

Exercise 13: Remarks on Witten index and SUSY QM ①

• Bosonic degree of freedom:

$$a := \frac{1}{\sqrt{2}}(Q + iP) \quad \text{acting on } \mathcal{H}_B := L^2(\mathbb{R})$$

$$[a, a^\dagger] = 1$$

$$N_B := a^\dagger a \quad \text{spec } N_B = \{0, 1, 2, 3, \dots\}$$

• Fermionic degree of freedom:

$$f := \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \quad \text{acting on } \mathcal{H}_F := \mathbb{C}^2$$

$$\{f, f^\dagger\} = 1$$

$$N_F := f^\dagger f \quad \text{spec } N_F = \{0, 1\}$$

• SUSY oscillator: on $\mathcal{H} := \mathcal{H}_B \otimes \mathcal{H}_F$

$$H_{\text{SUSY}} := \hbar\omega (a^\dagger a \otimes 1 + 1 \otimes f^\dagger f)$$

$$= \hbar\omega \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a \end{pmatrix} + \hbar\omega \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$= \hbar\omega \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a \end{pmatrix}$$

$$=: \begin{pmatrix} H_+ & 0 \\ 0 & H_- \end{pmatrix} \quad \begin{array}{l} \mathcal{H}^+ \\ \mathcal{H}^- \end{array} \quad \begin{array}{l} \text{spec } H_+ = \{\hbar\omega n \mid n \in \mathbb{N}\} \\ \text{spec } H_- = \{\hbar\omega n \mid n \in \mathbb{N}_0\} \end{array}$$

$$\mathcal{H} = \mathcal{H}^+ \oplus \mathcal{H}^-$$

• Witten - Index:

$$\begin{aligned}
 \Delta &:= \text{Tr} \left((-1)^{N_F} e^{-\beta H_{\text{Susy}}} \right) \\
 &= \text{Tr}_{\mathcal{H}^+} \left((-1)^1 e^{-\beta H_+} \right) + \text{Tr}_{\mathcal{H}^-} \left((-1)^0 e^{-\beta H_-} \right) \\
 &= \text{Tr}_{\mathcal{H}^-} e^{-\beta H_-} - \text{Tr}_{\mathcal{H}^+} e^{-\beta H_+} = 1 \\
 &= n_- - n_+ \quad n_{\pm} := \# \text{ zero energy eigenstates of } H_{\pm}
 \end{aligned}$$

• Fredholm-index: (formally) A linear operator $\mathcal{X} \rightarrow \mathcal{X}$

$$\begin{aligned}
 \text{ind } A &:= \dim \text{Ker } A - \dim \text{Ker } A^\dagger \\
 &= \dim \text{Ker } A^\dagger A - \dim \text{Ker } A A^\dagger
 \end{aligned}$$

let $A = a$

$$\begin{aligned}
 \leadsto \text{ind } a &= \dim \text{Ker } H_- - \dim \text{Ker } H_+ \\
 &= n_- - n_+ = \Delta
 \end{aligned}$$

Exercise 14: Pseudoclassical mechanics

Classical mechanics with Grassmann-valued degrees of freedom

Grassmann algebra:

Generators: $\psi_0, \bar{\psi}_0$ such that $\psi_0^2 = 0 = \bar{\psi}_0^2$

and $\{\bar{\psi}_0, \psi_0\} = 0$
anticommute

General element:

$$B := a_1 + \underbrace{a_2 \psi_0 + a_3 \bar{\psi}_0 + a_4 \bar{\psi}_0 \psi_0}_{\text{soul}}, \quad a_i \in \mathbb{C}$$

\downarrow
 body

$$\bar{B} = a_1^* + a_2^* \bar{\psi}_0 + a_3^* \psi_0 + a_4^* \bar{\psi}_0 \psi_0$$

Bosonic degree of freedom:

$$X = X_B + \eta \bar{\psi}_0 \psi_0 \quad \text{even } X_B, \eta \in \mathbb{R}$$

