



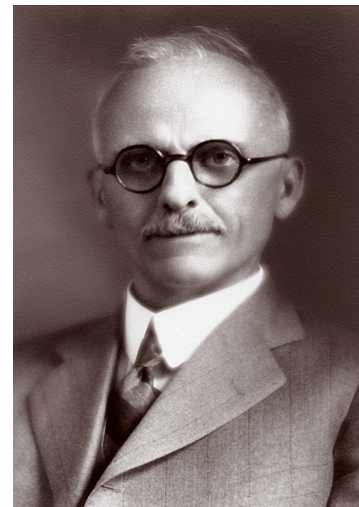
# $H_0$ and the never-ending story of the expansion rate of the Universe

Bruno Leibundgut  
ESO

# Great Debate: What is the size of the Universe?

Presentations at the Annual Meeting of the National Academy of Sciences in Washington DC, 26. April 1920

Harlow Shapley vs. Heber Curtis



<http://incubator.rockefeller.edu/geeks-of-the-week-harlow-shapley-heber-curtis/>



# Models of the Milky Way

## Jacobus Kapteyn (1922)

- based on stellar counts in selected areas
- Kapteyn contributed many stars in the southern hemisphere

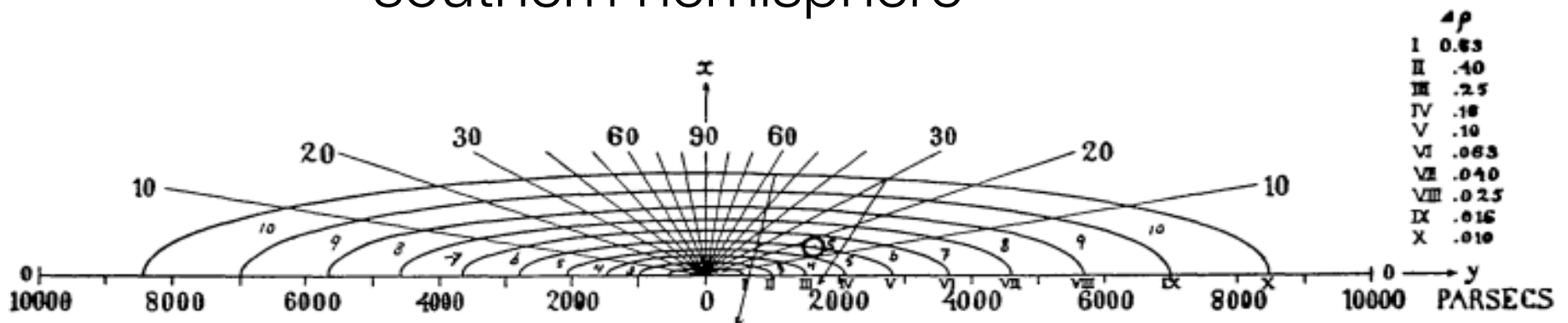
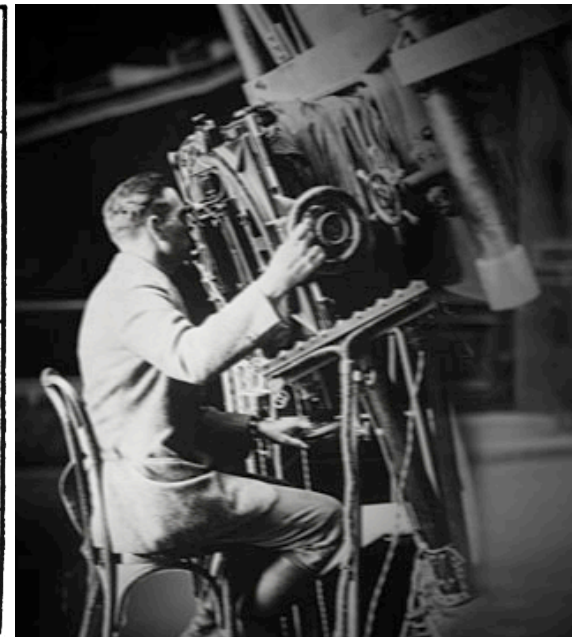
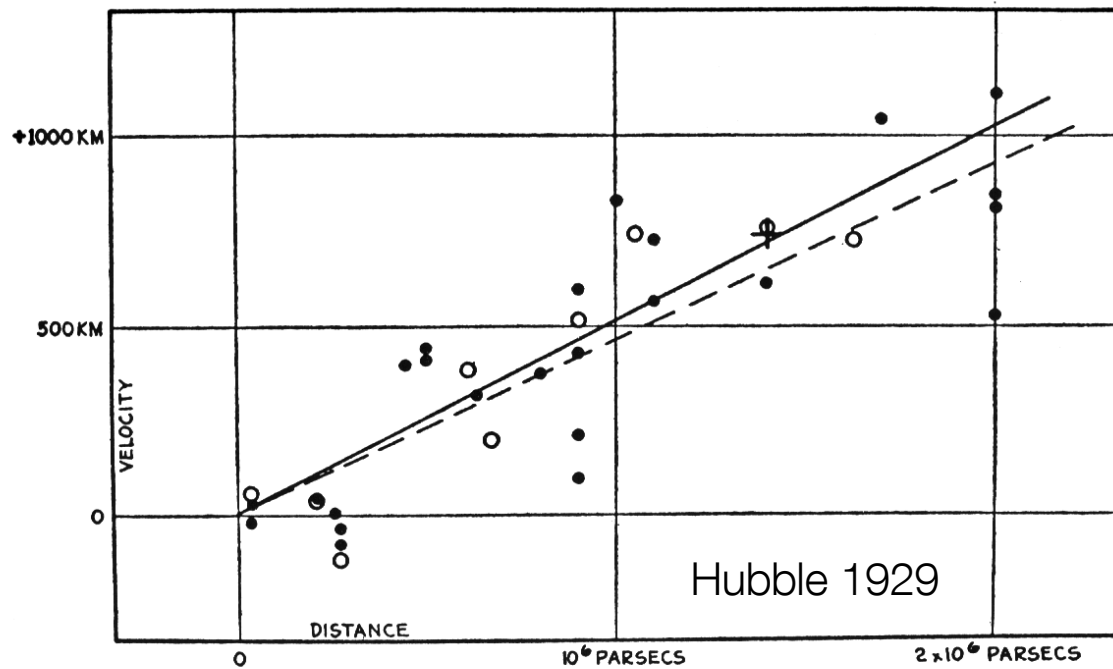


FIG. I

# Background

Expanding universe

→ expansion rate critical for cosmic evolution



STScI

FIG. 9. *The Formulation of the Velocity-Distance Relation.*

# Leading Theory of the Universe

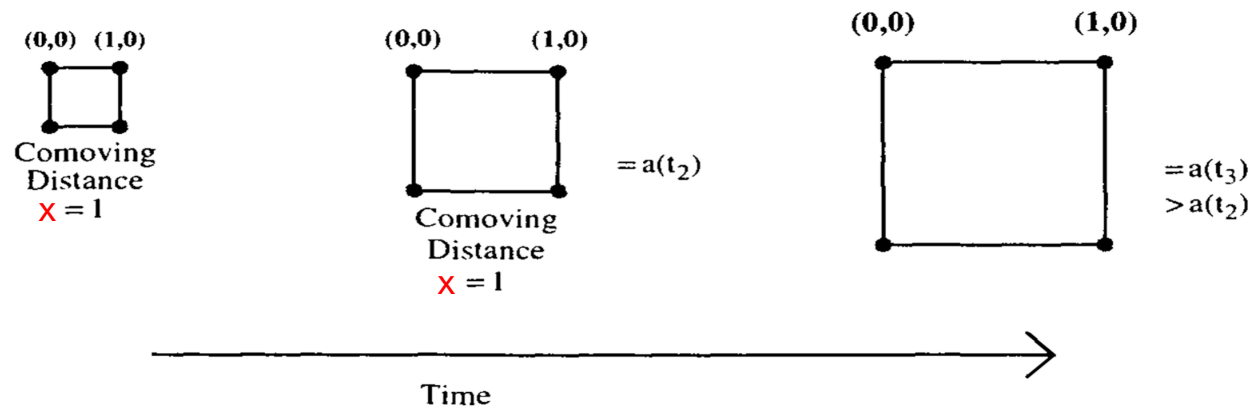


# Dealing with an expanding Universe

## Cosmic Distances

Separate the observed distances  $r(t)$  into the expansion factor  $a(t)$  and the fixed part  $x$  (called *comoving* distance)

$$r(t) = a(t)x$$



# Friedmann Equation

Time evolution of the scale factor is described through the time part of the Einstein equations

Assume a metric for a homogeneous and isotropic universe and a perfect fluid

$$\frac{\dot{a}^2}{a^2} + \frac{k}{a^2} = \frac{8\pi G}{3} \rho(t)$$

$$g_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & a^2 & 0 & 0 \\ 0 & 0 & a^2 & 0 \\ 0 & 0 & 0 & a^2 \end{pmatrix}$$



$$T^{\mu\nu} = \begin{pmatrix} \rho c^2 & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}$$

# Friedmann Equation

Put the various densities into the Friedmann equation

$$\frac{\dot{a}^2}{a^2} = H^2 = \frac{8\pi G}{3} \rho(t) - \frac{k}{a^2} = \frac{8\pi G}{3} (\rho_M + \rho_\gamma + \rho_{vac}) - \frac{k}{a^2}$$

Use the critical density  $\rho_{crit} = \frac{3H_0^2}{8\pi G} \approx 2 \cdot 10^{-29} \text{ g cm}^{-3}$

(flat universe),

define the ratio to the critical density  $\Omega = \frac{\rho}{\rho_{crit}}$

Most compact form of Friedmann equation

$$1 = \Omega_M + \Omega_\gamma + \Omega_{vac} + \Omega_k$$

with  $\Omega_k = -\frac{k}{a^2 H^2}$



# Dependence on Scale Parameter

For the different contents there were different dependencies for the scale parameter

$$\rho_M \propto a^{-3} \quad \rho_\gamma \propto a^{-4} \quad \rho_{vac} = \text{const}$$

Combining this with the critical densities we can write the density as

$$\rho = \frac{3H_0^2}{8\pi G} \left[ \Omega_M \left( \frac{a_0}{a} \right)^3 + \Omega_\gamma \left( \frac{a_0}{a} \right)^4 + \Omega_\Lambda + \Omega_k \left( \frac{a_0}{a} \right)^2 \right]$$

and the Friedmann equation

$$H^2 = H_0^2 \left[ \Omega_M (1+z)^3 + \Omega_\gamma (1+z)^4 + \Omega_\Lambda + \Omega_k (1+z)^2 \right]$$

# History of $H_0$

## Expansion rate by G. Lemaître (1927)

de l'observateur. En effet, la période de la lumière émise dans des conditions physiques semblables doit être partout la même lorsqu'elle est exprimée en temps propre.

$$\frac{v}{c} = \frac{\delta t_2}{\delta t_1} - 1 = \frac{R_2}{R_1} - 1 \quad (22)$$

mesure donc l'effet Doppler apparent dû à la variation du rayon de l'univers. *Il est égal à l'excès sur l'unité du rapport des rayons de l'univers à l'instant où la lumière est reçue et à l'instant où elle est émise.*  $v$  est la vitesse de l'observateur qui produirait le même effet. Lorsque la source est suffisamment proche nous pouvons écrire approximativement

$$\frac{v}{c} = \frac{R_2 - R_1}{R_1} = \frac{dR}{R} = \frac{R'}{R} dt = \frac{R'}{R} r$$

où  $r$  est la distance de la source. Nous avons donc

Footnote!

(<sup>2</sup>) En ne donnant pas de poids aux observations, on trouverait 670 Km./sec à  $1,16 \times 10^6$  parsecs, 575 Km./sec à  $10^6$  parsecs. Certains auteurs ont cherché à mettre en

# Intermezzo

## Age of the Universe

Matter-dominated universe has the following age

$$t_0 = \frac{2}{3H_0}$$

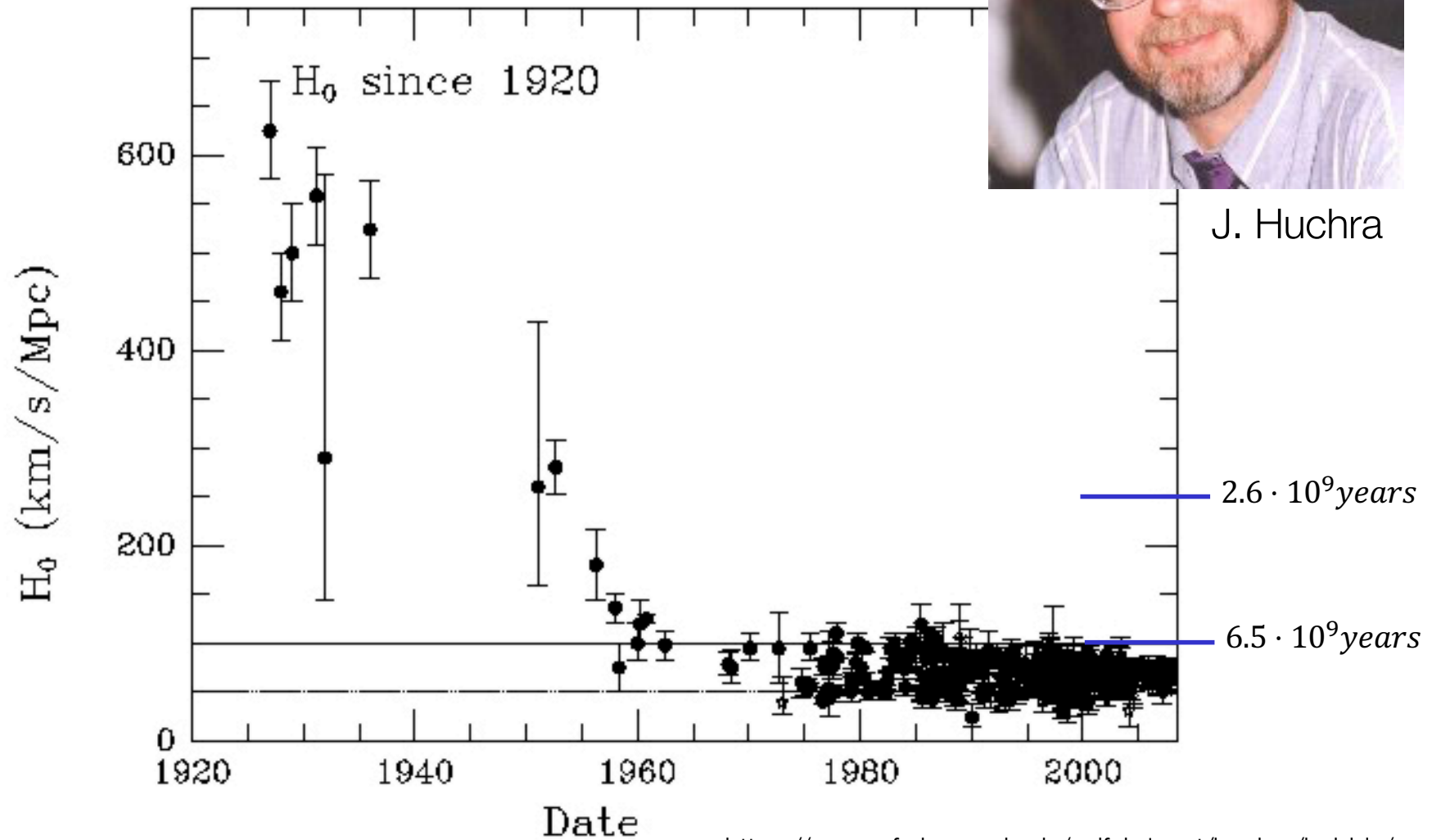
$H_0$ (km/s/Mpc)	$t_0$ (yr)
500	$1.30 \cdot 10^9$
250	$2.61 \cdot 10^9$
100	$6.52 \cdot 10^9$
80	$8.15 \cdot 10^9$
70	$9.32 \cdot 10^9$
60	$1.09 \cdot 10^{10}$
50	$1.30 \cdot 10^{10}$
30	$2.17 \cdot 10^{10}$

