

Lecture 6

23 Stochastic processes

23.1 Some basics from probability theory

Probability space:

Let

Ω : Sample space (set of events)

\mathcal{A} : σ -algebra of Ω

P : Probability measure $P : \mathcal{A} \rightarrow [0, 1]$

The set (Ω, \mathcal{A}, P) is called *probability space*

$\omega \in \Omega$ is called elementary event

$E \in \mathcal{A}$ is called *event*

$P(E)$ is the probability of event E

Conditions on probability measure:

(i) $0 \leq P(E) \leq 1 \quad \forall E \in \mathcal{A}$

(ii) $P(\Omega) = 1$

(iii) Let E_1, E_2, \dots be disjoint events $\Leftrightarrow E_i \cap E_j = \emptyset \quad \forall i \neq j$

(iv) $\bar{E} := \Omega \setminus E$ "not E " event $\implies P(\bar{E}) = 1 - P(E)$

(v) $P(\emptyset) = 0 \quad \emptyset = \text{impossible event}$

Random variable:

A measurable function

$$X : \Omega \rightarrow \Omega'$$

is called *random variable*. It induces a probability measure (push forward measure)

$$P_X(E') = P(\{\omega \in \Omega \mid X(\omega) \in E'\}), \quad \forall E' \in \Omega'.$$

P_X is called *distribution* of X .

23.2 Stochastic processes

Let T be an index set with $\tau \in T$ and

$$X_\tau : \begin{array}{l} \Omega \rightarrow \mathbb{R}^d = \Omega' \\ \omega \mapsto X_\tau(\omega) \end{array} \quad \forall \tau \in T$$

be a family of \mathbb{R}^d -valued random variables with (ω, \mathcal{A}, P) as probability space.

Then $(\omega, \mathcal{A}, P, \{X_\tau\}_{\tau \in T})$ is called a *stochastic process* with index set T and state (or sample) space \mathbb{R}^d .

Example: Brownian motion

Here $\Omega = \Omega' = \mathbb{R}^d$

$\omega \in \Omega$ denotes an elementary path taken by the particle during time $0 \leq \tau \leq T$.

$X_\tau(\omega) = \omega(\tau)$ position of Brownian particle at time τ = random variable for each τ

$X(\omega)$ for fixed $\omega \in \Omega$ is called *path* or *realisation* of the stochastic process

23.3 Finite-dimensional marginal distribution

The joint probability density

$$P_n(x_n t_n, x_{n-1} t_{n-1}, \dots, x_1 t_1) := \langle \delta(X_{t_n} - x_n) \delta(X_{t_{n-1}} - x_{n-1}) \cdots \delta(X_{t_1} - x_1) \rangle_{P_{X_\tau}}$$

of the random variables X_{t_i} with $t_i \in T$, $i = 1, 2, 3, \dots, n$, $n \in \mathbb{N}$ are called *marginal distributions*.

Properties:

- Positivity: $P_n(x_n t_n, x_{n-1} t_{n-1}, \dots, x_1 t_1) \geq 0$
- Symmetry: $P_n(x_{\pi(n)} t_{\pi(n)}, x_{\pi(n-1)} t_{\pi(n-1)}, \dots, x_{\pi(1)} t_{\pi(1)}) = P_n(x_n t_n, x_{n-1} t_{n-1}, \dots, x_1 t_1)$

$$\text{for all permutations } \begin{pmatrix} 1 & 2 & 3 & \cdots & n \\ \pi(1) & \pi(2) & \pi(3) & \cdots & \pi(n) \end{pmatrix}$$

- Compatability: $\int_{\mathbb{R}^d} d^d x_n P_n(x_n t_n, x_{n-1} t_{n-1}, \dots, x_1 t_1) = P_{n-1}(x_{n-1} t_{n-1}, \dots, x_1 t_1)$
- Normalisation: $\int_{\mathbb{R}^d} d^d x_1 P_1(x_1 t_1) = 1$

Kolmogorov's fundamental (or extension) theorem:

For each family $\{P_n(\cdots)\}_{n \in \mathbb{N}}$ of finite dimensional distributions, obeying above properties, exists a stochastic process $(\omega, \mathcal{A}, P, \{X_\tau\}_{\tau \in T})$ with these marginal distributions. The stochastic process is in essence unique.