Fermionic degree of freedom:

$$\begin{aligned} \psi &= a \psi_0 + b \bar{\psi}_0 \\ \bar{\psi} &= a^* \bar{\psi}_0 + b^* \psi_0 \end{aligned} \quad a, b \in \mathbb{C}$$

Simplest non-trivial model

$$L(\dot{x}, x, \dot{\psi}, \psi, \dot{\bar{\psi}}, \bar{\psi}) := \frac{1}{2} \dot{x}^2 - V_1(x) + \frac{i}{2} (\bar{\psi} \dot{\psi} - \dot{\bar{\psi}} \psi) - V_2(x) \bar{\psi} \psi$$

\uparrow real valued
 \uparrow real valued

Homework: Equations of motion

$$\ddot{x} = -V_1'(x) - V_2'(x) \bar{\psi} \psi$$

$$\dot{\bar{\psi}} = i V_2(x) \bar{\psi} \quad , \quad \dot{\psi} = -i V_2(x) \psi$$

Ansatz: $x(t) = x_B(t) + q(t) \bar{\psi}_0 \psi_0$
 with $\psi(0) = \psi_0 \quad \bar{\psi}(0) = \bar{\psi}_0$

$$\bar{\psi}(t) = \bar{\psi}_0 \exp\left\{ i \int_0^t d\tau V_2(x_B(\tau)) \right\}$$

$$\psi(t) = \psi_0 \exp\left\{ -i \int_0^t d\tau V_2(x_B(\tau)) \right\}$$

note that $V_2(x) = V_2(x_B) + V_2'(x_B) q \bar{\psi}_0 \psi_0$

Homework: Conserved energy $\mathcal{E} = E + F \bar{\psi}_0 \psi_0 \quad , \quad E, F \in \mathbb{R}$

$$\mathcal{E} := \frac{1}{2} \dot{x}^2 + V_1(x) + V_2(x) \bar{\psi}_0 \psi_0 = \text{const.}$$

$$\dot{x} = \dot{x}_B + \dot{q} \bar{\psi}_0 \psi_0 \quad \approx \quad \frac{1}{2} \dot{x}^2 = \frac{1}{2} \dot{x}_B^2 + \dot{x}_B \dot{q} \bar{\psi}_0 \psi_0$$

$$V_1(x) = V_1(x_B) + V_1'(x_B) q \bar{\psi}_0 \psi_0$$

$$V_2(x) \bar{\psi}_0 \psi_0 = V_2(x_B) \bar{\psi}_0 \psi_0$$

Compare body and soul

$$\bullet \quad \boxed{E = \frac{1}{2} \dot{x}_B^2 + V_1(x_B)} \quad \Leftrightarrow \quad \ddot{x}_B = -V_1'(x_B)$$

integrable std. classical dynamics

$$\bullet \quad F = \dot{x}_B \dot{q} + V_1'(x_B) q + V_2(x_B)$$

Solution:

$$\boxed{q(t) = \dot{x}_B(t) \left[\frac{q(0)}{\dot{x}_B(0)} + \int_0^t d\tau \frac{F - V_2(x_B(\tau))}{\dot{x}_B^2(\tau)} \right]}$$

Proof:
$$\dot{q}(t) = \frac{\ddot{x}_B(t)}{\dot{x}_B(t)} q(t) + \dot{x}_B(t) \frac{F - V_2(x_B(t))}{\dot{x}_B^2(t)}$$

$$\begin{aligned} \dot{q}(t) \dot{x}_B(t) &= \ddot{x}_B(t) q(t) + F - V_2(x_B(t)) \\ &= -V_1'(x_B(t)) q(t) + F - V_2(x_B(t)) \quad \# \end{aligned}$$

Conclusion:

1-dim. pseudo-classical mechanics is integrable

Solution is in essence characterised by solution of the std. classical x_B , the body of the bosonic degree of freedom