

3. Homework

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Problem 1: According to Section 4.2 for $\epsilon = -\frac{1}{2} - 2i$, $\beta = 0$ the solution reads

$$u(x) = e^{x^2/2} H_{2k}(ix) \quad \text{here } k=1 \sim H_2(z) = 4z^2 - 2 \quad \text{2. Hermite polynomial}$$

Hence $u(x) = (-4x^2 - 2) e^{x^2/2}$

$$u'(x) = x u(x) + e^{-x^2/2} (-8x) = x u(x) + \frac{8x}{4x^2+2} u(x)$$

$$\leadsto \frac{u'(x)}{u(x)} = x + \frac{4x}{2x^2+1}$$

With $V_-(x) = \left(\frac{u'(x)}{u(x)} \right)^2 - \frac{x^2}{2} + 2\epsilon$

$$= \left(x + \frac{4x}{2x^2+1} \right)^2 - \frac{x^2}{2} - 5 = \frac{x^2}{2} + \frac{8x^2}{2x^2+1} + \left(\frac{4x}{2x^2+1} \right)^2 - 5$$

$$\text{Spec}(H_-) = \left\{ -\frac{5}{2}, \frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots \right\}$$

Compare with general formula for arbitrary $k \in \mathbb{N}$

$$V_-(x) = \frac{x^2}{2} + 2ix \frac{H_{2k+1}(ix)}{H_{2k}(ix)} - \left(\frac{H_{2k+1}(ix)}{H_{2k}(ix)} \right)^2 - 4k - 1$$

Note $H_3(ix) = -i8x^3 - i12x$

$$2ix H_3(ix) = 16x^4 + 24x^2 = 8x^2(2x^2+3)$$

$$\leadsto \frac{2ix H_3(ix)}{H_2(ix)} = 8x^2 \frac{2x^2+3}{4x^2-2} = -4x^2 \frac{2x^2+3}{2x^2+1}$$

$$\left(\frac{H_3(ix)}{H_2(ix)} \right)^2 = \left(\frac{8ix^3 + 12ix}{4x^2-2} \right)^2 = -4x^2 \left(\frac{2x^2+3}{2x^2+1} \right)^2$$

$$V_-(x) = \frac{x^2}{2} - 4x^2 \frac{2x^2+3}{2x^2+1} + 4x^2 \frac{(2x^2+3)^2}{(2x^2+1)^2} - 5$$

$$= \frac{x^2}{2} + \frac{4x^2}{(2x^2+1)^2} (2x^2+3)(2x^2+3 - 2x^2-1) = \frac{x^2}{2} + \frac{4x^2}{(2x^2+1)^2} (2x^2+3)2 - 5$$

$$= \frac{x^2}{2} + \frac{8x^2}{(2x^2+1)^2} (2x^2+1+2) - 5 = \frac{x^2}{2} + \frac{8x^2}{2x^2+1} + \frac{16x^2}{(2x^2+1)^2} \quad \#$$

Problem 2: Schrödinger like eq for $V_+(x) = 0$ and $\epsilon = -\frac{\hbar^2 k^2}{2m} < 0 = E_0$ (2)

$$-\frac{\hbar^2}{2m} \partial_x^2 u(x) = -\frac{\hbar^2 k^2}{2m} u(x) \Rightarrow u''(x) = k^2 u(x)$$

General solution $u(x) = \alpha e^{kx} + \beta e^{-kx}$ or w/o loss of generality $\alpha = 1$

$$\underline{u(x) = e^{kx} + \beta e^{-kx}}$$

Check: $u(x) = 0 \leadsto e^{kx} = -\beta e^{-kx} \leadsto \beta = -e^{2kx} \leq 0$