In other words, a stochastic process is in essence uniquely defined via its family of marginal distributions.

23.4 Brownian motion

1827 *Robert Brown*: Observed irregular motion of pollen on a static liquid originating from random hits of pollen by molecules of the liquid.

1905 *Albert Einstein*: Provided first correct theoretical description:

- Assume motions within two successive time intervals τ (small) are independent (no correlation of hits)
- Within τ position x varies by amount Δ
- Δ is a random variable with a probability distribution $\Phi(\Delta)$ such that $\Phi(-\Delta) = \Phi(\Delta)$ and $\langle \Delta^2 \rangle \sim \tau$.

Let $f(x, t)$ be the probability distribution density to find particle at position x at time t . Then

$$f(x, t + \tau) = \int d\Delta \Phi(\Delta) f(x + \Delta, t)$$

and for small τ and Δ we have

$$f(x, t + \tau) \approx f(x, t) + \tau \dot{f}(x, t) \\ \int d\Delta \Phi(\Delta) f(x + \Delta, t) \approx f(x, t) \underbrace{\int d\Delta \Phi(\Delta)}_{=1} + f'(x, t) \underbrace{\int d\Delta \Delta \Phi(\Delta)}_{=0} + f''(x, t) \underbrace{\int d\Delta \frac{1}{2} \Delta^2 \Phi(\Delta)}_{=: D\tau}$$

Comparing both sides results in the diffusion equation

$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$$

Einstein: $D = \frac{k_B T}{6\pi\eta a}$ with T temperature, η viscosity and a radius of pollen.

Obviously $\langle \Delta^2 \rangle = 2D\tau$ and with initial condition $f(x, 0) = \delta(x)$ the solution of the diffusion equation is given by the Gaussian distribution

$$f(x, t) = \frac{1}{\sqrt{4\pi D\tau}} e^{-\frac{x^2}{4D\tau}}$$

characterising the probability distribution of X_t .

1907 Paul Langevin: Starts with Newton's eq. for X_t

$$m\ddot{X}_t + \underbrace{6\pi\eta a \dot{X}_t}_{=\text{friction}} = F \quad \leftarrow \quad \text{random force}$$

Multiply by X_t

$$m\ddot{X}_t X_t + 6\pi\eta a \dot{X}_t X_t = F X_t$$

Formally we have $\frac{1}{2} \frac{d^2}{dt^2} X_t^2 = \dot{X}_t^2 + \ddot{X}_t X_t$ resulting in

$$\frac{m}{2} \frac{d^2}{dt^2} X_t^2 - m\dot{X}_t^2 + 3\pi\eta a \frac{d}{dt} X_t^2 = F X_t$$

Now we take the mean value, assume thermodyn. equilibrium $\frac{m}{2} \langle \dot{X}_t^2 \rangle = \frac{1}{2} k_B T$ (1-dim.) and set $x^2(t) := \langle X_t^2 \rangle$

$$\frac{m}{2} \frac{d^2}{dt^2} x^2(t) - k_B T + 3\pi\eta a \frac{d}{dt} x^2(t) = \langle F X_t \rangle \stackrel{\text{independent}}{=} \langle F \rangle \underbrace{\langle X_t \rangle}_{=0} = 0$$

Solution is same as Einstein's: $x^2(t) = \langle X_t^2 \rangle = \frac{k_B T}{3\pi\eta a} t = 2Dt$

23.5 Stationary Markov processes

Conditional probability density: Let $\tau_1 \leq \tau_2 \leq \dots \leq \tau_m \leq t_1 \leq \dots \leq t_n$ then

$$P_{n|m}(x_n t_n, \dots, x_1 t_1 | \xi_m \tau_m, \dots, \xi_1 \tau_1) := \frac{P_{n+m}(x_n t_n, \dots, x_n t_n, \xi_m \tau_m, \dots, \xi_1 \tau_1)}{P_m(\xi_m \tau_m, \dots, \xi_1 \tau_1)}$$

is the *conditional probability density* for X_t taking values x_i at times t_i under condition that the values ξ_j at τ_j were realised by X_t .