- age of the Earth:  $4.5 \cdot 10^9$  years
- oldest stars:  $\sim 1.2 \cdot 10^{10}$  years

# History of $H_0$

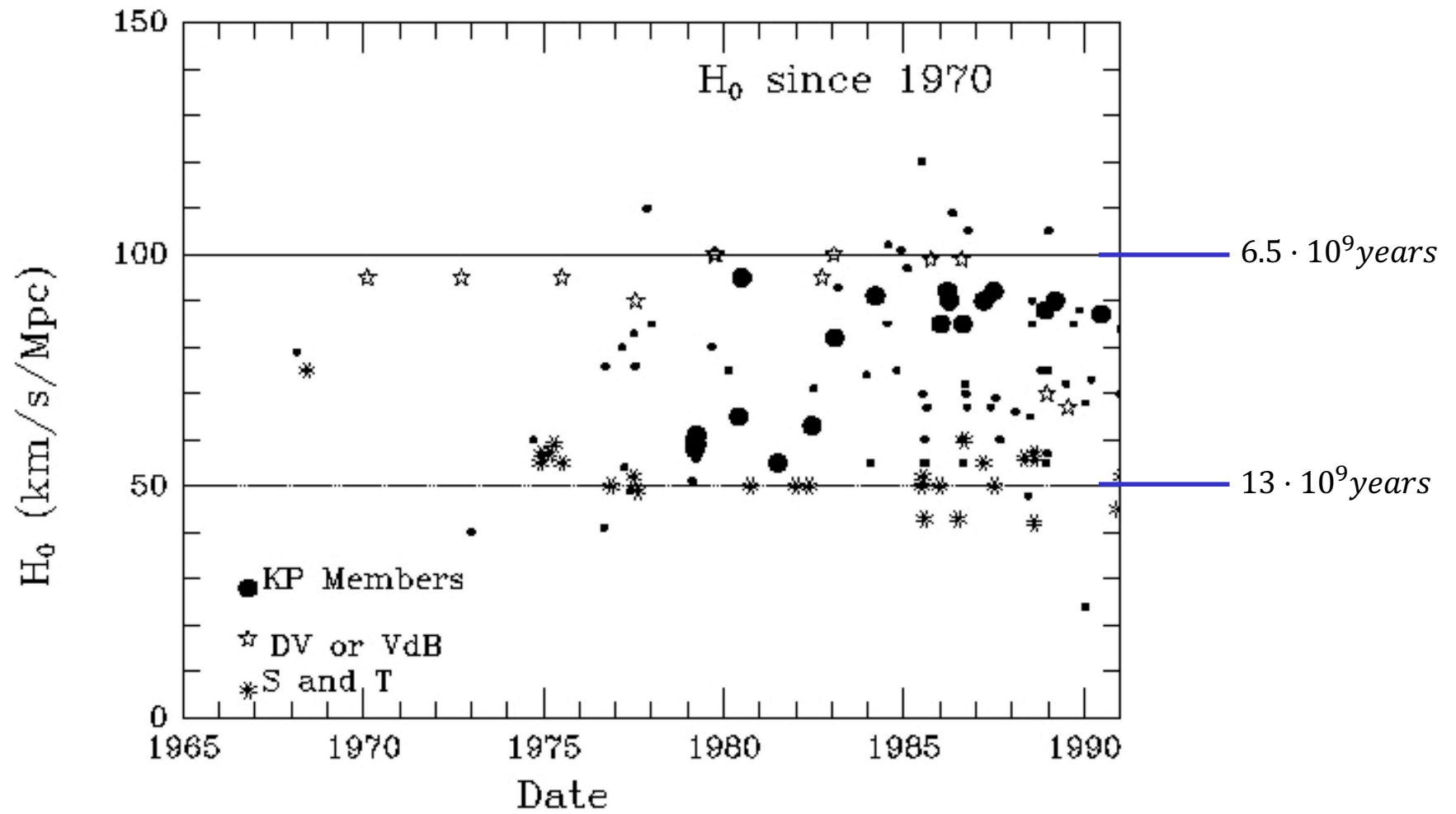


J. Huchra



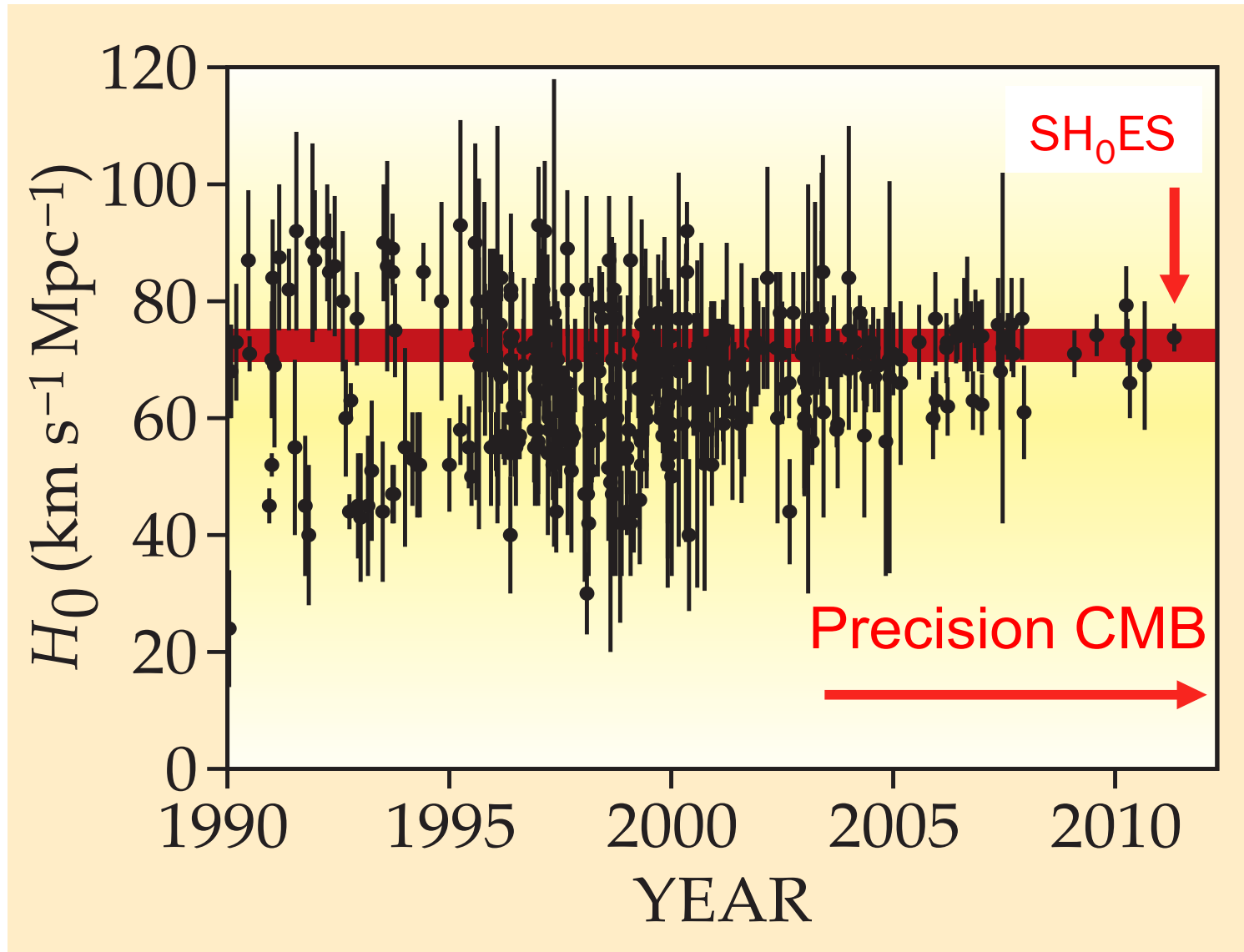
<https://www.cfa.harvard.edu/~dfabricant/huchra/hubble/>

# History of $H_0$



Copyright J. Huchra 2008

# History of $H_0$



A. Riess

# Expansion of the Universe

Luminosity distance in an isotropic, homogeneous universe as a Taylor expansion

$$D_L = \frac{cz}{H_0} \left\{ 1 + \frac{1}{2}(1 - q_0)z - \frac{1}{6} \left[ 1 - q_0 - 3q_0^2 + j_0 \pm \frac{c^2}{H_0^2 R^2} \right] z^2 + O(z^3) \right\}$$

Hubble-  
Lemaître Law

deceleration

jerk/equation of state

$$H_0 = \frac{\dot{a}}{a} \quad q_0 = -\frac{\ddot{a}}{a} H_0^{-2} \quad j_0 = \frac{\dddot{a}}{a} H_0^{-3}$$

THE ABILITY OF THE 200-INCH TELESCOPE TO DISCRIMINATE  
BETWEEN SELECTED WORLD MODELS

ALLAN SANDAGE

Mount Wilson and Palomar Observatories  
Carnegie Institution of Washington, California Institute of Technology

Received October 14, 1960; revised November 5, 1960

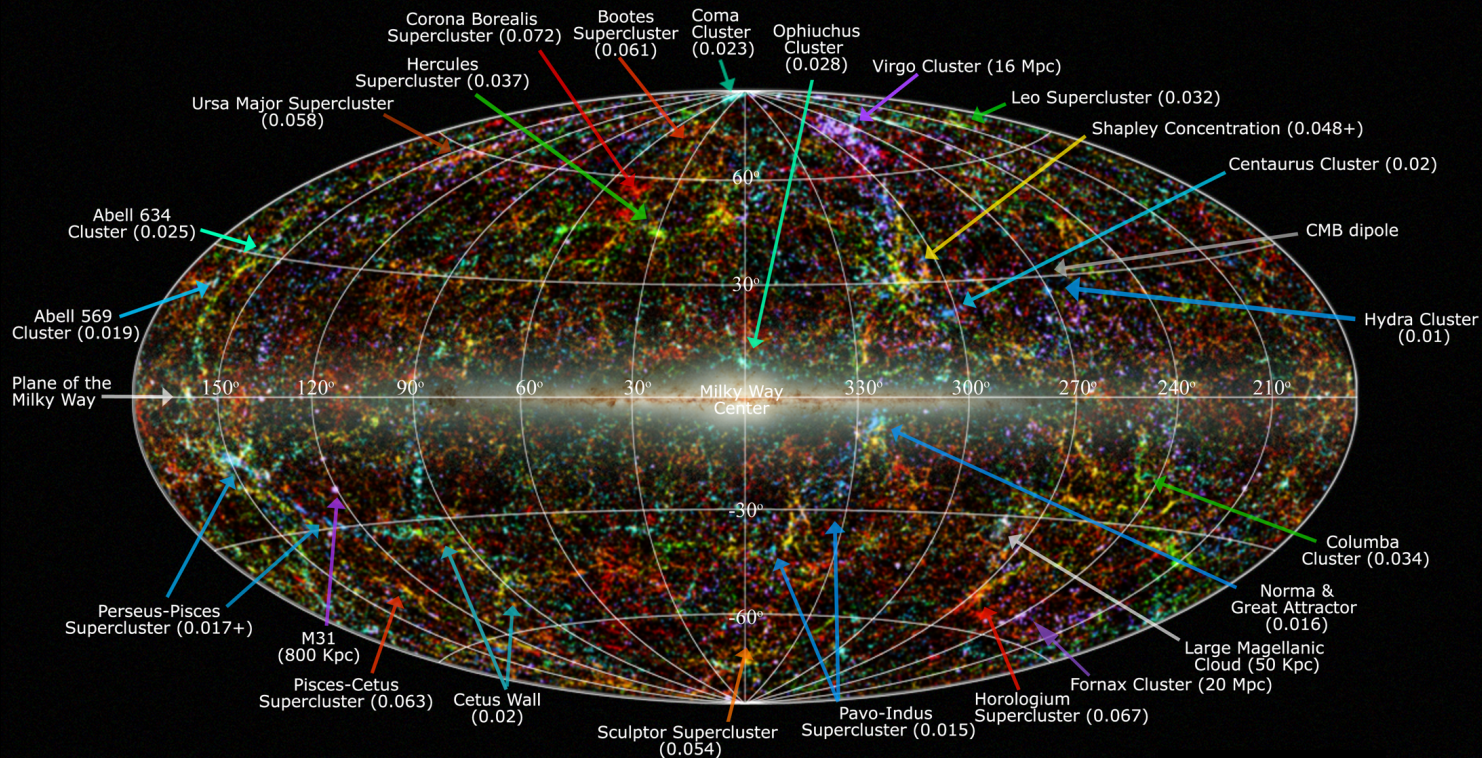
$$m_{\text{bol}} = 5 \log z + 1.086 (1 - q_0) z + \dots + \text{const.},$$



(15) ApJ 1961

# Extragalactic Distances Required for a 3D picture of the (local) universe

## 2MASS Redshift Survey



**Legend:** image shows 2MASS galaxies color coded by the 2MRS redshift (Huchra et al 2011); familiar galaxy clusters/superclusters are labeled (numbers in parenthesis represent redshift).

