

Problem 11: Generators of $SO(3)$ in $L^2(\mathbb{R}^3)$

a) Use result of problem 10c)

$$\exp\{-i\delta\alpha^a L_a\} = 1 + L_a^2(\cos \delta\alpha^a - 1) - iL_a \sin \delta\alpha^a \approx 1 - iL_a \delta\alpha^a + O((\delta\alpha)^2)$$

Hence

$$g(\vec{\delta\alpha}) \approx 1 - i \sum_{a=1}^3 \delta\alpha^a L_a = 1 + \begin{pmatrix} 0 & -\delta\alpha^3 & \delta\alpha^2 \\ \delta\alpha^3 & 0 & -\delta\alpha^1 \\ -\delta\alpha^2 & \delta\alpha^1 & 0 \end{pmatrix}$$

b)

$$\delta\vec{x} = \vec{x}' - \vec{x} = \begin{pmatrix} 0 & -\delta\alpha^3 & \delta\alpha^2 \\ \delta\alpha^3 & 0 & -\delta\alpha^1 \\ -\delta\alpha^2 & \delta\alpha^1 & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} -x_2\delta\alpha^3 + x_3\delta\alpha^2 \\ x_1\delta\alpha^3 - x_3\delta\alpha^1 \\ -x_1\delta\alpha^2 + x_2\delta\alpha^1 \end{pmatrix} =: \begin{pmatrix} \delta x_1 \\ \delta x_2 \\ \delta x_3 \end{pmatrix}$$

Hence

$$\left. \begin{aligned} X_1 &= - \sum_{k=1}^3 \frac{\delta x_k}{\delta\alpha_1} \partial_k = x_3\partial_2 - x_2\partial_3 \\ X_2 &= - \sum_{k=1}^3 \frac{\delta x_k}{\delta\alpha_2} \partial_k = x_1\partial_3 - x_3\partial_1 \\ X_3 &= - \sum_{k=1}^3 \frac{\delta x_k}{\delta\alpha_3} \partial_k = x_2\partial_1 - x_1\partial_2 \end{aligned} \right\} \vec{X} = -\vec{x} \times \vec{\nabla} = \frac{1}{i\hbar} \vec{x} \times \vec{p} = \frac{1}{i\hbar} \vec{L}$$

c) As $X_i = -\varepsilon_{ijk} x_j \partial_k$ (sum over repeated indices) \implies

$$[X_i, X_j] = \varepsilon_{ilm} \varepsilon_{jkn} [x_l \partial_m, x_k \partial_n]$$

Consider

$$\begin{aligned} [x_l \partial_m, x_k \partial_n] &= x_l \partial_m x_k \partial_n - x_k \partial_n x_l \partial_m \\ &= x_l (\delta_{km} + \underline{x_k \partial_m}) \partial_n - x_k (\delta_{nl} + \underline{x_l \partial_n}) \partial_m \\ &= \delta_{km} x_l \partial_n - \delta_{nl} x_k \partial_m \end{aligned}$$

Hence

$$\begin{aligned}
[X_i, X_j] &= \varepsilon_{ilm} \varepsilon_{jkn} (\delta_{km} x_l \partial_n - \delta_{nl} x_k \partial_m) \\
&= \varepsilon_{ilk} \varepsilon_{jkn} x_l \partial_n - \varepsilon_{ilm} \varepsilon_{jkl} x_k \partial_m \\
&\quad \text{use } \varepsilon_{jkn} = -\varepsilon_{jnk} \quad \text{and} \quad \varepsilon_{ilm} = -\varepsilon_{iml} \\
&= \varepsilon_{iml} \varepsilon_{jkl} x_k \partial_m - \varepsilon_{ilk} \varepsilon_{jnk} x_l \partial_n \\
&\quad \text{use } \varepsilon_{rst} \varepsilon_{uvt} = (\delta_{ru} \delta_{sv} - \delta_{rv} \delta_{su}) \\
&= \left(\delta_{ij} \delta_{mk} - \delta_{ik} \delta_{mj} \right) x_k \partial_m - \left(\delta_{ij} \delta_{ln} - \delta_{in} \delta_{lj} \right) x_l \partial_n \\
&= \delta_{in} \delta_{lj} x_l \partial_n - \delta_{ik} \delta_{mj} x_k \partial_m \\
&\quad \text{replace in 1. term } n \rightarrow m, \quad l \rightarrow k \\
&= \underbrace{(\delta_{im} \delta_{jk} - \delta_{ik} \delta_{jm})}_{\varepsilon_{ijr} \varepsilon_{mkr}} x_k \partial_m \\
&= -\varepsilon_{ijr} \varepsilon_{rkm} x_k \partial_m = \varepsilon_{ijr} X_r \quad \Longrightarrow \quad \boxed{c_{ij}{}^r = \varepsilon_{ijr}} \quad \#
\end{aligned}$$