Hence $\underline{\beta > 0} \Rightarrow u(x) > 0$

Potential:

$$V_-(x) = \frac{\hbar^2}{m} \left(\frac{u'(x)}{u(x)} \right)^2 - V_+(x) + 2\epsilon, \quad u'(x) = k(e^{kx} - \beta e^{-kx})$$

$$= \frac{\hbar^2 k^2}{m} \left(\frac{e^{kx} - \beta e^{-kx}}{e^{kx} + \beta e^{-kx}} \right)^2 - \frac{\hbar^2 k^2}{m}$$

$$= \frac{\hbar^2 k^2}{m} \left[\left(\frac{e^{kx} - \beta e^{-kx}}{e^{kx} + \beta e^{-kx}} \right)^2 - 1 \right] \quad \#$$

has 1 bound state $\epsilon = -\frac{\hbar^2 k^2}{2m}$

with $\psi_e^-(x) = \frac{C}{e^{kx} + \beta e^{-kx}} \quad \#$

Special case $\beta = 1$:

$$\frac{e^{kx} - e^{-kx}}{e^{kx} + e^{-kx}} = \tanh kx$$

$$\leadsto V_-(x) = \frac{\hbar^2 k^2}{m} (\tanh^2 kx - 1) = \frac{\hbar^2 k^2}{2m} \frac{2}{\cosh^2 kx}$$

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Problem 3:

$$E[\phi] = \int dx \left[\frac{1}{2} \dot{\phi}^2 + \frac{1}{2} \phi'^2 + U(\phi) \right]$$

$$\frac{dE}{dt} = \int dx \left[\frac{1}{2} \partial_t \dot{\phi}^2 + \frac{1}{2} \partial_t \phi'^2 + \partial_t U(\phi) \right]$$

$$= \int dx \left[\frac{1}{2} 2 \dot{\phi} \ddot{\phi} + \frac{1}{2} 2 \phi' \dot{\phi}' + U'(\phi) \dot{\phi} \right]$$

$$= \int dx \left[\dot{\phi} \ddot{\phi} + \phi' \dot{\phi}' + U'(\phi) \dot{\phi} \right]$$

$$\int dx \phi' \dot{\phi}' = \underbrace{\phi' \dot{\phi}'}_{=0} \Big|_{-\infty}^{+\infty} - \int dx \phi'' \dot{\phi}$$

$$= \int dx \left[\underbrace{\dot{\phi} \ddot{\phi} - \phi'' \dot{\phi}}_{=0} + U'(\phi) \dot{\phi} \right] \dot{\phi} = 0$$

Problem 4:

$$W(x) = \text{sgn } x \quad \leadsto \quad V_-(x) = W^2(x) - W'(x) = 1 - 2\delta(x)$$

Hint: $\int_{-\infty}^z dx \, d_x \text{sgn } x = \text{sgn } x \Big|_{-\infty}^z = \text{sgn } z + 1 = \begin{cases} 0 & z < 0 \\ 2 & z > 0 \end{cases}$

$$\int_{-\infty}^z dx \, 2\delta(x) = \begin{cases} 0 & z < 0 \\ 2 & z > 0 \end{cases} \quad \neq$$

$$\psi_0(x) = N \exp\left\{-\int dx W(x)\right\} = N e^{-|x|}$$

Proof: $\int_0^z dx \, \text{sgn } x = \begin{cases} z & z > 0 \\ -(-z) & z < 0 \end{cases} = |z|$

$$\phi_{st}(x) = \int dx \, \psi_0(x) = \underline{\underline{\text{sgn } x (1 - e^{-|x|})}} \quad \text{with } N=1 \quad \underline{\underline{\phi_{\pm} = \pm 1}}$$

Proof: $\phi_{st}'(x) = \underbrace{2\delta(x)(1 - e^{-|x|})}_{=0} + \text{sgn } x (e^{-|x|}) \text{sgn } x = e^{-|x|}$

Hence: $U(\phi_{st}) = \frac{1}{2} (\phi_{st}'(x))^2 = \frac{1}{2} e^{-2|x|}$

Consider $|\phi_{st}(x)| = 1 - e^{-|x|} \leadsto e^{-|x|} = 1 - |\phi_{st}(x)|$

$$\Rightarrow U(\phi_{st}) = \frac{1}{2} (1 - |\phi_{st}(x)|)^2$$

analytically continue beyond $\phi_{\pm} = \pm 1$

$$\underline{\underline{U(\phi) = \frac{1}{2} (1 - |\phi|)^2}}$$