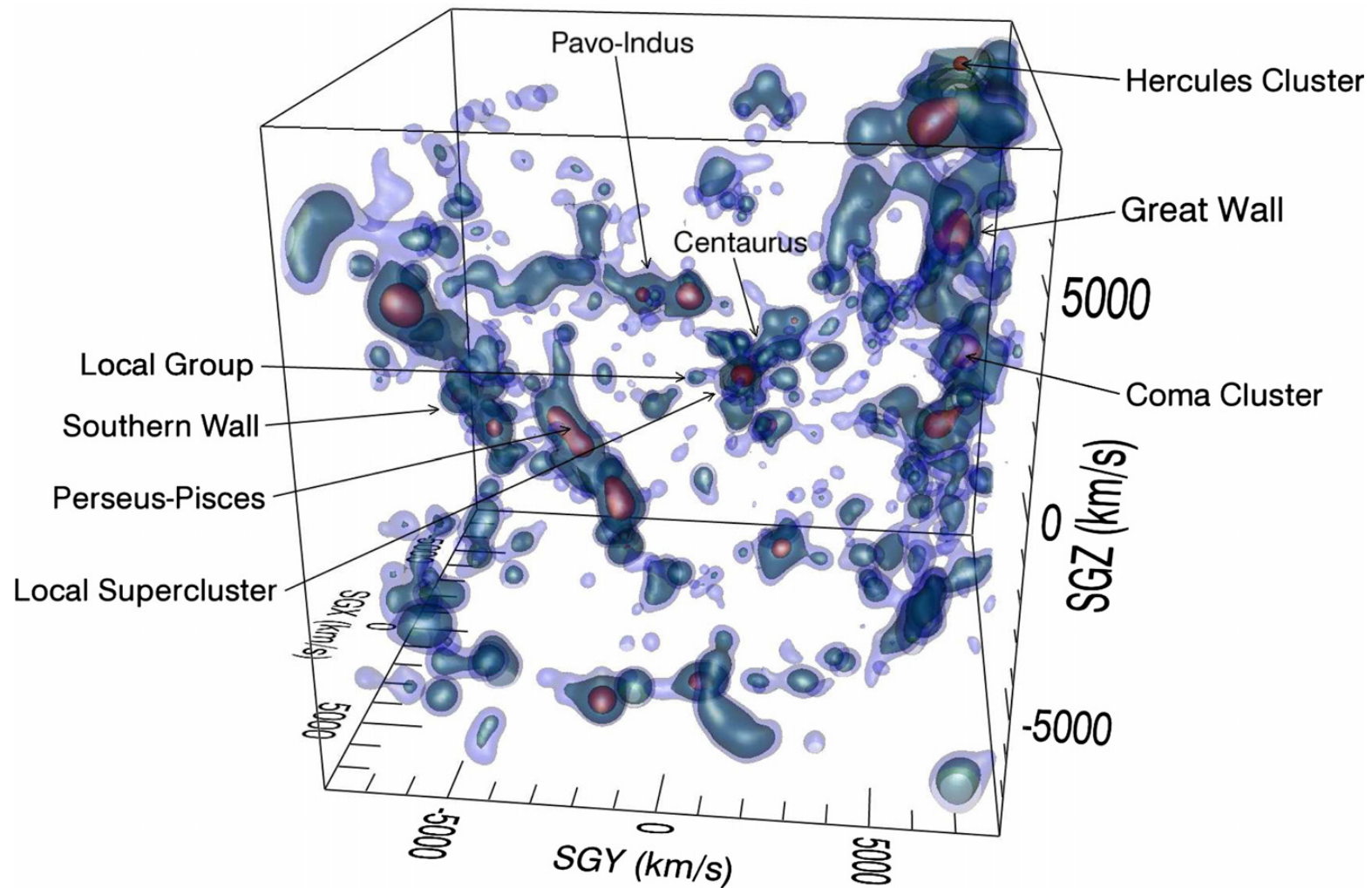
Graphic created by T. Jarrett (IPAC/Caltech)



# Extragalactic Distances

THE ASTRONOMICAL JOURNAL, 146:69 (14pp), 2013 September

COURTOIS ET AL.

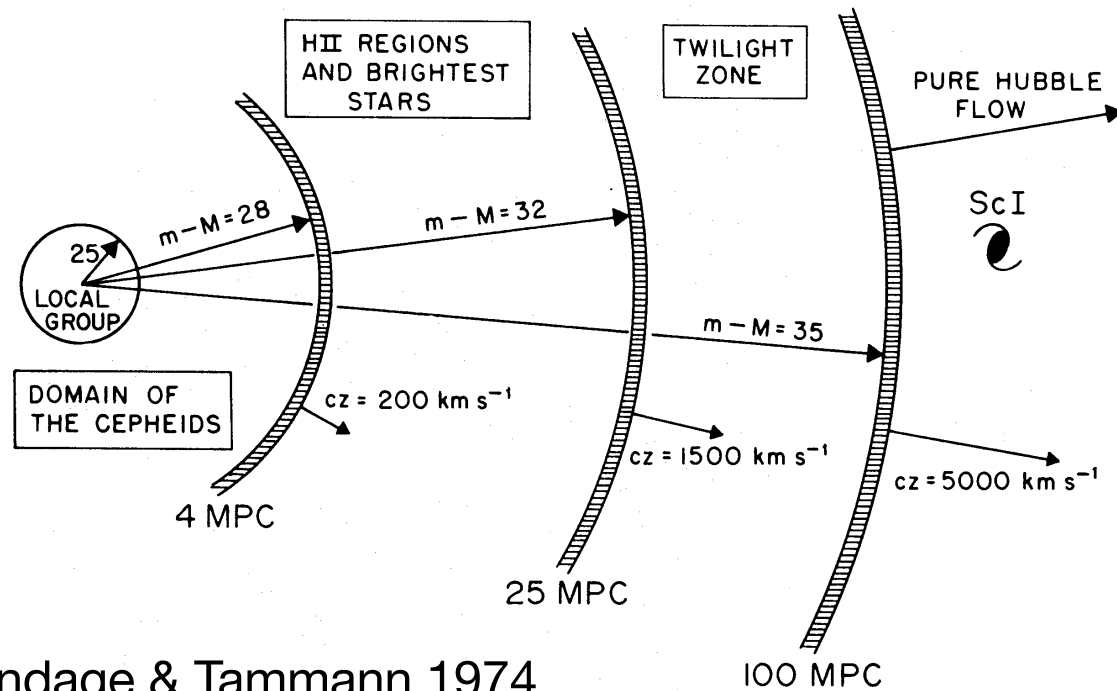


**Figure 8.** Perspective view of the V8k catalog after correction for incompleteness and represented by three layers of isodensity contours. The region in the vicinity of the Virgo Cluster now appears considerably diminished in importance. The dominant structures are the Great Wall and the Perseus–Pisces chain, with the Pavo–Indus feature of significance.

# Measuring $H_0$

## Classical approach

→ distance ladder to reach (smooth) Hubble flow



Sandage & Tamman 1974

# Hubble Constant

## Three different methods

### 1. Distance ladder

- Calibrate next distance indicator with the previous

### 2. Physical methods

- Determine either luminosity or length through physical quantities
  - Sunyaev-Zeldovich effect (galaxy clusters)
  - Expanding photosphere method in supernovae
  - Physical calibration of thermonuclear supernovae
  - Geometric methods, e.g. masers

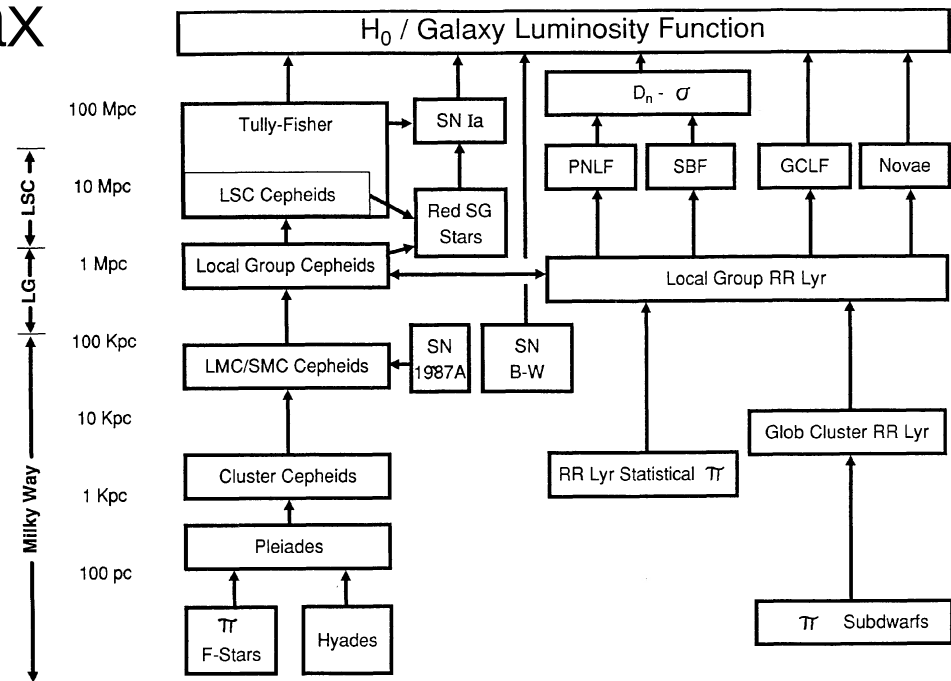
### 3. Global solutions

- Use knowledge of all cosmological parameters
  - Cosmic Microwave Background

# Classical Distance Ladder

Primary distance indicators (within the Milky Way)

- trigonometric parallax
- proper motion
- apparent luminosity
  - main sequence
  - red clump stars
  - RR Lyrae stars
  - eclipsing binaries
  - Cepheid stars



Pathways to Extragalactic Distances

Jacoby et al. 1992

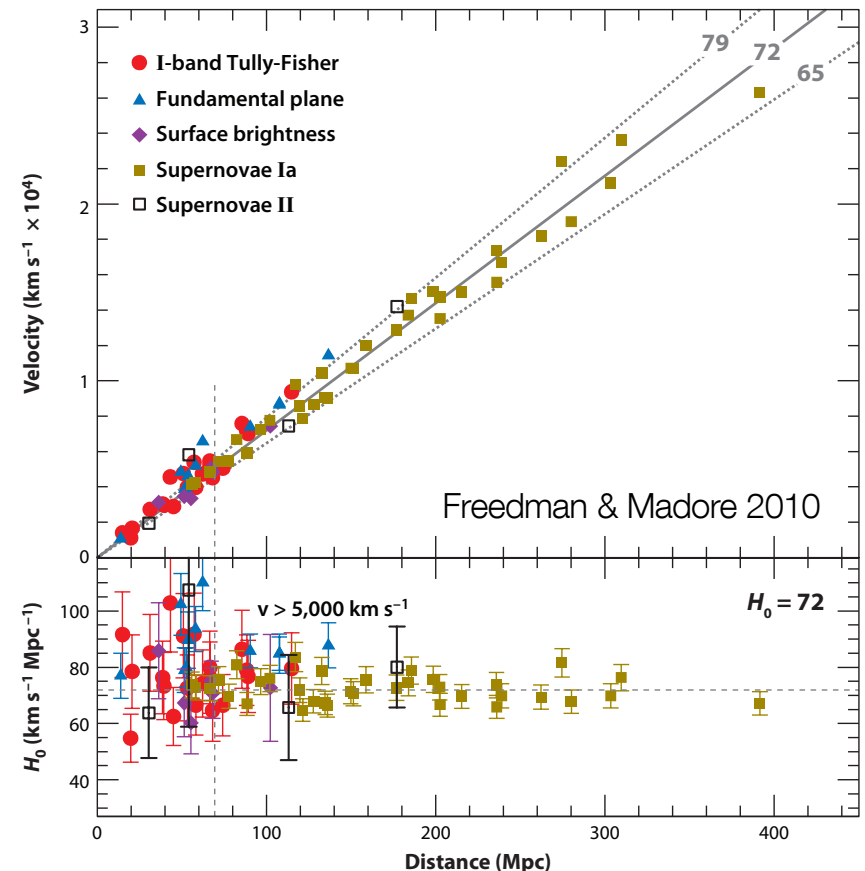
# Classical Distance Ladder

Secondary distance indicators (beyond the Local Group)

- Important check
  - Large Magellanic Cloud
- Tully-Fisher relation
- Fundamental Plane
- Supernovae (mostly SN Ia)
- Surface Brightness Fluctuations