Cartan metric:

$$\begin{aligned}
g_{kl} &:= c_{kr}{}^s c_{ls}{}^r = \varepsilon_{krs} \varepsilon_{lsr} = -\varepsilon_{krs} \varepsilon_{lrs} \\
&= -(\delta_{kl} \delta_{rr} - \delta_{kr} \delta_{rl}) \\
&= -(3\delta_{kl} - \delta_{kl}) = -2\delta_{kl} \quad \#
\end{aligned}$$

Metric tensors:

$$(g_{kl}) = \begin{pmatrix} -2 & 0 & 0 \\ 0 & -2 & 0 \\ 0 & 0 & -2 \end{pmatrix} \quad (g^{kl}) = \begin{pmatrix} -\frac{1}{2} & 0 & 0 \\ 0 & -\frac{1}{2} & 0 \\ 0 & 0 & -\frac{1}{2} \end{pmatrix}$$

Casimir operator:

$$C := g^{kl} X_k X_l = -\frac{1}{2} \vec{X}^2 = \frac{1}{2\hbar} \vec{L}^2$$

Problem 12: Generators of $SO(4)$ in $L^2(\mathbb{R}^4)$

General considerations:

Let $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n$, x_μ and ∂_μ denote coordinate and derivative, respectively.

Consider operator

$$L_{\mu\nu} := x_\mu \partial_\nu - x_\nu \partial_\mu$$

Obviously, $L_{\mu\nu}$ is generator of rotation in $x_\mu - x_\nu$ -plane

\implies set of all $L_{\mu\nu} = -L_{\nu\mu}$ generates rotations in \mathbb{R}^n

There exist $n(n-1)/2$ independent generators of $SO(n)$.

Derivation of the $so(n)$ algebra:

$$\begin{aligned} [L_{\mu\nu}, L_{\rho\sigma}] &= (x_\mu \partial_\nu - x_\nu \partial_\mu)(x_\rho \partial_\sigma - x_\sigma \partial_\rho) - (x_\rho \partial_\sigma - x_\sigma \partial_\rho)(x_\mu \partial_\nu - x_\nu \partial_\mu) \\ &= x_\mu \partial_\nu x_\rho \partial_\sigma - x_\mu \partial_\nu x_\sigma \partial_\rho - x_\nu \partial_\mu x_\rho \partial_\sigma + x_\nu \partial_\mu x_\sigma \partial_\rho \\ &\quad - x_\rho \partial_\sigma x_\mu \partial_\nu + x_\rho \partial_\sigma x_\nu \partial_\mu + x_\sigma \partial_\rho x_\mu \partial_\nu - x_\sigma \partial_\rho x_\nu \partial_\mu \\ &= x_\mu \left(\delta_{\nu\rho} + \underline{x_\rho \partial_\nu} \right) \partial_\sigma - x_\mu \left(\delta_{\nu\sigma} + \underline{x_\sigma \partial_\nu} \right) \partial_\rho - x_\nu \left(\delta_{\mu\rho} + \underline{x_\rho \partial_\mu} \right) \partial_\sigma + x_\nu \left(\delta_{\mu\sigma} + \underline{x_\sigma \partial_\mu} \right) \partial_\rho \\ &\quad - x_\rho \left(\delta_{\mu\sigma} + \underline{x_\mu \partial_\sigma} \right) \partial_\nu + x_\rho \left(\delta_{\nu\sigma} + \underline{x_\nu \partial_\sigma} \right) \partial_\mu + x_\sigma \left(\delta_{\mu\rho} + \underline{x_\mu \partial_\rho} \right) \partial_\nu - x_\sigma \left(\delta_{\nu\rho} + \underline{x_\nu \partial_\rho} \right) \partial_\mu \\ &= \delta_{\nu\rho} (x_\mu \partial_\sigma - x_\sigma \partial_\mu) - \delta_{\nu\sigma} (x_\mu \partial_\rho - x_\rho \partial_\mu) - \delta_{\mu\rho} (x_\nu \partial_\sigma - x_\sigma \partial_\nu) + \delta_{\mu\sigma} (x_\nu \partial_\rho - x_\rho \partial_\nu) \\ &= \delta_{\nu\rho} L_{\mu\sigma} - \delta_{\nu\sigma} L_{\mu\rho} - \delta_{\mu\rho} L_{\nu\sigma} + \delta_{\mu\sigma} L_{\nu\rho} \end{aligned}$$

a) Now $n = 4$ and $(x_1, x_2, x_3, x_4) = (x, y, z, t) \in \mathbb{R}^4$

We identify

$$\begin{array}{lll} M_1 = L_{32} & M_2 = L_{13} & M_3 = L_{21} \\ N_1 = L_{14} & N_2 = L_{24} & N_3 = L_{34} \end{array}$$