Obvious relations:

$$P_1(x_2 t_2) = \int dx_1 P_2(x_2 t_2, x_1 t_1) = \int dx_1 P_{1|1}(x_2 t_2 | x_1 t_1) P_1(x_1 t_1)$$

$$P_{1|1}(x_3 t_3 | x_1 t_1) = \int dx_2 P_{2|1}(x_3 t_3, x_2 t_2 | x_1 t_1) = \int dx_2 P_{1|2}(x_3 t_3 | x_2 t_2, x_1 t_1) P_{1|1}(x_2 t_2 | x_1 t_1)$$

Markov assumption: If

$$P_{n|m}(x_n t_n, \dots, x_n t_n | \xi_m \tau_m, \dots, \xi_1 \tau_1) = P_{n|1}(x_n t_n, \dots, x_n t_n | \xi_m \tau_m)$$

depends only on the **last** condition \iff : **Markov process**

Chapman-Kolmogorov equation:

$$\boxed{P_{1|1}(x_3 t_3 | x_1 t_1) = \int dx_2 P_{1|1}(x_3 t_3 | x_2 t_2) P_{1|1}(x_2 t_2 | x_1 t_1)}$$

Note:

$$P_n(x_n t_n, x_{n-1} t_{n-1}, \dots, x_1 t_1) = P_{1|1}(x_n t_n | x_{n-1} t_{n-1}) P_{1|1}(x_{n-1} t_{n-1} | x_{n-2} t_{n-2}) \dots P_{1|1}(x_2 t_2 | x_1 t_1)$$

$P_{1|1}$ is called the *transition probability density* of the Markov process. If it depends only of the time difference then it is called a *stationary Markov process* with

$$m_{t_2-t_1}(x_2, x_1) := P_{1|1}(x_2 t_2 | x_1 t_1)$$

Stationary Markov process: Is uniquely defined by

- State space: $Q \subseteq \mathbb{R}^d$
- Initial distribution: $P_0(x)$, $x \in Q$, at $t = 0$
- Transition probability density: $m_t(x_2, x_1)$, $t > 0$, obeying below requirements

- $m_t(x_2, x_1) \geq 0$
- $\int_Q dx_2 m_t(x_2, x_1) = 1$
- $\lim_{t \searrow 0} m_t(x_2, x_1) = \delta(x_2 - x_1)$
- $\int_Q dx_2 m_t(x_3, x_2) m_t(x_2, x_1) = m_t(x_3, x_1)$

- If, in addition, the *Lindeberg condition* holds, i.e.,

$$\lim_{t \searrow 0} \frac{1}{t} \int_{|x_2-x_1|>\varepsilon} dx_2 m_t(x_2, x_1) = 0 \quad \forall \varepsilon > 0, x_1 \in Q$$

then this is a *continuous Markov process* who's paths (realisations) are almost sure continuous (no jumps)

Realisations of a continuous stationary markov process, with initial distribution $P_0(x) = \delta(x - q_a)$ are all continuous open-end paths living in Q and starting in $q_a \in Q$.

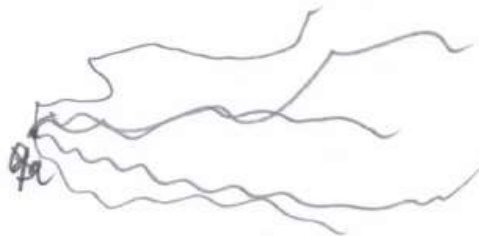
We denote this set of paths by the *path space* $\mathcal{C}(Q, q_a)$.

From now on we use the notation $M_t(\omega) \equiv X_t(\omega) \in \mathcal{C}(Q, q_a)$ for such Markov processes.

23.6 Path integral representation of stationary Markov processes

Ingredients:

- Continuous stationary markov process $M_t(\omega)$, $\omega \in \Omega$, $t \geq 0$
- State (or confinguration) space $Q \subseteq \mathbb{R}^d$
- Initial distribution $P_0(q) = \delta(q - q_a)$
- Path space $\mathcal{C}(Q, q_a)$, all open-end paths $M_t(\omega)$ starting at q_a



Path: for fixed $\omega \in \Omega$ with $M_0(\omega) = q_a$

$$x : \begin{array}{l} \mathbb{R}_+^0 \rightarrow Q \\ t \mapsto M_t(\omega) \end{array}$$