Gruber Cosmology Prize



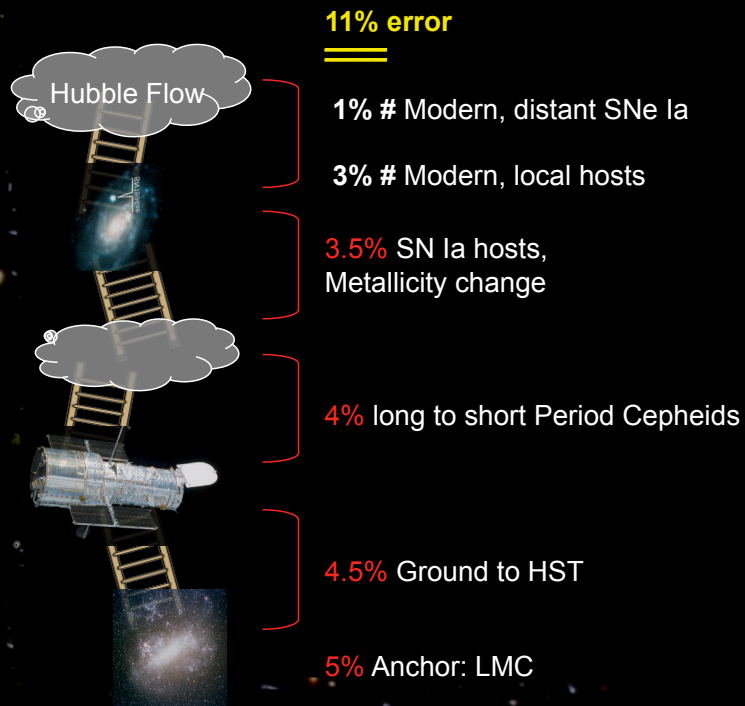


# Hubble Constant

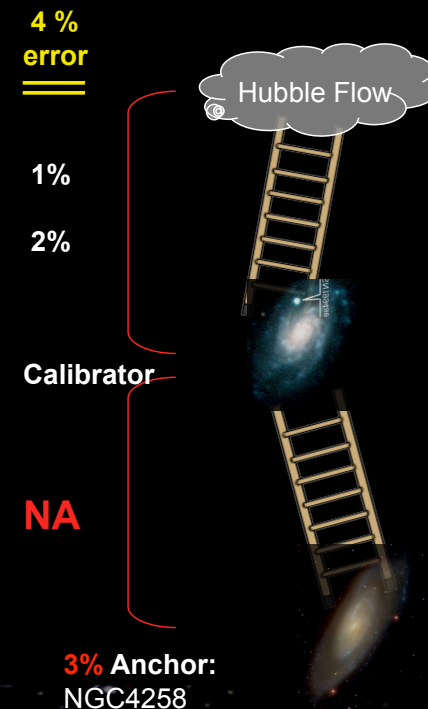
## Calibration of $M(SN Ia @ max)$

### Distance ladder

#### PAST DISTANCE LADDER (100 Mpc)



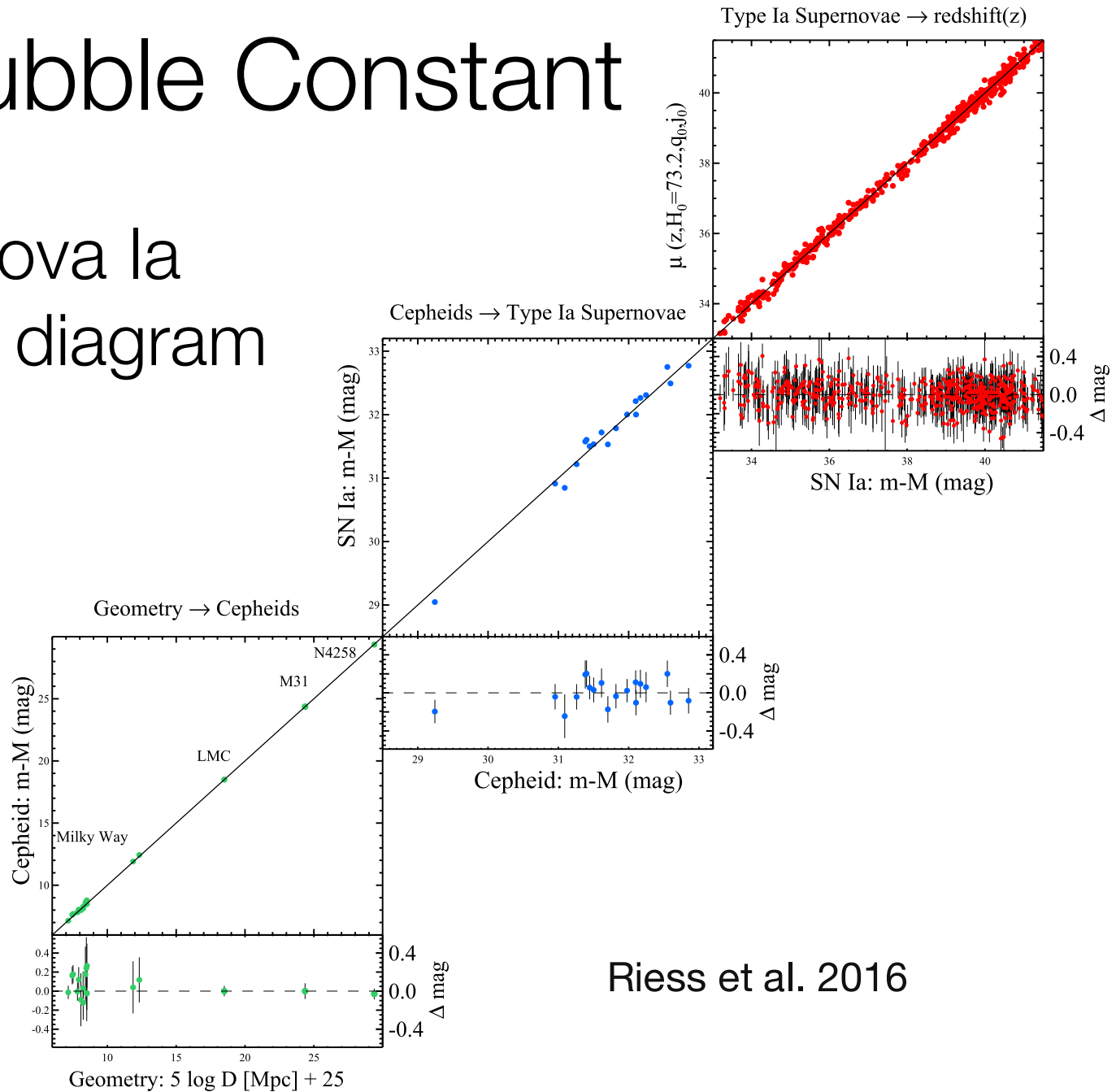
#### NEW LADDER (100 Mpc)



Adam Riess

# Hubble Constant

## Supernova Ia Hubble diagram



# $H_0$ with Supernovae

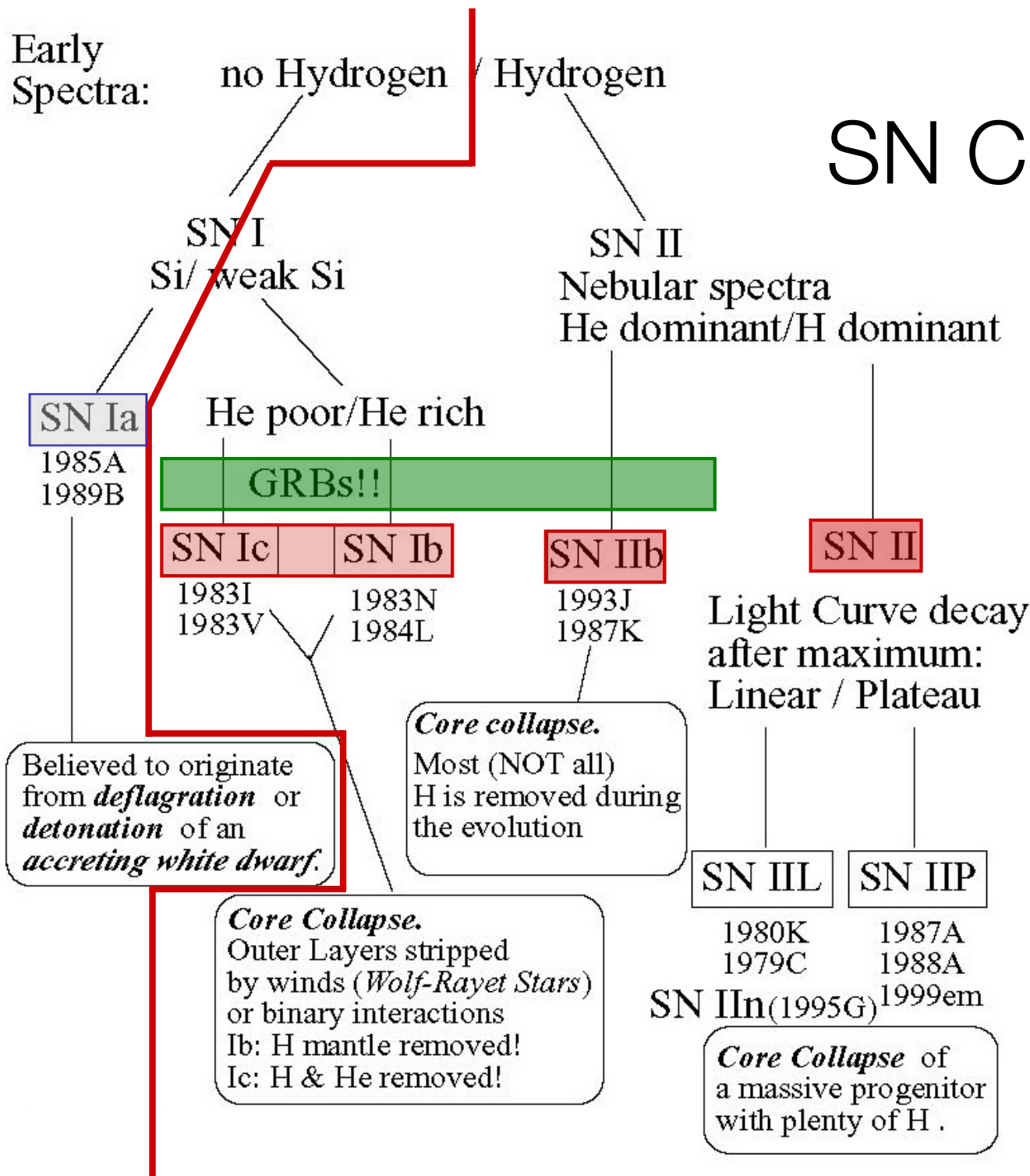
- Local calibrators (calibrate the Cepheid L-P rel.)
  - Large Magellanic Cloud
    - 1% accuracy with eclipsing binaries (Pietrzyński et al. (2019))
  - Maser in NGC 4258
    - 3% accuracy (Humphreys et al. 2013)
  - geometric distances (parallaxes) to nearby Cepheids
- Extinction
  - absorption in the Milky Way and the host galaxy
  - corrections not always certain
- Peculiar velocities of galaxies
  - typically around 300 km/s



# Gaia and $H_0$

- Calibrate Cepheid distances with parallaxes
  - long-period Cepheids so far not accessible
- Single step to the SNe Ia
  - reduced uncertainty on  $H_0$
- Currently problems with systematic offsets between HST and Gaia (Riess et al. 2018)
  - Discrepancy of  $(-46 \pm 13)\mu as$
- Goal: uncertainty less than 1%

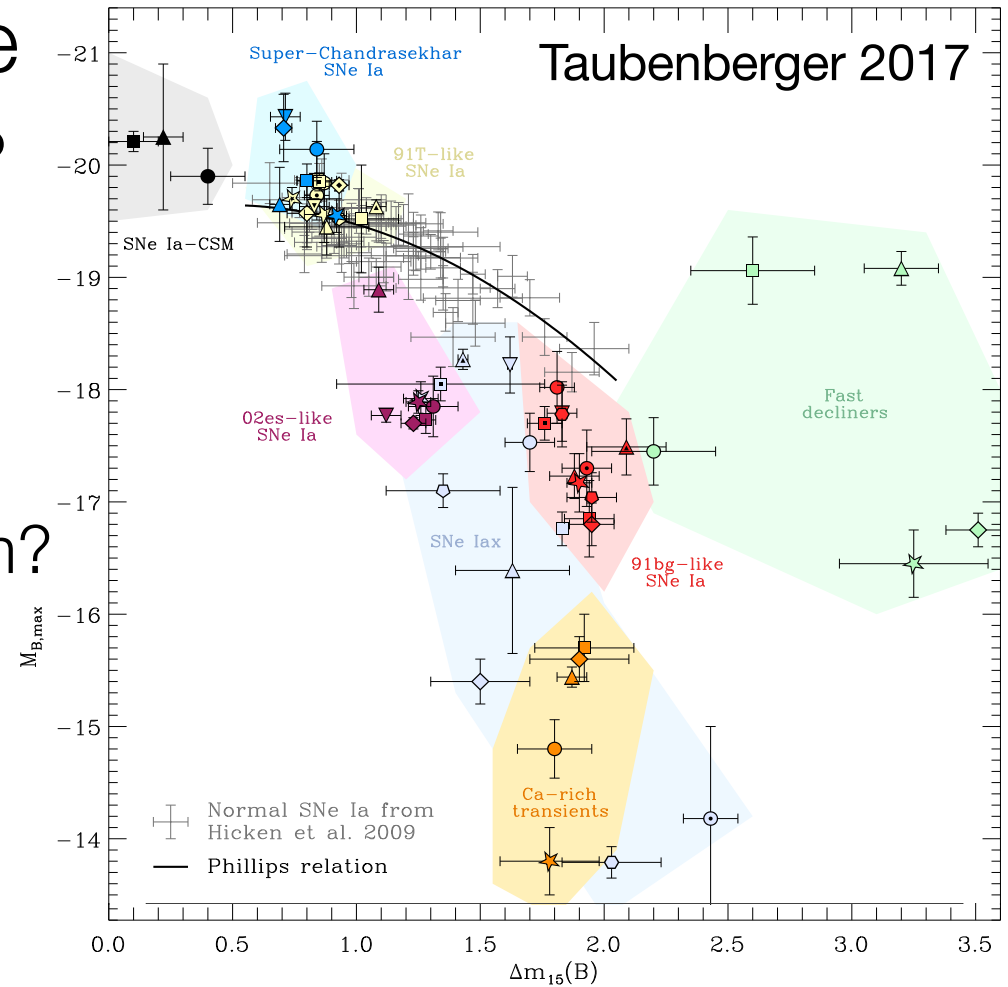
# SN Classification



# Type Ia Supernovae

## Variations on a theme

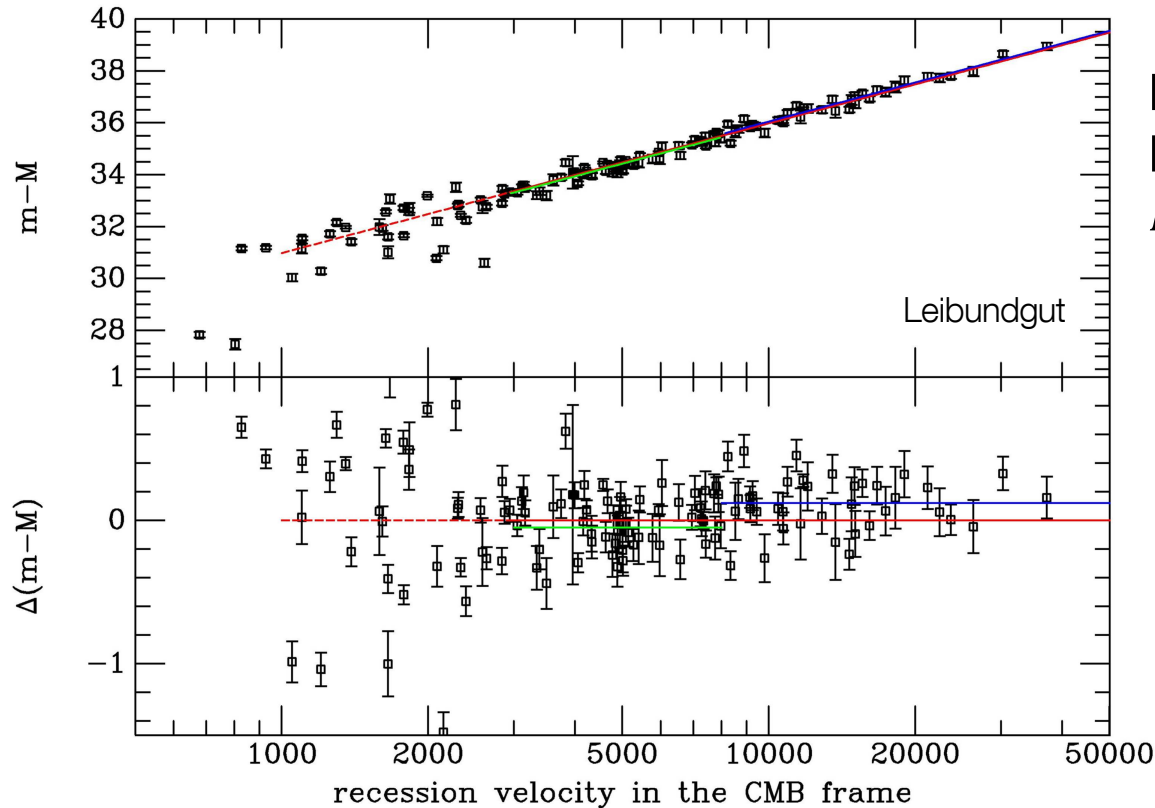
- critical parameters?
  - nickel mass
  - ejecta mass
  - explosion energy(?)
  - explosion mechanism?
  - progenitor evolution?



