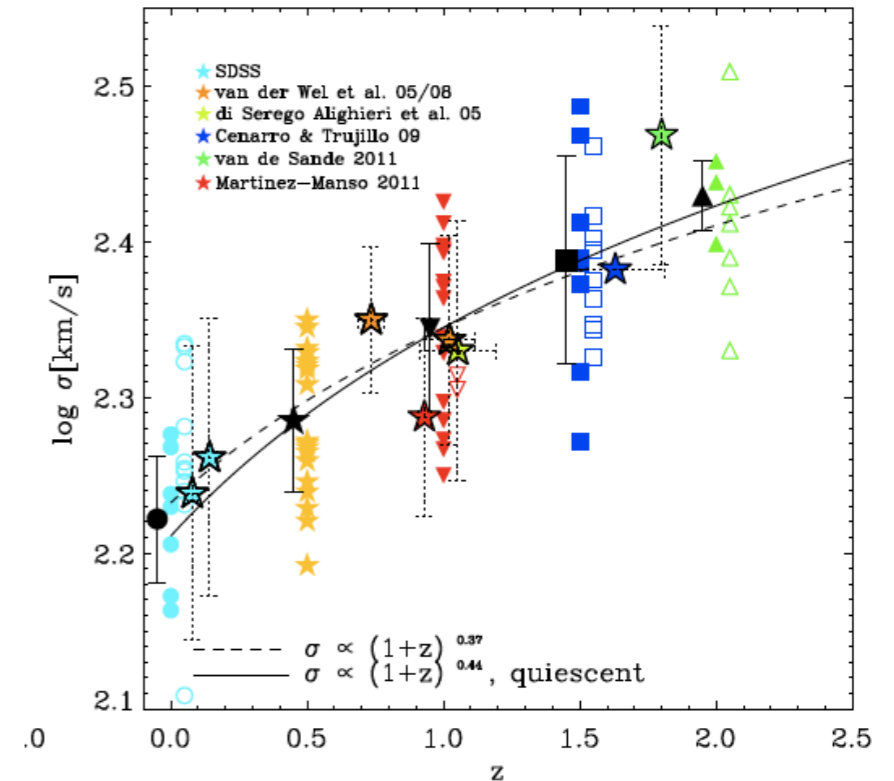
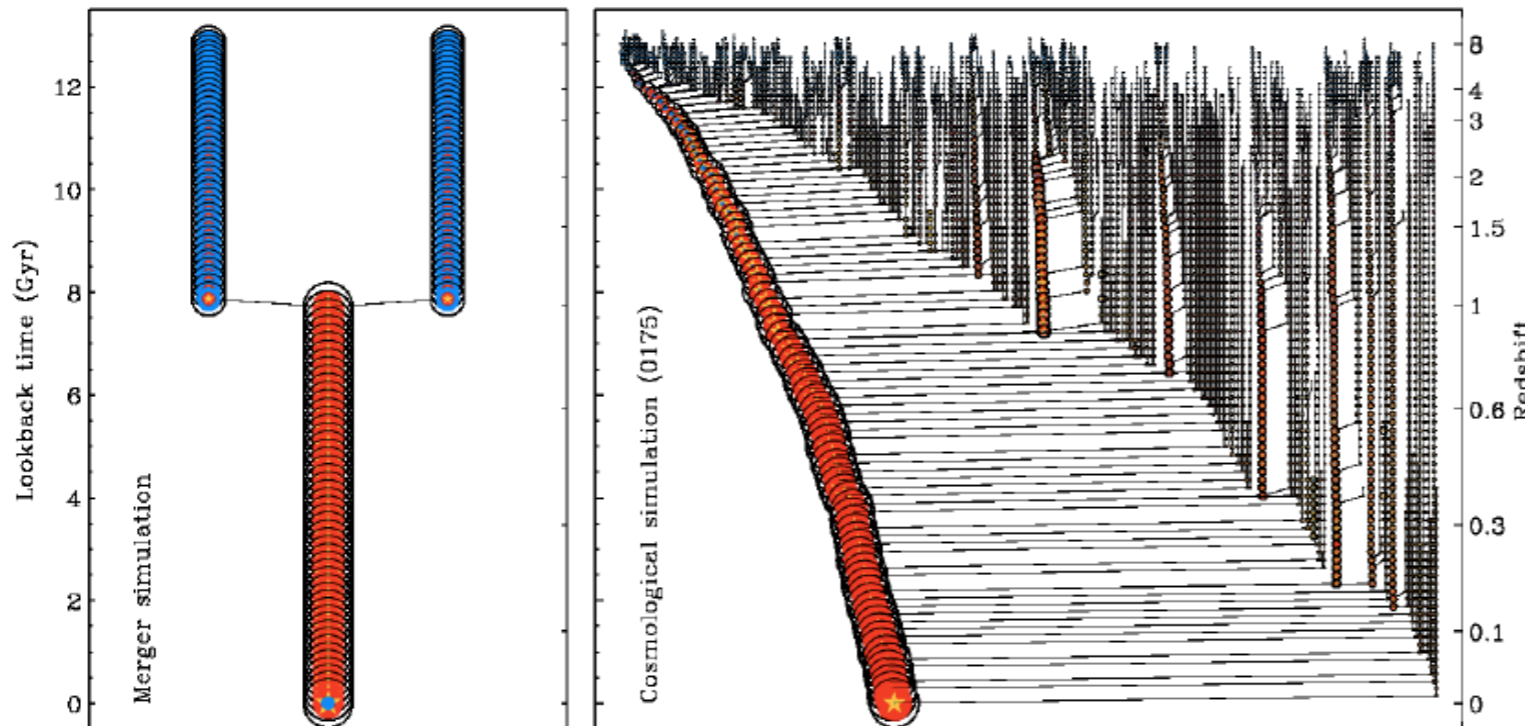
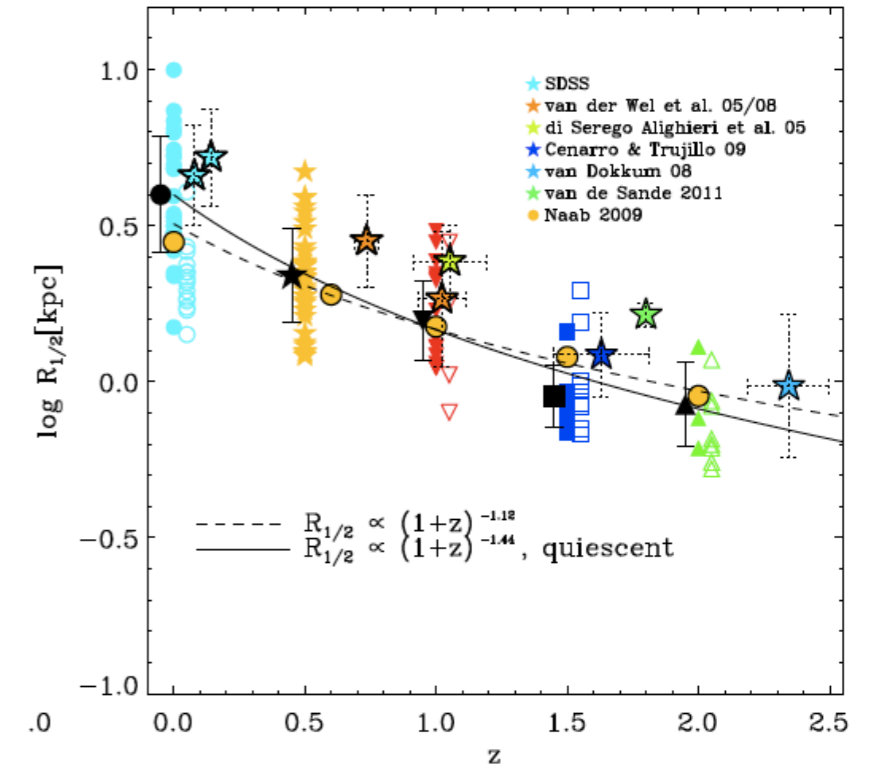