Probability measure: P

- Initial condition: $P(\{\omega \in \Omega | M_0(\omega) = q_a\}) = 1$
- Continuity: $P(\{\omega \in \Omega | \lim_{t' \rightarrow t} M_{t'}(\omega) = M_t(\omega)\}) = 1$
- Marginal distribution:

$$P(\{\omega \in \Omega | M_{t_n}(\omega) = q_n, \dots, M_{t_1}(\omega) = q_1\}) = m_{t_n - t_{n-1}}(q_n, q_{n-1}) \cdots m_{t_1 - t_0}(q_1, q_a)$$

with

$$0 = t_0 < t_1 < \cdots < t_n, \quad q_i \in Q, \quad i = 1, 2, \dots, n, \quad \forall n \in \mathbb{N}.$$

- Induced mapping:

$$\phi : \begin{array}{l} \Omega \rightarrow \mathcal{C}(Q, q_a) \\ \omega \mapsto x \text{ Realisation} \end{array}$$

induces a probability measure on $\mathcal{C}(Q, q_a)$: $M := P \circ \phi^{-1}$

The M -path integral:

Consider arbitrary real-valued functional F

$$F : \begin{array}{l} \mathcal{C}(Q, q_a) \rightarrow \mathbb{R} \\ x \mapsto F[x] \end{array}$$

Then the expectation value is given by the M -path integral

$$\langle F[M_t] \rangle = \int_{\mathcal{C}(Q, q_a)} dM[x] F[x]$$

The measure can explicitly be defined via so-called cylinder functionals

$$F[x] = f_n(x(t_n), \dots, x(t_1)) \quad 0 < t_1 < \cdots < t_n, \quad n \in \mathbb{N}$$

as follows

$$\begin{aligned} & \int_{\mathcal{C}(Q, q_a)} dM[x] f_n(x(t_n), \dots, x(t_1)) \\ & := \int_Q dq_n \cdots \int_Q dq_1 m_{t_n - t_{n-1}}(q_n, q_{n-1}) \cdots m_{t_1 - t_0}(q_1, q_a) f_n(q_n, \dots, q_1) \end{aligned}$$

This is similar to the definition of a Riemann or Lebesgue integral.

Extension to general functional is then via approximation by suitable cylinder functions, $F = \lim_{n \rightarrow \infty} f_n$, i.e. via a time-lattice approximation

$$\int_{\mathcal{C}(Q, q_a)} dM[x] F[x] = \lim_{n \rightarrow \infty} \int_{\mathcal{C}(Q, q_a)} dM[x] f_n(x(t_n), \dots, x(t_1))$$

Properties:

- $\int_{\mathcal{C}(Q, q_a)} dM[x] = 1$

- $\int_{\mathcal{C}(Q, q_a)} dM[x] \delta(x(t_n) - q_n) \cdots \delta(x(t_1) - q_1) = m_{t_n - t_{n-1}}(q_n, q_{n-1}) \cdots m_{t_1 - t_0}(q_1, q_a)$

Generator T_M : Is local operator on dense subset of $L^2(Q)$ defined by its q' -representation

$$T_M^{(q')} \psi(q') = \langle q' | T_M | \psi \rangle \quad \text{via} \quad \langle q' | e^{-tT_M} | q \rangle := m_t(q', q)$$

It obeys the

Fokker-Planck equation:

$$-\partial_t m_t(q', q) = T_M^{(q')} m_t(q', q)$$

Typically structure:

$$T_M = \frac{\vec{P}^2}{2} + i\vec{P} \cdot \vec{m}(\vec{Q})$$

Diffusion constant $D = \frac{1}{2}$.

Drift $\vec{m}(\vec{Q})$ due to external force

Feynman-Kac formula: Let $H := T_M + V(Q)$ then

$$\boxed{\langle q_b | e^{-tH} | q_a \rangle = \int_{\mathcal{C}(Q, q_a)} dM[x] \delta(x(t) - q_b) \exp \left\{ - \int_0^t d\tau V(x(\tau)) \right\}}$$

Conclusion: There are at least three equivalent ways (generator, transition density, path integral) to characterise a Markov process, they are related by

$$\langle q' | e^{-tT_M} | q \rangle = m_t(q', q) = \int_{\mathcal{C}(Q, q)} dM[x] \delta(x(t) - q')$$

Examples:

Wiener process

Bessel process

Legendre process