# Hubble Constant

## SN Hubble diagram

$$m - M = 5 \log v + 25 - 5 \log H_0$$

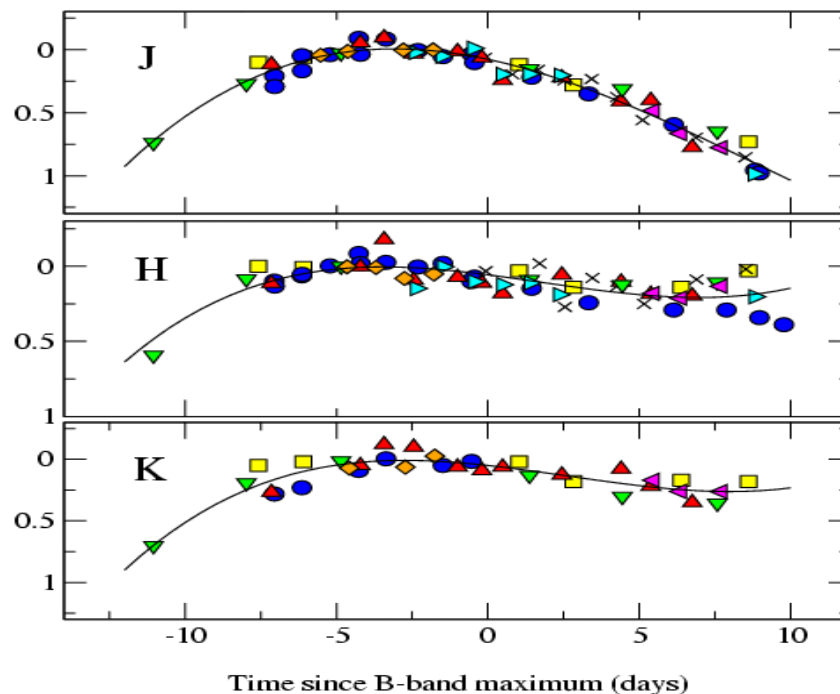


Proves  $M$  is constant  
Direct connection of  $M$  and  
 $H_0$

Leibundgut

# The Promise of the (Near-)Infrared

- Extinction is much reduced in the near-IR
  - $A_H/A_V \cong 0.19$  (Cardelli et al. 1989)
- SNe Ia much better behaved



	SN	$m_{15}(B)$
▲	1980N	(1.29)
■	1986G	(1.79)
▲	1998bu	(1.05)
✕	1999aw	(0.81)
●	1999ee	(0.94)
▼	2000ca	(1.01)
◆	2001el	(1.15)

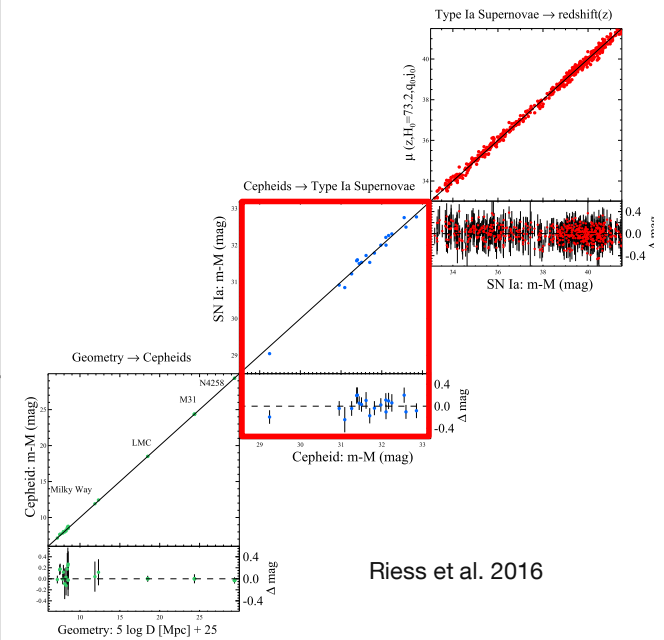
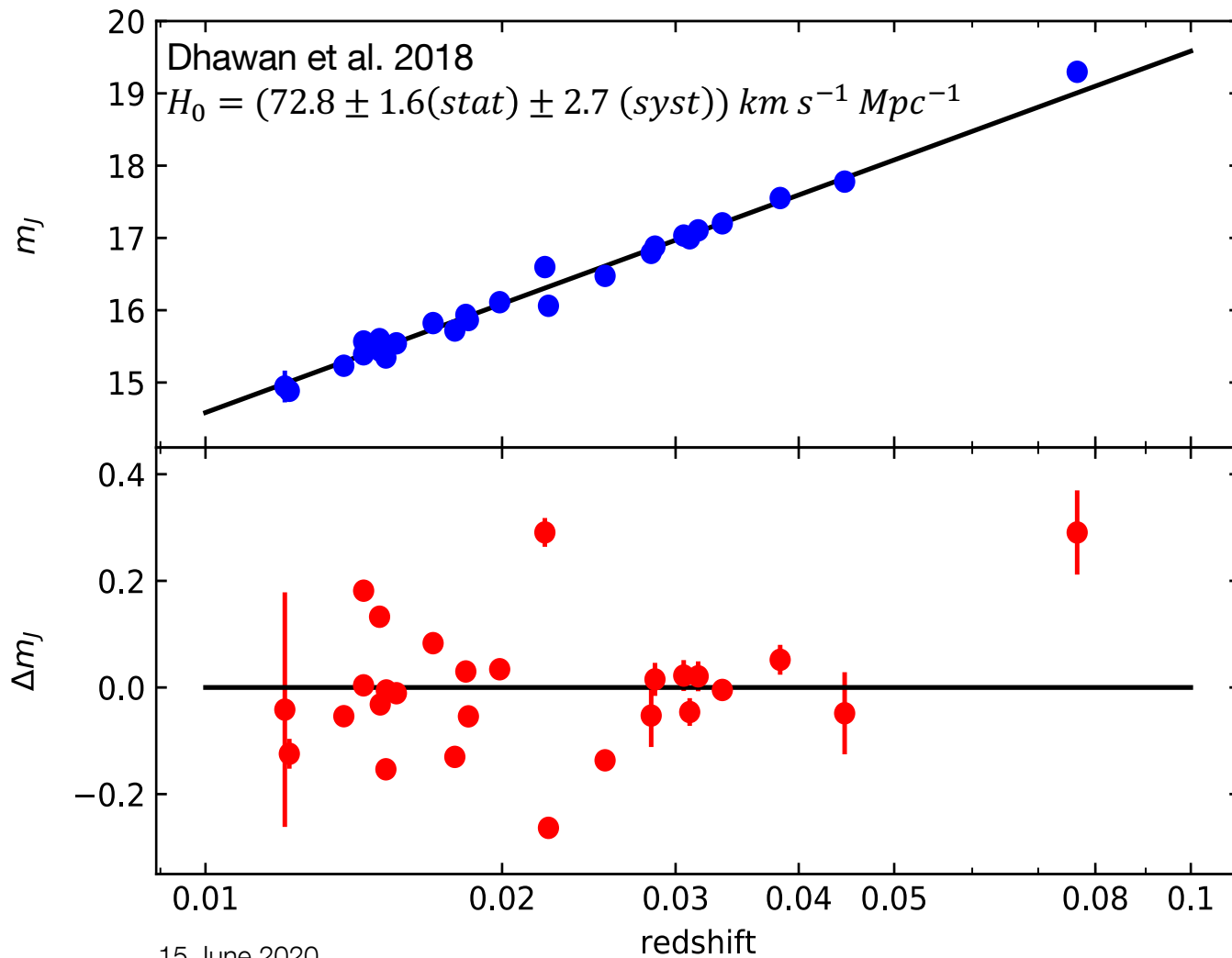
Krisciunas et al. (2004)

Mark Phillips

Bruno Leibundgut

# Current Status (NIR)

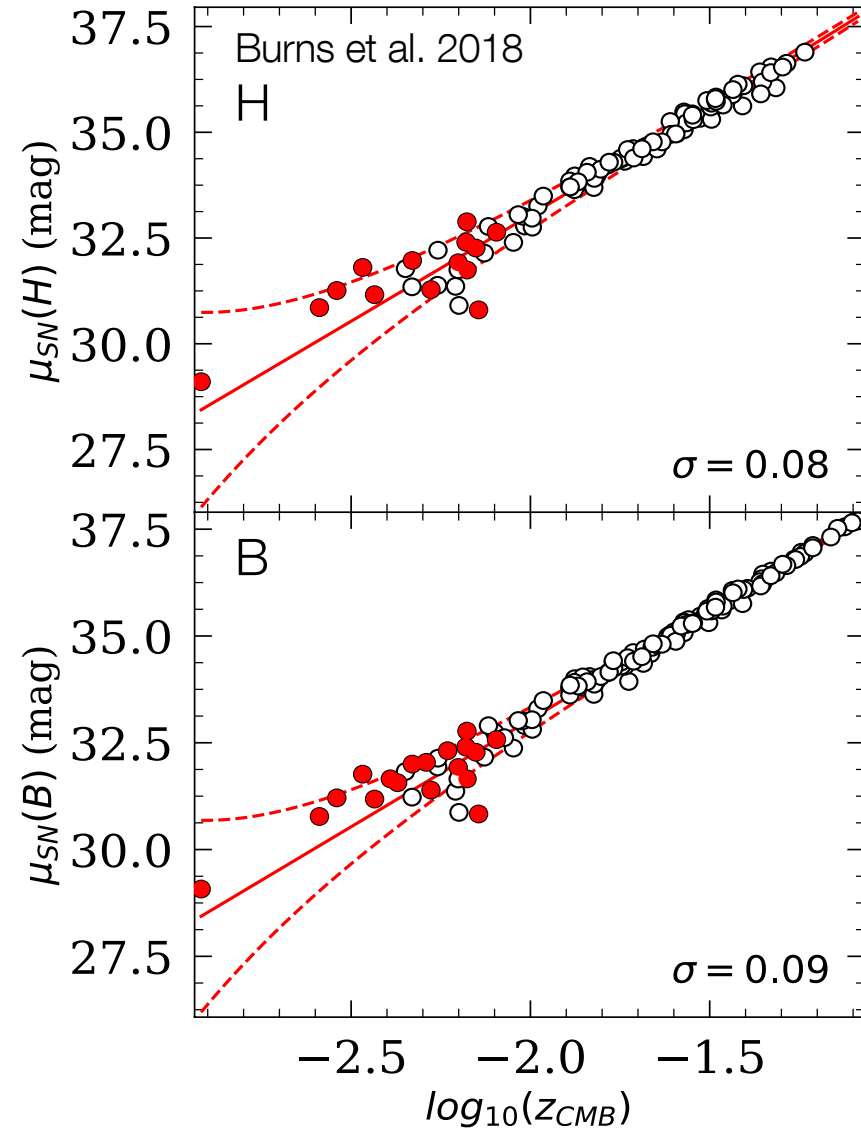
## 9 calibrators + 27 Hubble flow SNe



Riess et al. 2016

# Infrared SN Hubble Diagrams

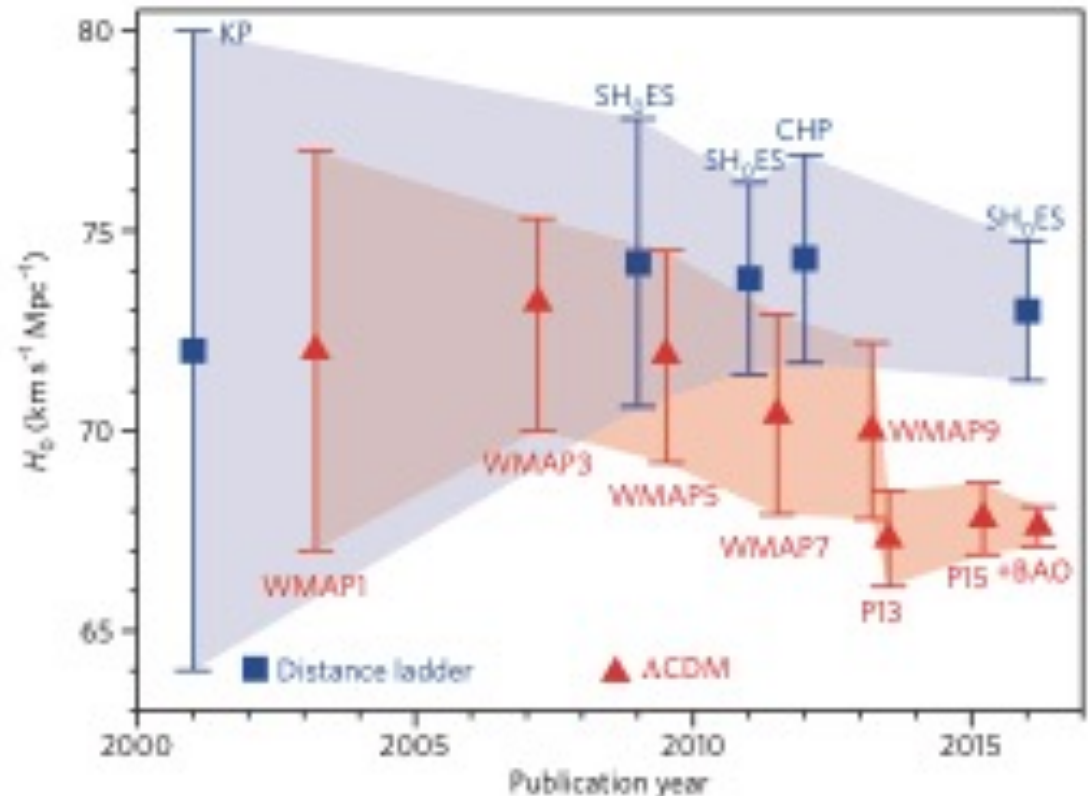
NIR Hubble diagram  
becoming  
competitive