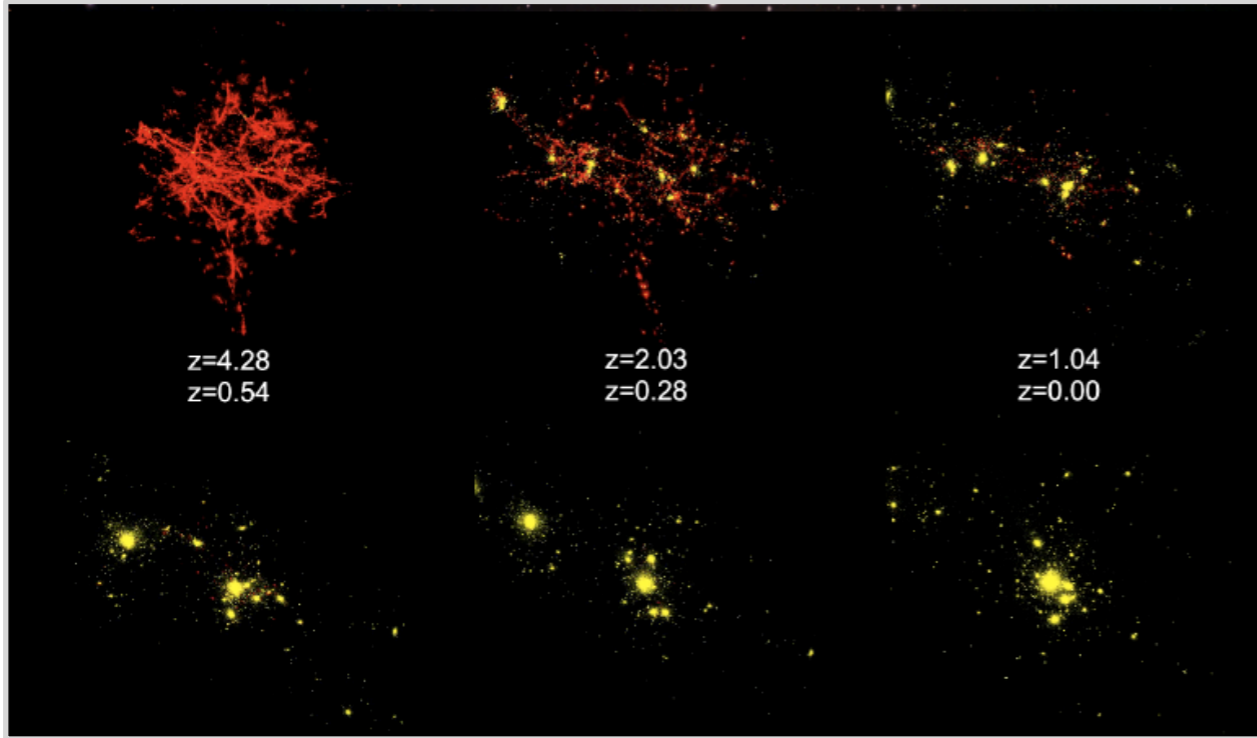




Oser 10,12

The evolution of massive ellipticals