Now we can use above result of $so(n)$ algebra:

$$\left. \begin{array}{l} [M_1, M_2] = [L_{32}, L_{13}] = L_{21} = M_3 \\ [M_2, M_3] = [L_{13}, L_{21}] = L_{32} = M_1 \\ [M_3, M_1] = [L_{21}, L_{32}] = L_{13} = M_2 \end{array} \right\} [M_i, M_j] = \varepsilon_{ijk} M_k$$

$$\left. \begin{array}{l} [M_1, N_1] = [L_{32}, L_{14}] = 0 \\ [M_2, N_2] = [L_{13}, L_{24}] = 0 \\ [M_3, N_3] = [L_{21}, L_{34}] = 0 \end{array} \right\} [M_i, N_i] = 0$$

$$\left. \begin{aligned} [M_1, N_2] &= [L_{32}, L_{24}] = L_{34} = N_3 \\ [M_2, N_3] &= [L_{13}, L_{34}] = L_{14} = N_1 \\ [M_3, N_1] &= [L_{21}, L_{14}] = L_{24} = N_2 \\ [M_1, N_3] &= [L_{32}, L_{34}] = L_{24} = -N_3 \\ [M_2, N_1] &= [L_{13}, L_{14}] = L_{34} = -N_1 \\ [M_3, N_2] &= [L_{21}, L_{24}] = L_{14} = -N_2 \end{aligned} \right\} [M_i, N_j] = \varepsilon_{ijk} N_k$$

$$[N_i, N_j] = [L_{i4}, L_{j4}] = -L_{ij} = L_{ji} = \varepsilon_{ijk} M_k$$

b) Let $J_k := \frac{1}{2}(M_k + N_k)$

$$\begin{aligned} 4[J_k, J_l] &= [M_k + N_k, M_l + N_l] \\ &= [M_k, M_l] + [M_k, N_l] + [N_k, M_l] + [N_k, N_l] \\ &= \varepsilon_{klm} M_m + \varepsilon_{klm} N_m - \varepsilon_{lkm} N_m + \varepsilon_{klm} M_m \\ &= \varepsilon_{klm} (2M_m + 2N_m) = 4\varepsilon_{klm} J_m \end{aligned}$$

$$\implies [J_k, J_l] = \varepsilon_{klm} J_m \quad so(3) - \text{algebra}$$

Let $K_i := \frac{1}{2}(M_i - N_i)$

$$\begin{aligned} 4[K_i, K_j] &= [M_i - N_i, M_j - N_j] \\ &= [M_i, M_j] - [M_i, N_j] - [N_i, M_j] + [N_i, N_j] \\ &= \varepsilon_{ijk} M_k - \varepsilon_{ijk} N_k + \varepsilon_{jik} N_k + \varepsilon_{ijk} M_k \\ &= 2\varepsilon_{ijk} (M_k - N_k) = 4\varepsilon_{ijk} K_k \end{aligned}$$

$$\implies [K_i, K_j] = \varepsilon_{ijk} K_k \quad so(3) - \text{algebra}$$

Finally we show the decoupling of both subalgebras

$$\begin{aligned} [J_i, K_j] &= [M_i + N_i, M_j - N_j] \\ &= [M_i, M_j] - [M_i, N_j] + [N_i, M_j] - [N_i, N_j] \\ &= \varepsilon_{ijk} M_k - \varepsilon_{ijk} N_k - \varepsilon_{jik} N_k - \varepsilon_{ijk} M_k \\ &= 0 \end{aligned}$$

$$\implies so(4) = so(3) \oplus so(3)$$

Casimir operators of $so(4)$

We have two Casimir from both $so(3)$ algebras:

$$\vec{J}^2 = J_1^2 + J_2^2 + J_3^2, \quad \vec{K}^2 = K_1^2 + K_2^2 + K_3^2$$

The corresponding labels of their UIR are $j, k = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$

Hence, the UIR of $so(4)$ can be denoted by the pair (j, k)

Some relations

$$\begin{aligned} \vec{J}^2 &= \frac{1}{4}(\vec{M} + \vec{N})^2 = \frac{1}{4}(\vec{M}^2 + \vec{M} \cdot \vec{N} + \vec{N} \cdot \vec{M} + \vec{N}^2) \\ &\quad \text{but } \vec{M} \cdot \vec{N} = \vec{N} \cdot \vec{M} \text{ as } [M_i, N_i] = 0 \\ &= \frac{1}{4}(\vec{M}^2 + 2\vec{M} \cdot \vec{N} + \vec{N}^2) \end{aligned}$$

$$\vec{K}^2 = \frac{1}{4}(\vec{M} - \vec{N})^2 = \frac{1}{4}(\vec{M}^2 - 2\vec{M} \cdot \vec{N} + \vec{N}^2)$$

$$\vec{J}^2 - \vec{K}^2 = \vec{M} \cdot \vec{N}$$

$$\vec{J}^2 + \vec{K}^2 = \frac{1}{4}(\vec{M}^2 + \vec{N}^2)$$