# Problem solved?

New discrepancy between the measurements of the local  $H_0$  (distance ladder) and early universe (CMB)

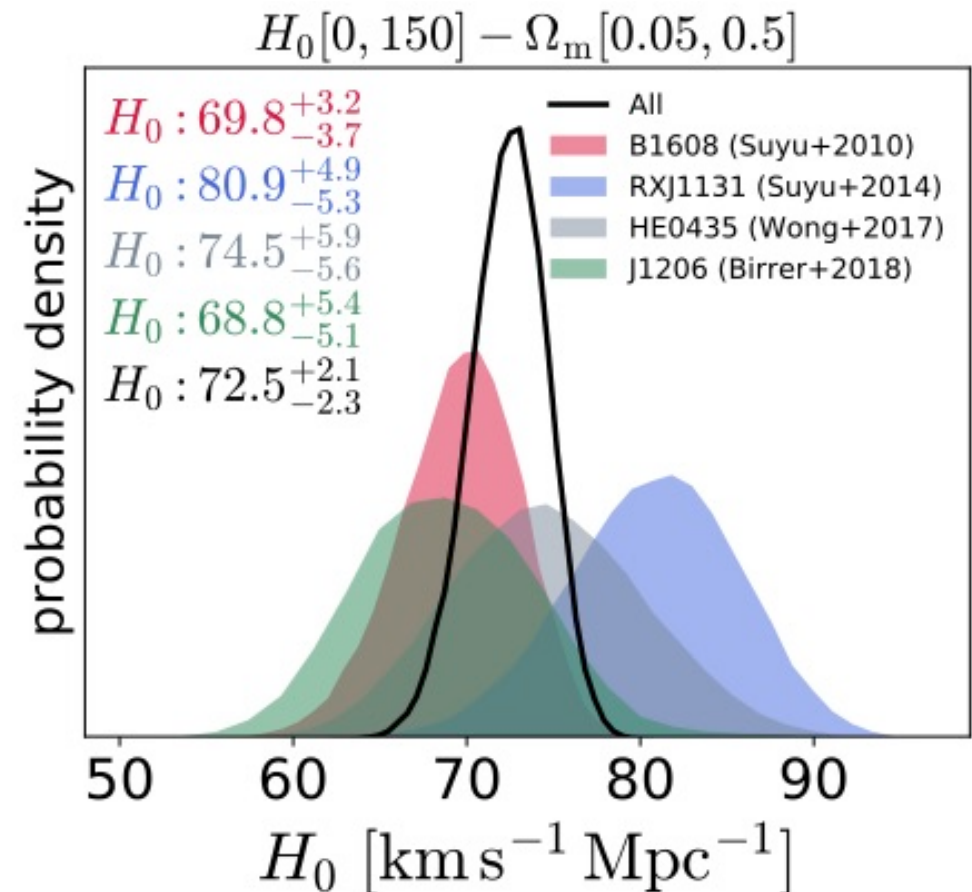
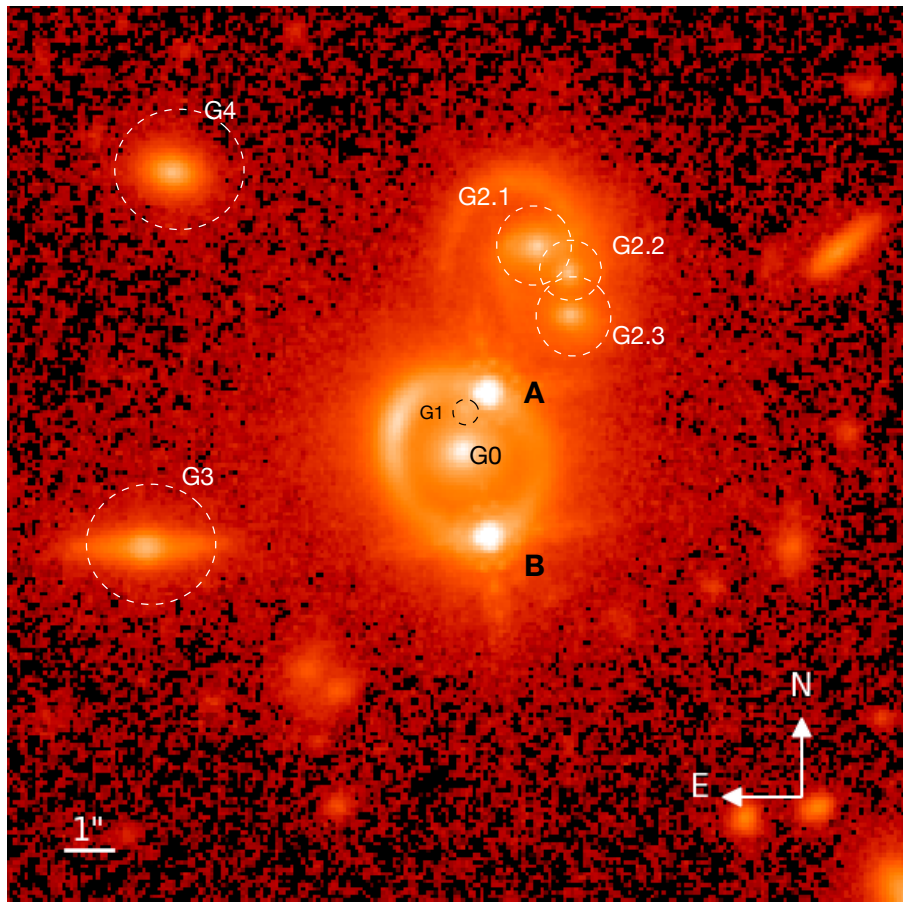
Indication of an incomplete model of cosmology?





# Promising Results from strong lensing

## Time delays in lensed quasars



Birrer et al. 2018

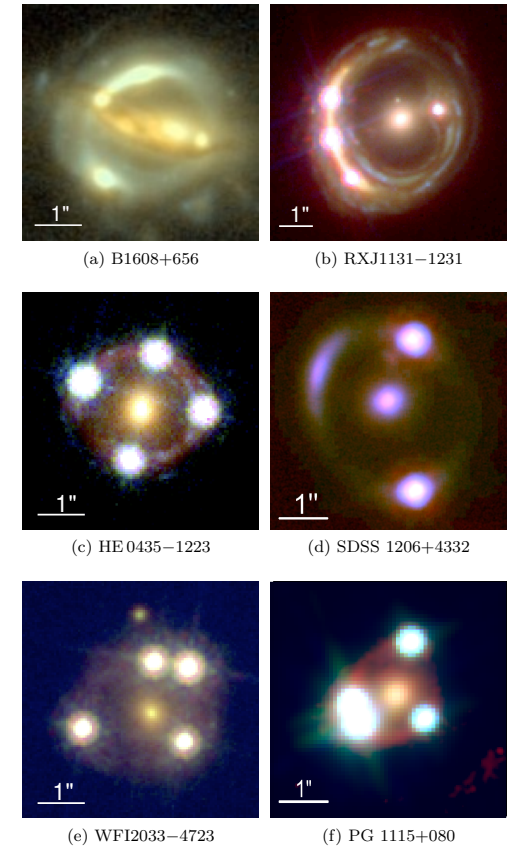
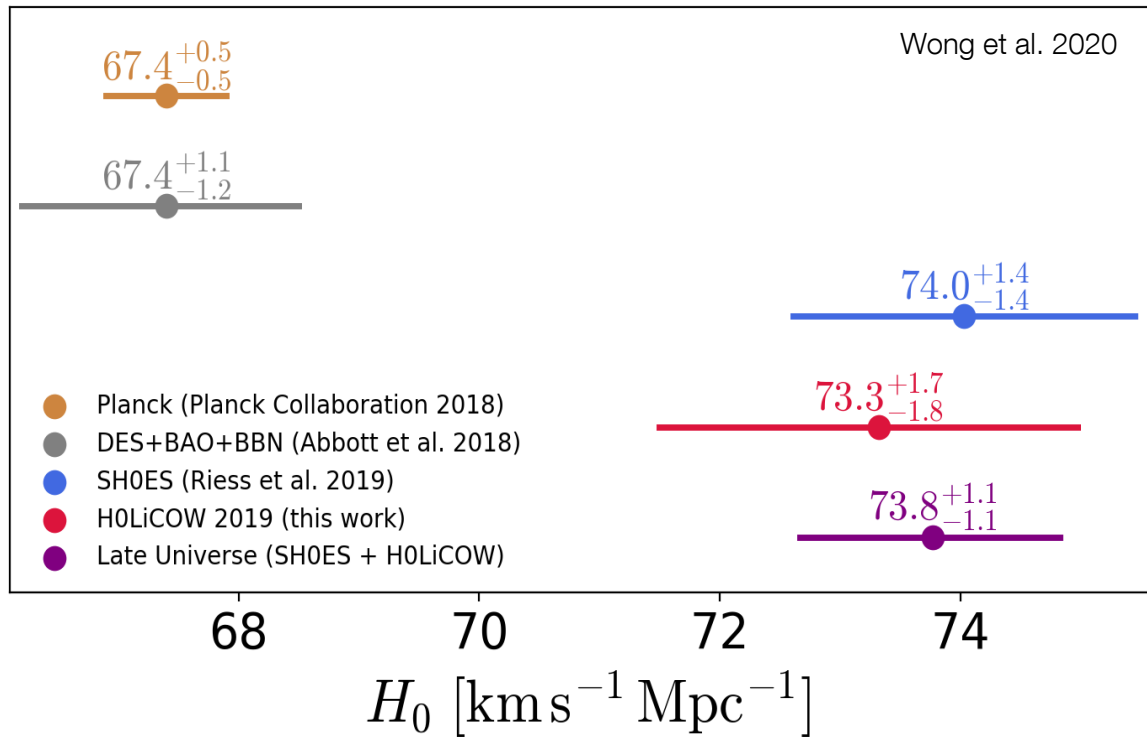
Bruno Leibundgut



# Gravitationslinsen

## H0LICOW collaboration

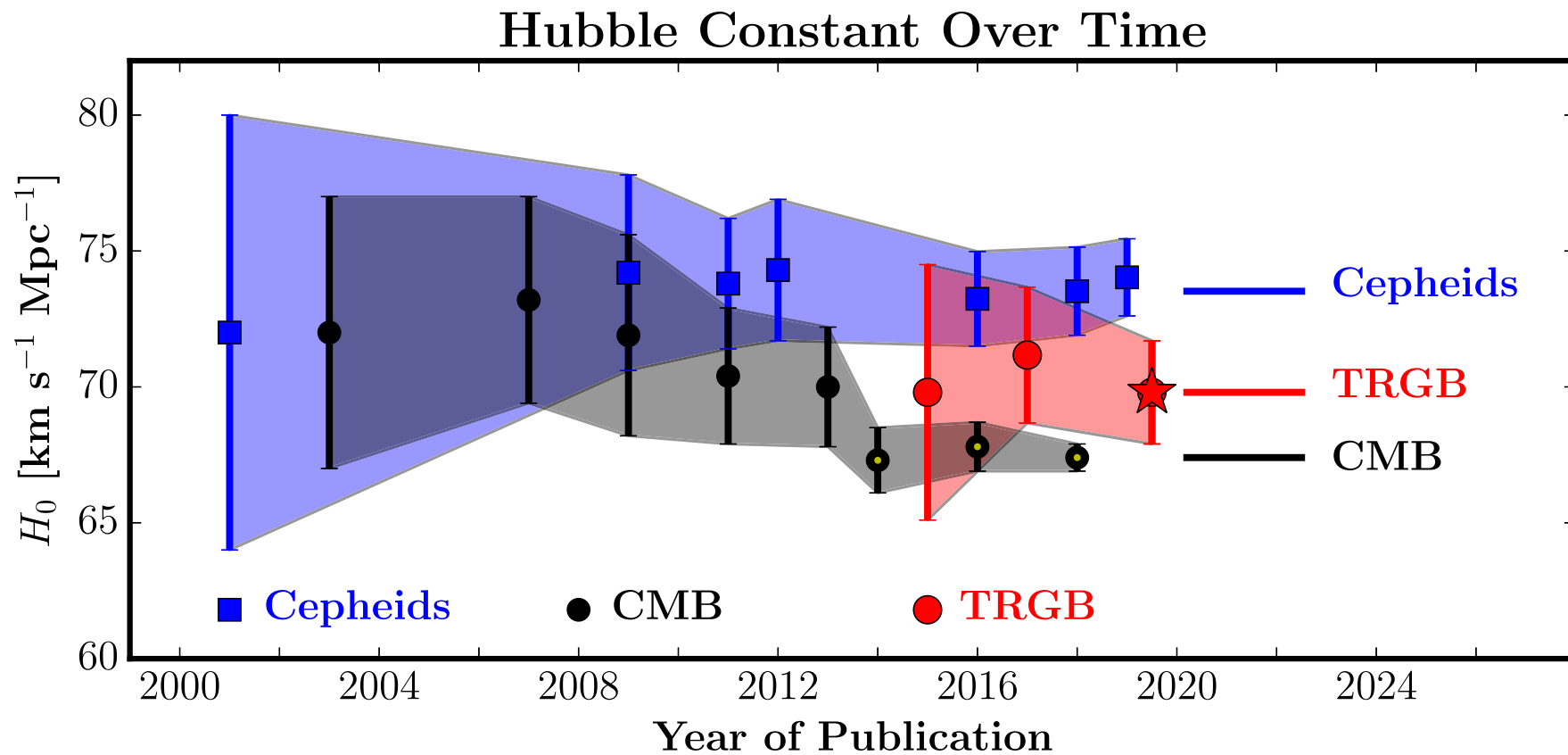
flat  $\Lambda$ CDM



# ...and another attempt



## Tip of the Red Giant Branch



# Hubble Constant(s)

Planck satellite (CMB; 2018)

measurement at  $z \approx 1000$

$$H_0 = (67.4 \pm 0.5) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Riess et al. (local; 2016)

$$H_0 = (73.24 \pm 1.74) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Dhawan et al. (local; NIR; 2018)

same zero-point as Riess et al. (2016)

$$H_0 = (72.8 \pm 1.6(\text{stat}) \pm 2.7(\text{syst})) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Riess et al. (local; 2018)

$$H_0 = (73.53 \pm 1.62) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Riess et al. (local; 2019)