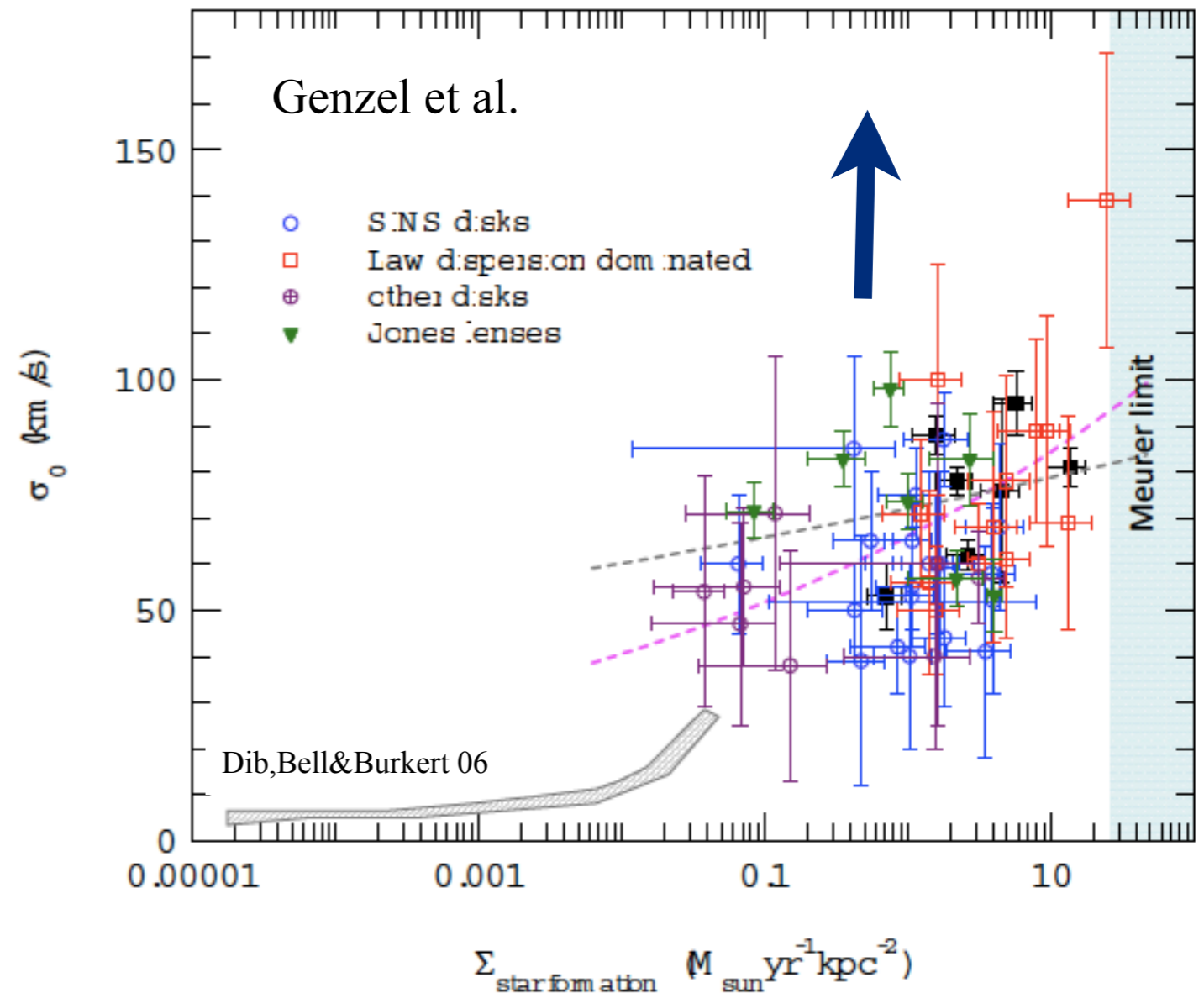
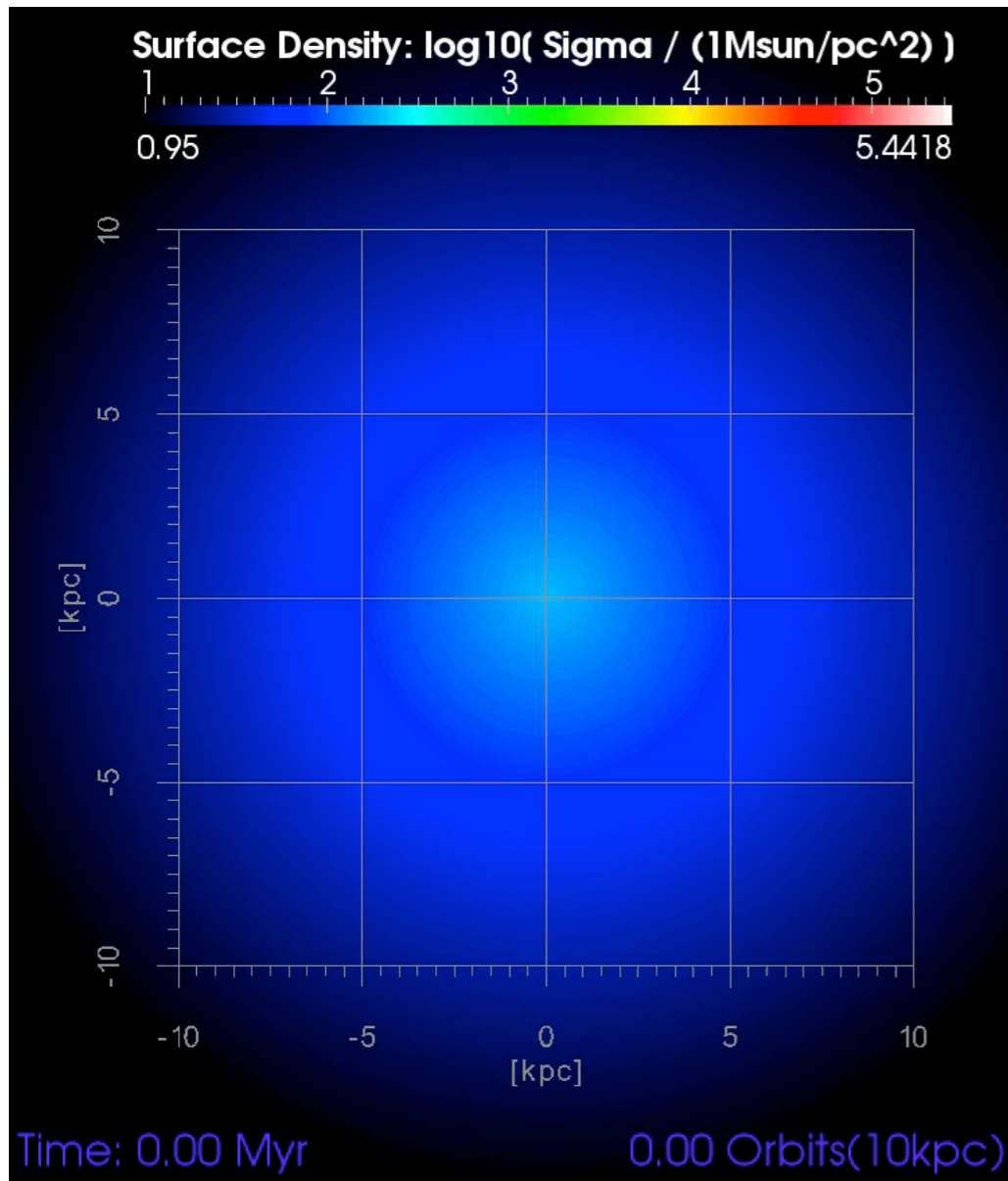
(de Lucia+06; Oser+10,12; Johansson+09,12; Hirschmann+12; Hiltz+12,13; Naab+07, 09,13; see however e.g. Gallego+12; Posti+13)



Connecting high-z and low-z

- High-z disks that did not experience VDI form fast rotating Es.