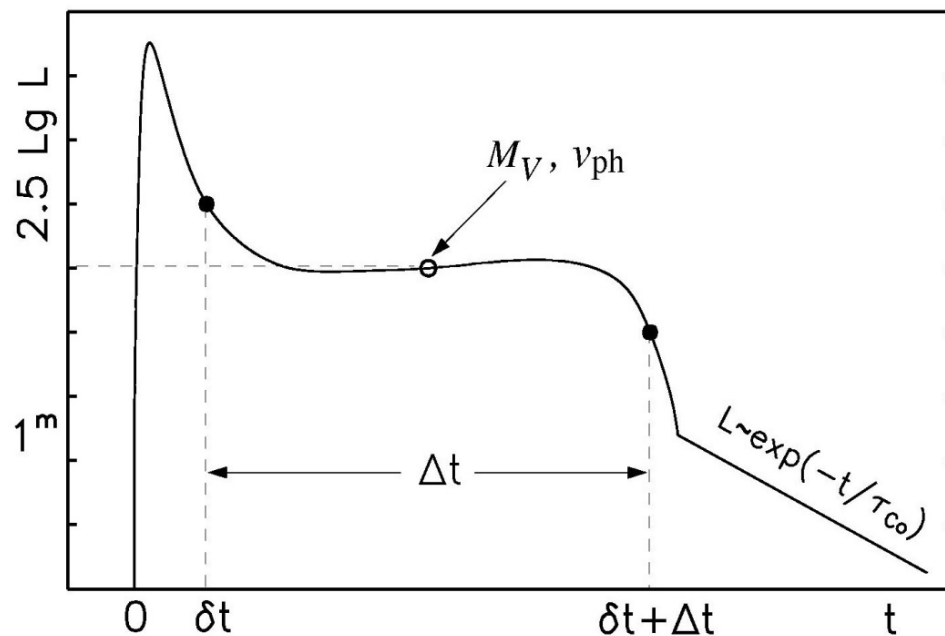
$$H_0 = (74.22 \pm 1.82) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

Freedman et al. (local, TRGB; 2019)

$$H_0 = (69.8 \pm 0.8 \pm 1.7) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

# Physical parameters of core collapse SNe

Light curve shape and the velocity evolution can give an indication of the total explosion energy, the mass and the initial radius of the explosion



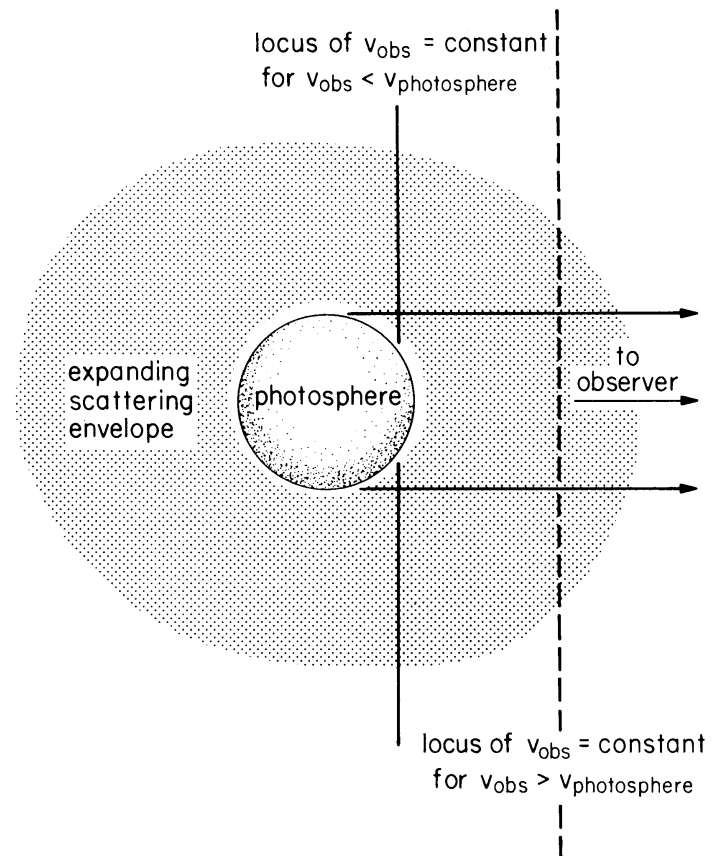
Observables:

- length of plateau phase  $\Delta t$
- luminosity of the plateau  $L_V$
- velocity of the ejecta  $v_{ph}$
  
- $E \propto \Delta t^4 \cdot v_{ph}^5 \cdot L^{-1}$
- $M \propto \Delta t^4 \cdot v_{ph}^3 \cdot L^{-1}$
- $R \propto \Delta t^2 \cdot v_{ph}^{-4} \cdot L^2$

# Expanding Photosphere Method

Modification of Baade-Wesselink method  
for variable stars

- Assumes
  - Sharp photosphere  
→ thermal equilibrium
  - Spherical symmetry  
→ radial velocity
  - Free expansion



Kirshner & Kwan 1974

# Expanding Photosphere Method

$$\theta = \frac{R}{D} = \sqrt{\frac{f_\lambda}{\zeta_\lambda^2 \pi B_\Lambda(T)}}; R = v(t - t_0) + R_0; D_A = \frac{v}{\theta} (t - t_0)$$

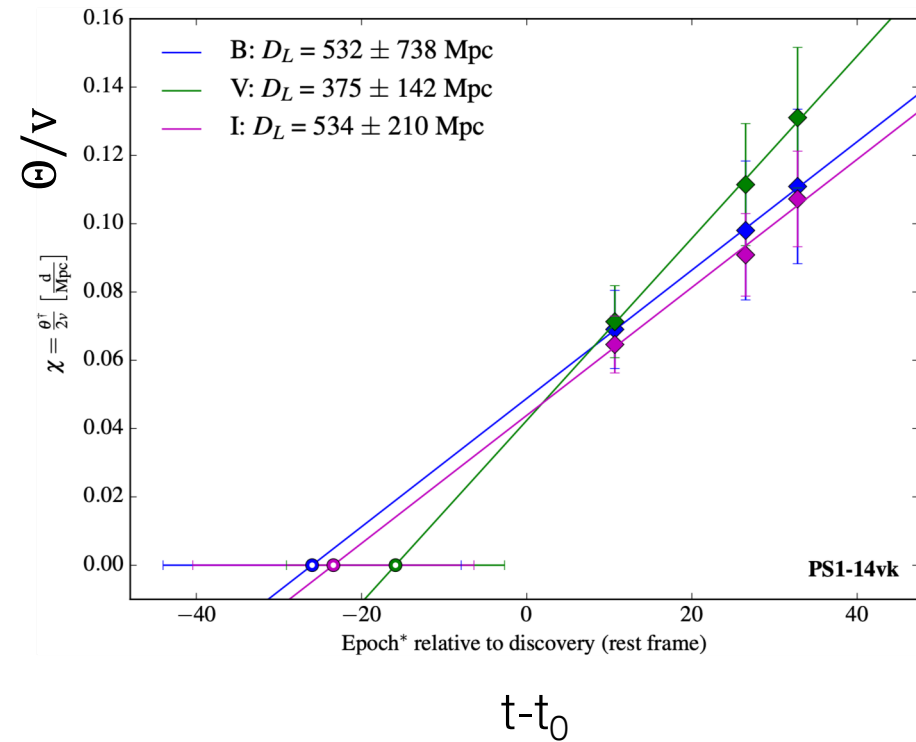
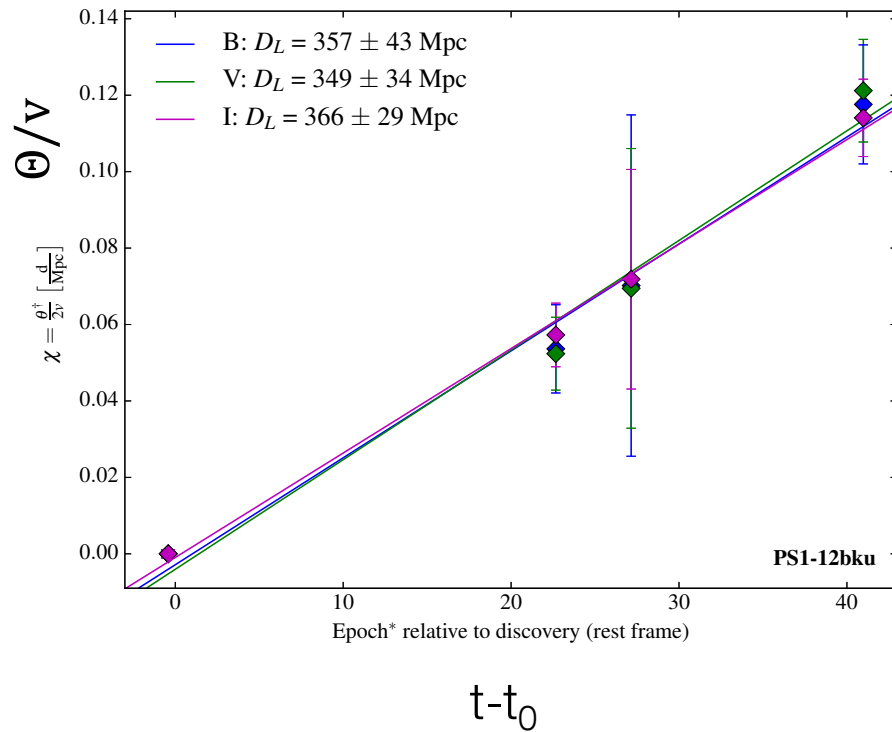
- $R$  from radial velocity
  - Requires lines formed close to the photosphere
- $D$  from the surface brightness of the black body
  - Deviation from black body due to line opacities
  - Encompassed in the dilution factor  $\zeta^2$

# Expanding Photosphere Method

- It's all in the data...

$$\frac{\Theta}{v} = \frac{1}{D_A} (t - t_0)$$

Gall et al. 2018



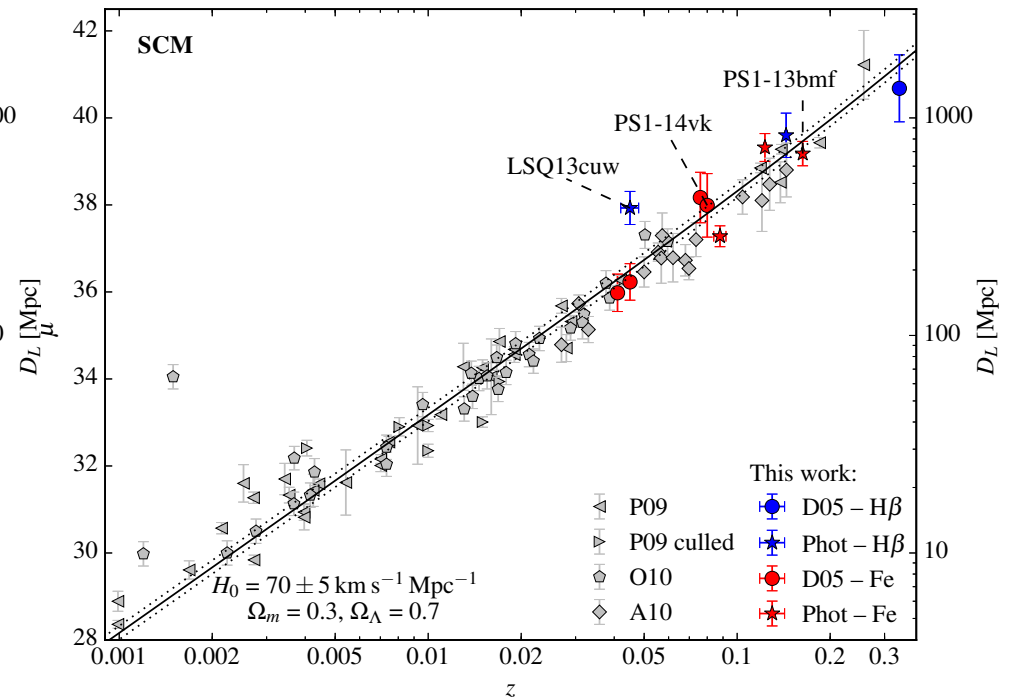
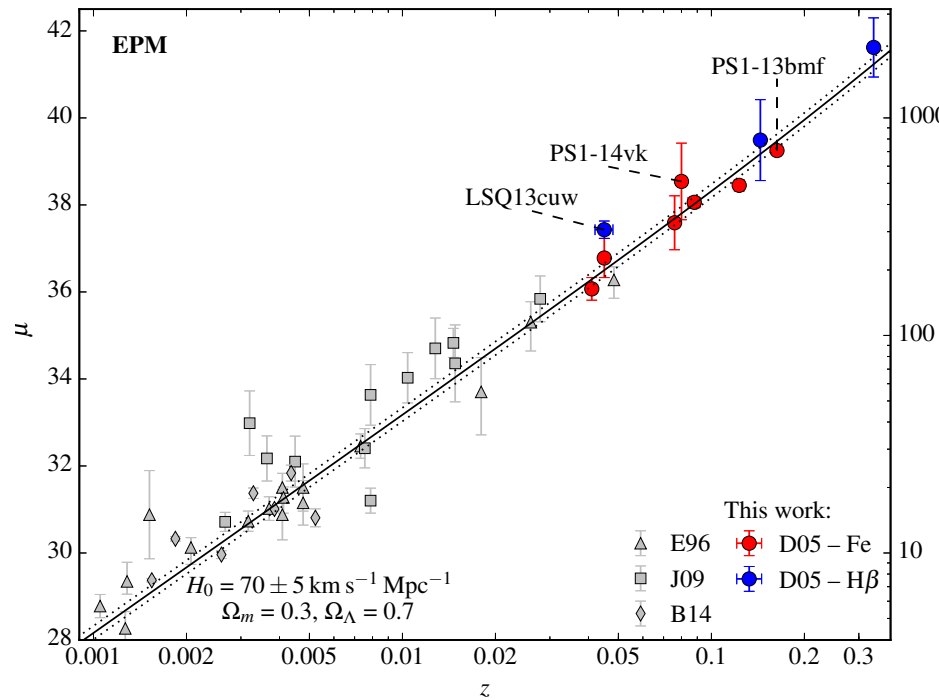


# Preliminary Results

## Consistent results

– not independent as local calibration required

Gall et al. 2018

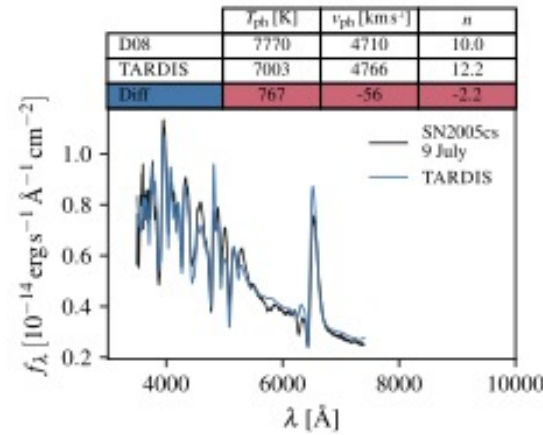


# Expanded Photosphere Method Reloaded

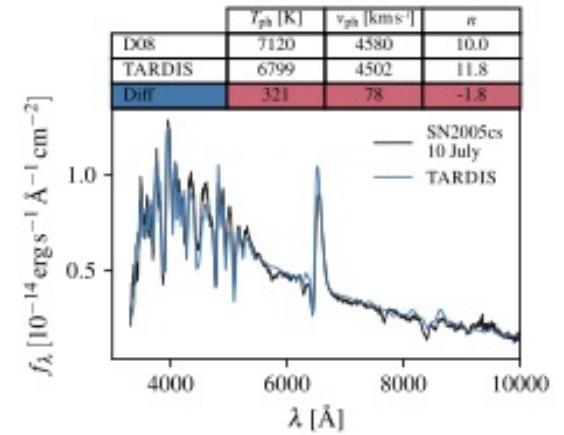
- Use individual atmospheric models for the spectral fits
  - use of the TARDIS radiation transport model
  - absolute flux emitted
- Accurate explosion date
  - accurate zero point
- At least 5 epochs per supernova

# Atmosphere Models

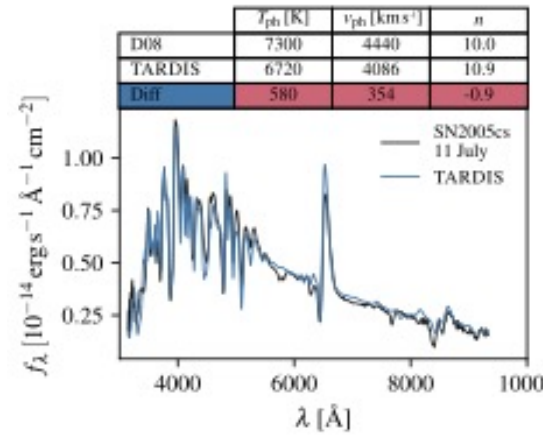
TARDIS fits for different epochs



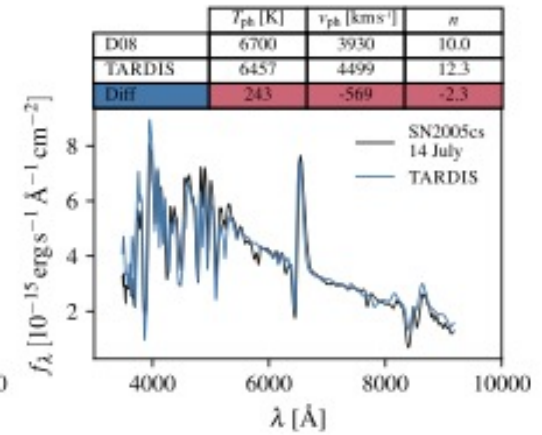
(a) 9 July 2005



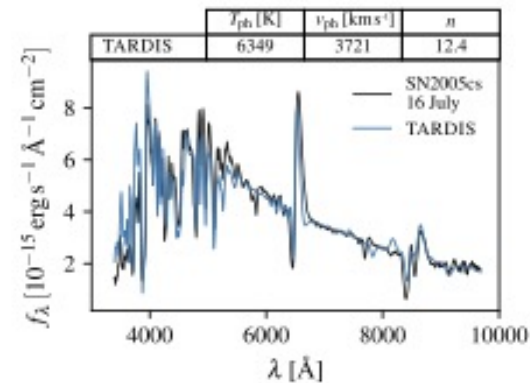
(b) 10 July 2005



(c) 11 July 2005



(d) 14 July 2005

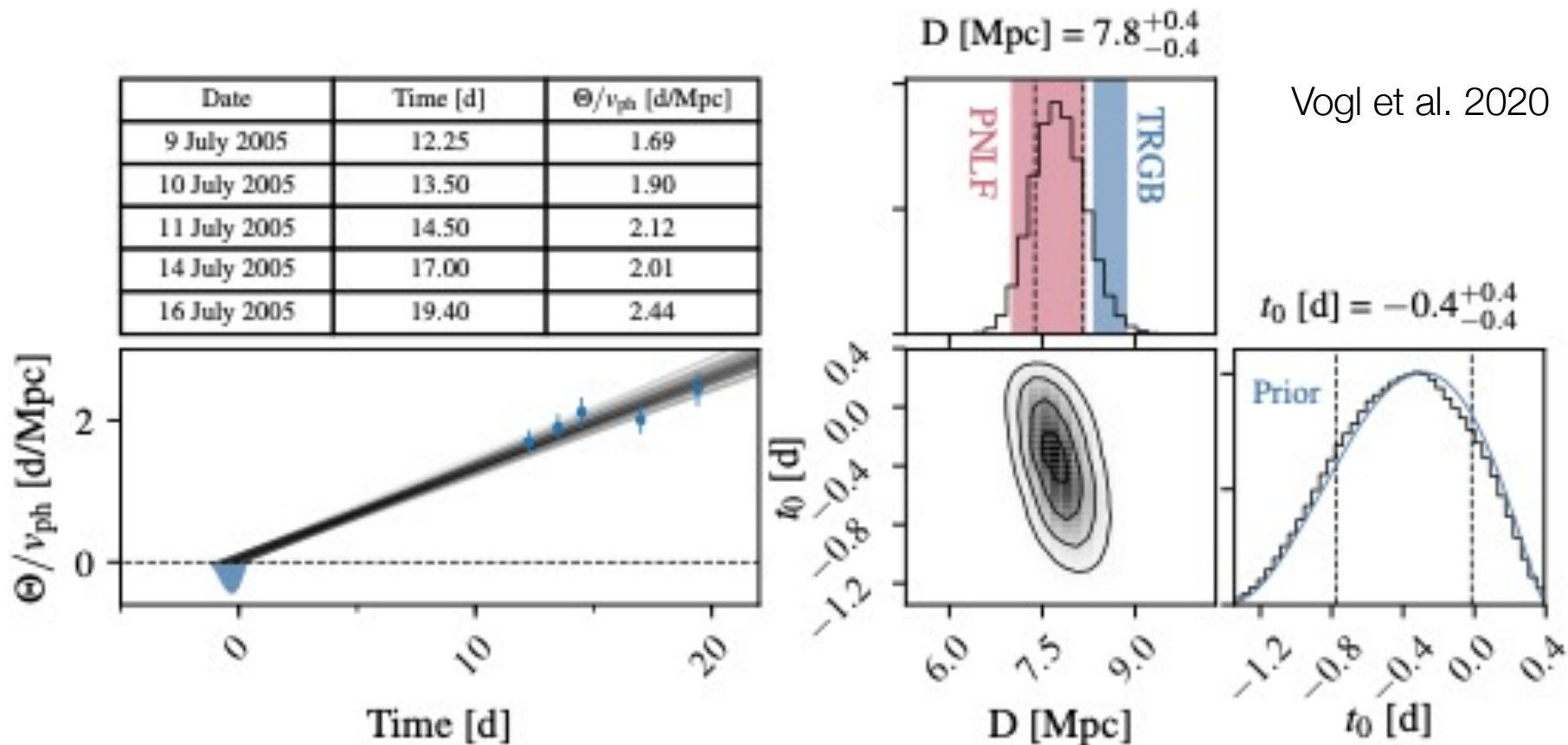


(e) 16 July 2005

Vogl et al. 2020

# Distance Determination

Slope is inverse distance:  $\frac{\Theta}{v} = \frac{1}{D_A} (t - t_0)$





# adH0cc

“accurate determination of  $H_0$  with core-collapse supernovae”

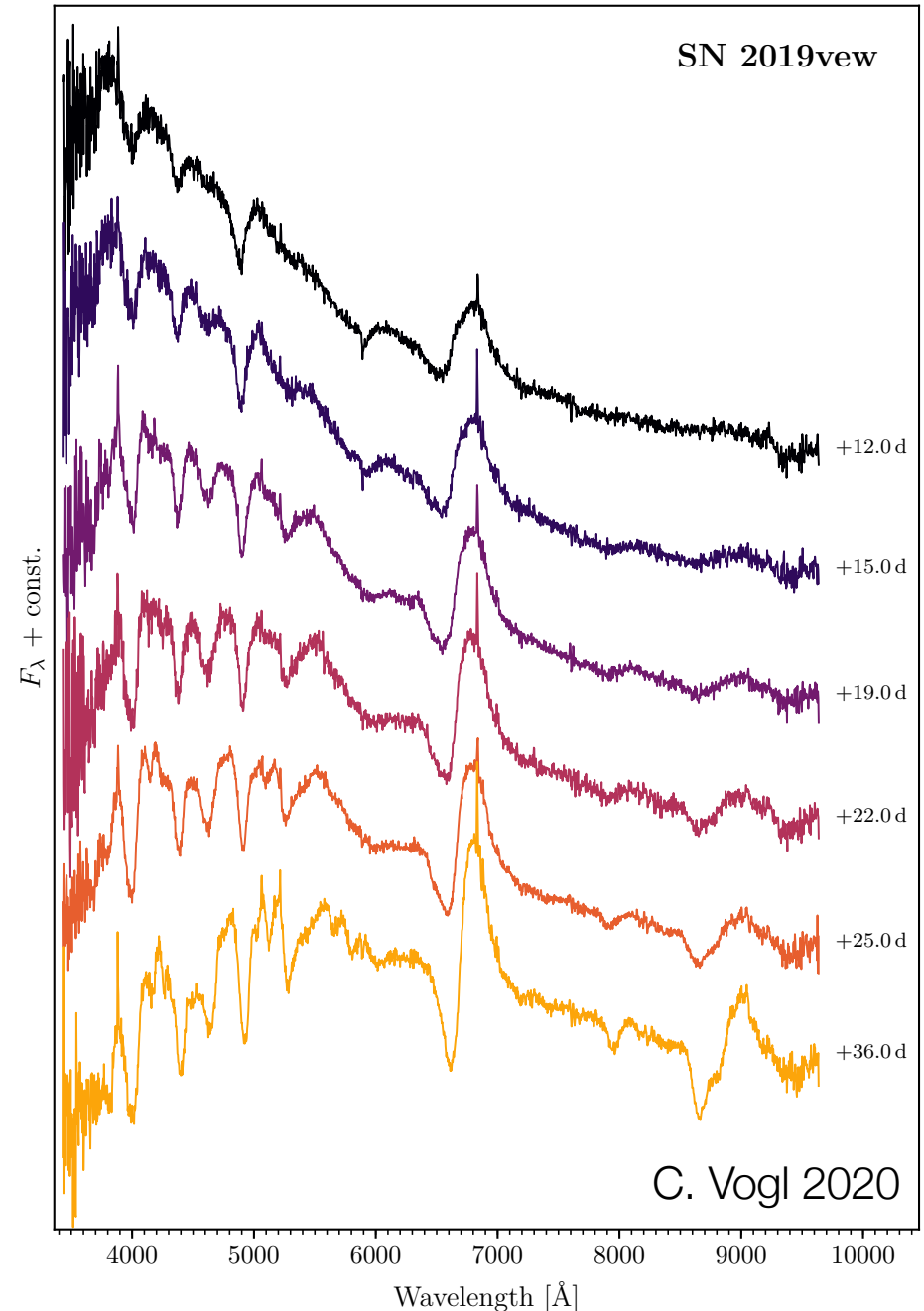
- Use the Expanding Photosphere Method to ~30 Type II supernovae in the Hubble flow ( $0.03 < z < 0.1$ )
- Independent of distance ladder
  - no parallaxes, no Cepheids, no Type Ia supernovae
- FORS2 Large Programme over 3 semesters
  - 6 epochs spectroscopy and photometry per supernova
  - 8 SNe followed in first semester (P104)
  - currently on hold
- SNFactory data
  - about 15 SNe with  $0.01 < z < 0.05$



# adH0cc

## Critical observables

- time of explosion
- spectral coverage
  - before max until well into the plateau
- photometry
  - simulatenously to spectroscopy



# Conclusions

Hubble constant sets absolute scale  
(and age) of the universe

- Past conflicts resolved
  - Age of Universe is bigger than age of the Earth
    - recognition of different stellar populations
  - Age of Universe bigger than oldest stars
    - cosmological constant



# Conclusions

Current discrepancy of  $\sim 4.4\sigma$  between

- $H_0$  measured locally (distance ladder) and
- $H_0$  measured at  $z=1100$  (CMB)

Significance?

- systematics based on Cepheid calibration

Extreme accuracy required

- e.g. Cepheid parallax zeropoint

Independent measurement needed

- Expanding Photosphere Method