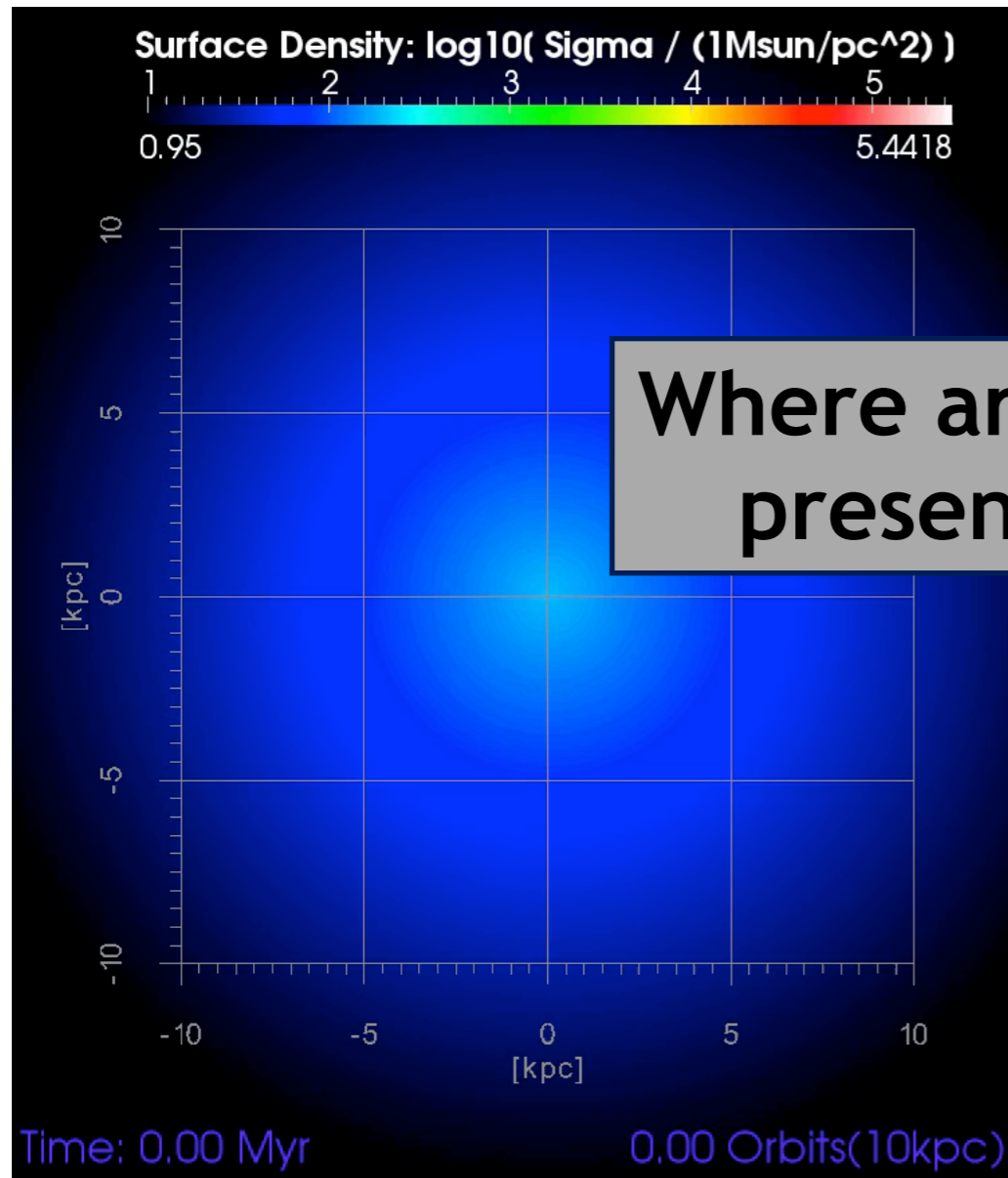
Behrendt+13



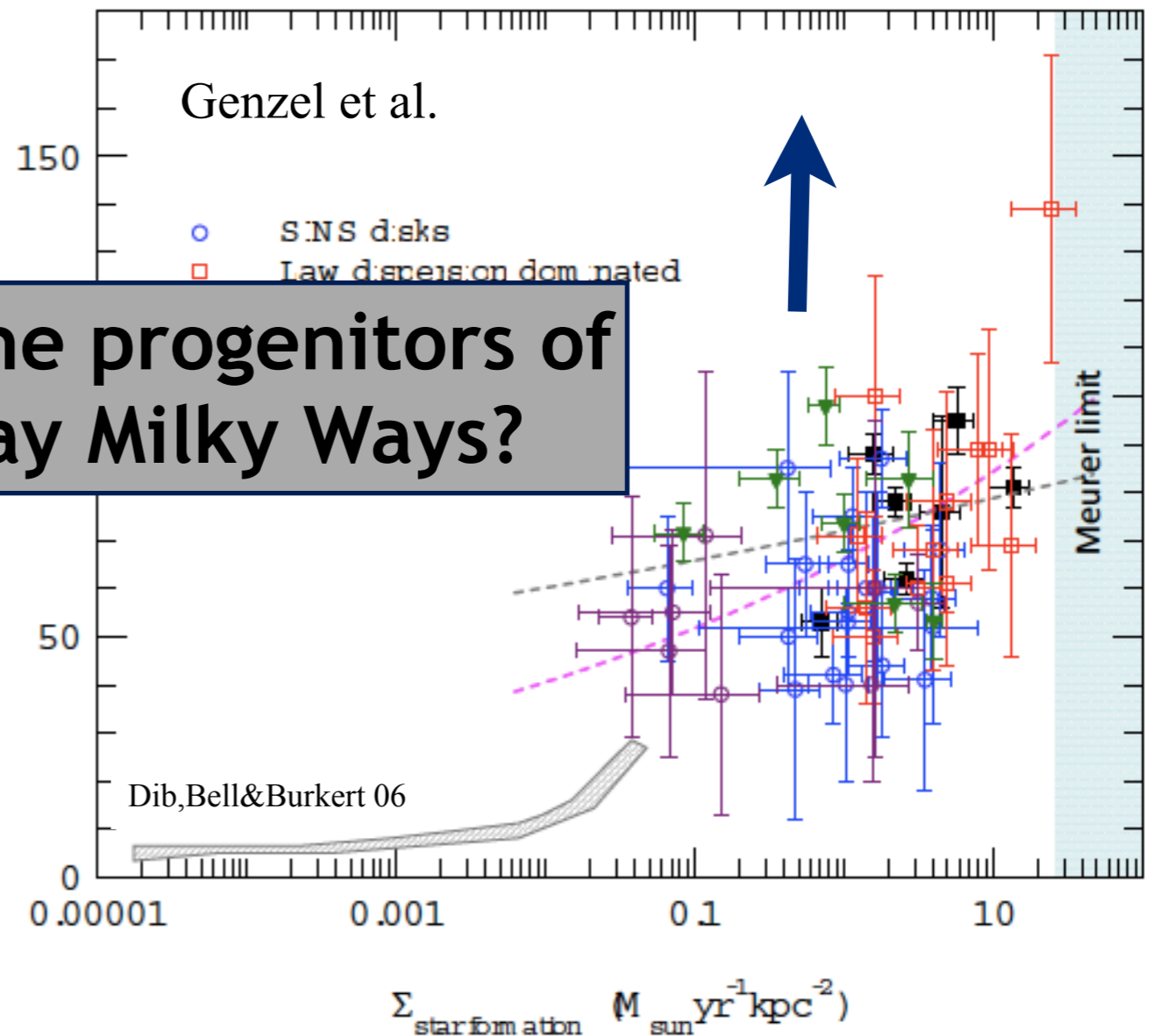
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Behrendt+13



Where are the progenitors of present-day Milky Ways?



Galaxy formation is a boundary condition problem

- Self-regulated disk evolution versus VDIs
- Depletion versus accretion timescale
- In situ versus ex-situ star formation
- Major versus minor mergers

